

# An introduction to NMR Hyperpolarization

## The struggle towards higher NMR sensitivity

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## How to increase sensitivity?



Physics

Chemistry

Engineering

## NMR sensitivity

Most important: signal to noise ratio (SNR)

Other factors:

- high homogeneity (e.g. for spectral or spatial resolution)  
high gradient strengths (e.g. in MRI)
- Noise reduction: (e.g. Johnson noise)
- sensitive coils ( $B_1$  field)

Time

$$SNR = \frac{S}{N} \sim \sqrt{t}$$

Is there another way to get more signal?

## How to increase an NMR signal:

### Spin density

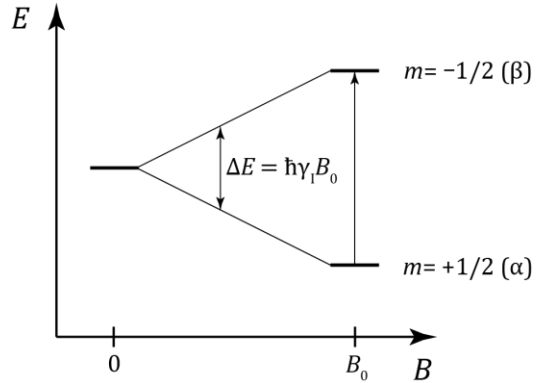
Concentration  
Abundance

### Polarization

Gyromagnetic ratio  
Temperature  
Magnetic field

### Hyperpolarization

DNP  
SEOP  
PHIP

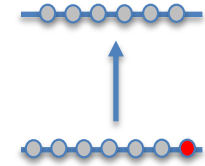
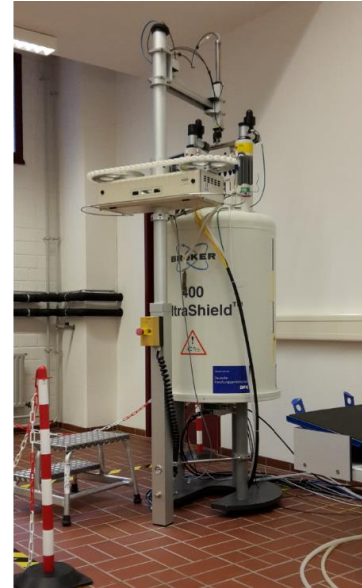


## Zeeman splitting

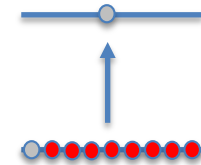
at thermal equilibrium Boltzmann distribution

Energy difference between the two spin states:  $1.6 \times 10^{-4} \text{ kJ mol}^{-1}$

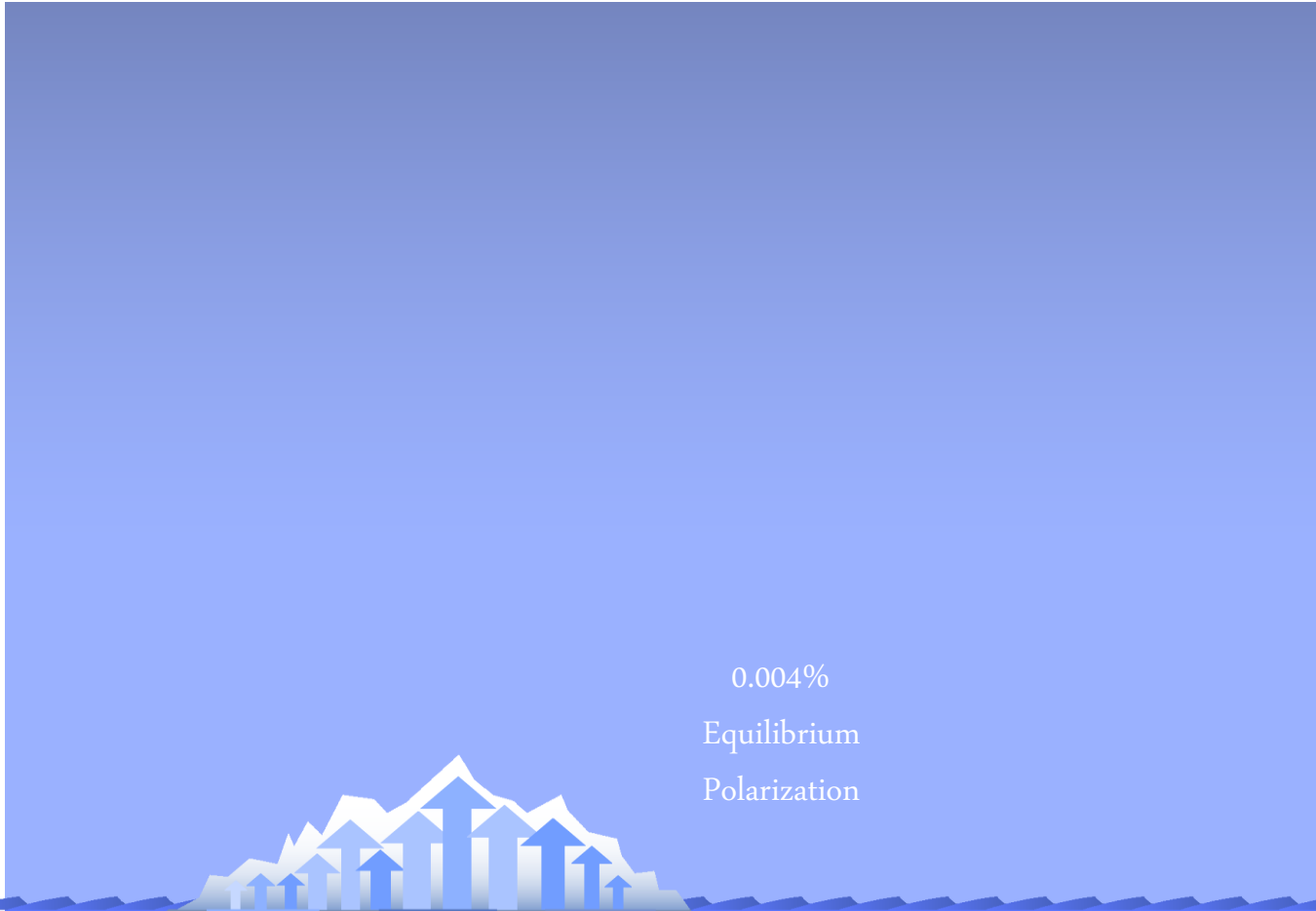
→ population difference of 1 : 31000 for 9.4 T at 300K



Standard NMR



Hyperpolarization



## NMR signal

- In the simple case of  $I=1/2$  the resulting magnetization is

$$M = \frac{\hbar}{2} N \gamma P \qquad P = \frac{N_+ - N_-}{N_+ + N_-}$$

- Ratios between populations of successive magnetic energy levels

$$\frac{N_m}{N_{m+1}} = \exp\left(\frac{\gamma \hbar B_0}{kT}\right)$$

- At high temperatures ( $T > 1K$ ) and average fields (1-10 Tesla):

$$M = \frac{N \gamma \hbar}{2} \left( \frac{1 + \frac{\gamma \hbar B_0}{2kT} - 1 + \frac{\gamma \hbar B_0}{2kT}}{1 + \frac{\gamma \hbar B_0}{2kT} + 1 - \frac{\gamma \hbar B_0}{2kT}} \right) = \frac{N}{4} \gamma^2 \hbar^2 \frac{B_0}{kT}$$

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# NMR Hyperpolarization

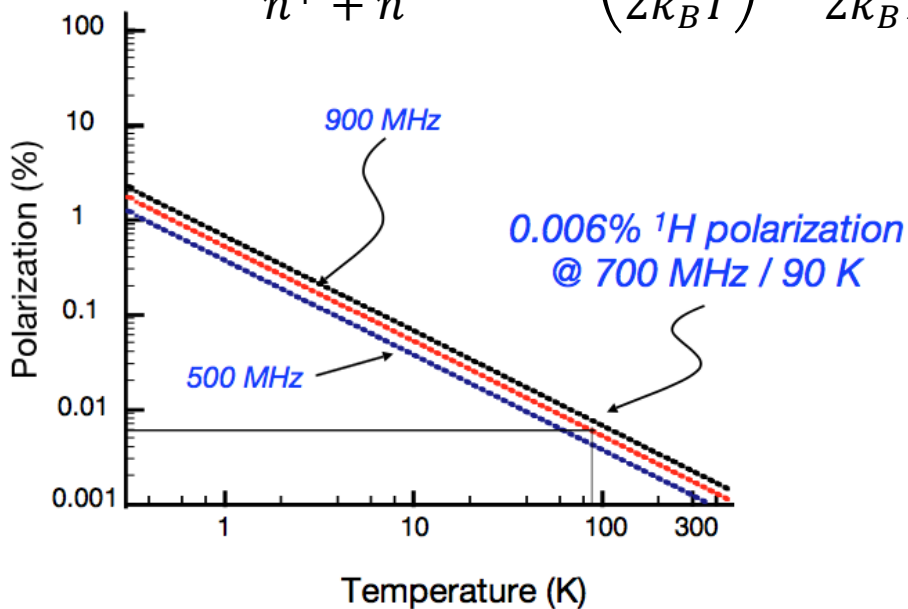
- Brute Force NMR
  - Very low temperatures
  - Very high fields
- Dynamic nuclear polarisation (DNP)
- Spin-exchange optical pumping (SEOP)
- Parahydrogen induced polarisation (PHIP)



## Nuclear Spin Polarization

### Temperature and Field Dependence

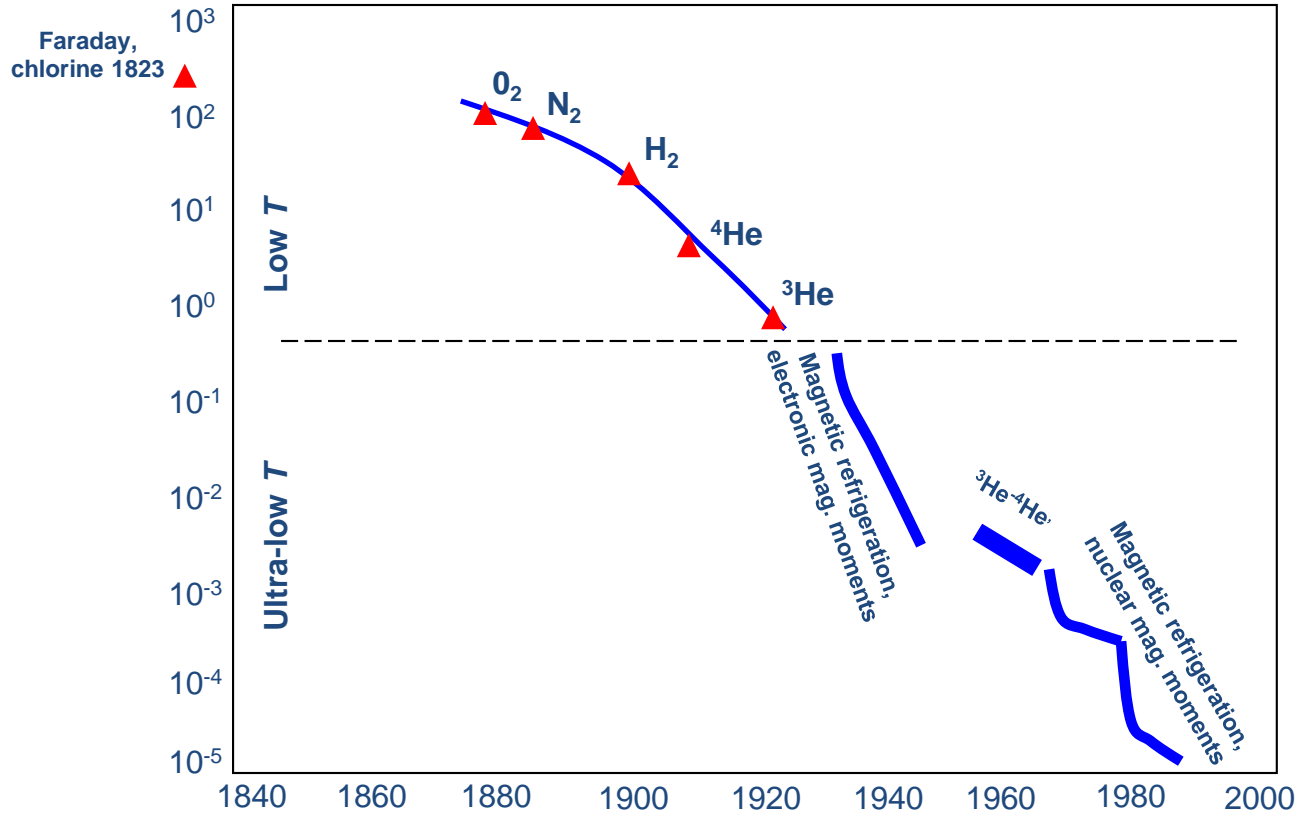
$$P = \frac{n^+ - n^-}{n^+ + n^-} = \tanh\left(\frac{\gamma\hbar B_0}{2k_B T}\right) \approx \frac{\gamma\hbar B_0}{2k_B T}$$



| Substance        | boiling T<br>(P=1 bar) | melting T<br>(P=1 bar) | Latent heat<br>kJ/liter | Price<br>\$/ liter |
|------------------|------------------------|------------------------|-------------------------|--------------------|
| H <sub>2</sub> O | 373.15                 | 273.15                 | 2252                    |                    |
| Xe               | 165.1                  | 161.3                  | 303                     |                    |
| O <sub>2</sub>   | 90.2                   | 54.4                   | 245                     |                    |
| N <sub>2</sub>   | 77.4                   | 63.3                   | 160                     | 0.3                |
| H <sub>2</sub>   | 20.3                   | 14.0                   | 31.8                    |                    |
| <sup>4</sup> He  | 4.21                   | --                     | 2.56                    | 4                  |
| <sup>3</sup> He  | 3.19                   | --                     | 0.48                    | 5x10 <sup>4</sup>  |

the cooling power  
diminishes rapidly  
with decreasing  $T$

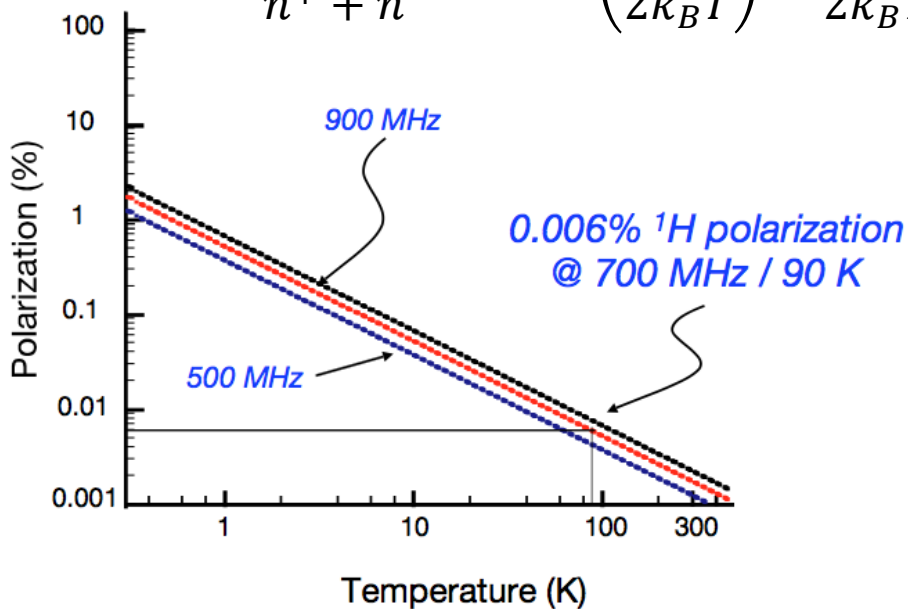
|                      |                  |                  |                  |                  |      |      |      |
|----------------------|------------------|------------------|------------------|------------------|------|------|------|
| $P$ , torr           | 10 <sup>-4</sup> | 10 <sup>-3</sup> | 10 <sup>-2</sup> | 10 <sup>-1</sup> | 1    | 10   | 100  |
| $T(^4\text{He})$ , K | 0.56             | 0.66             | 0.79             | 0.98             | 1.27 | 1.74 | 2.64 |
| $T(^3\text{He})$ , K | 0.23             | 0.28             | 0.36             | 0.47             | 0.66 | 1.03 | 1.79 |



## Nuclear Spin Polarization

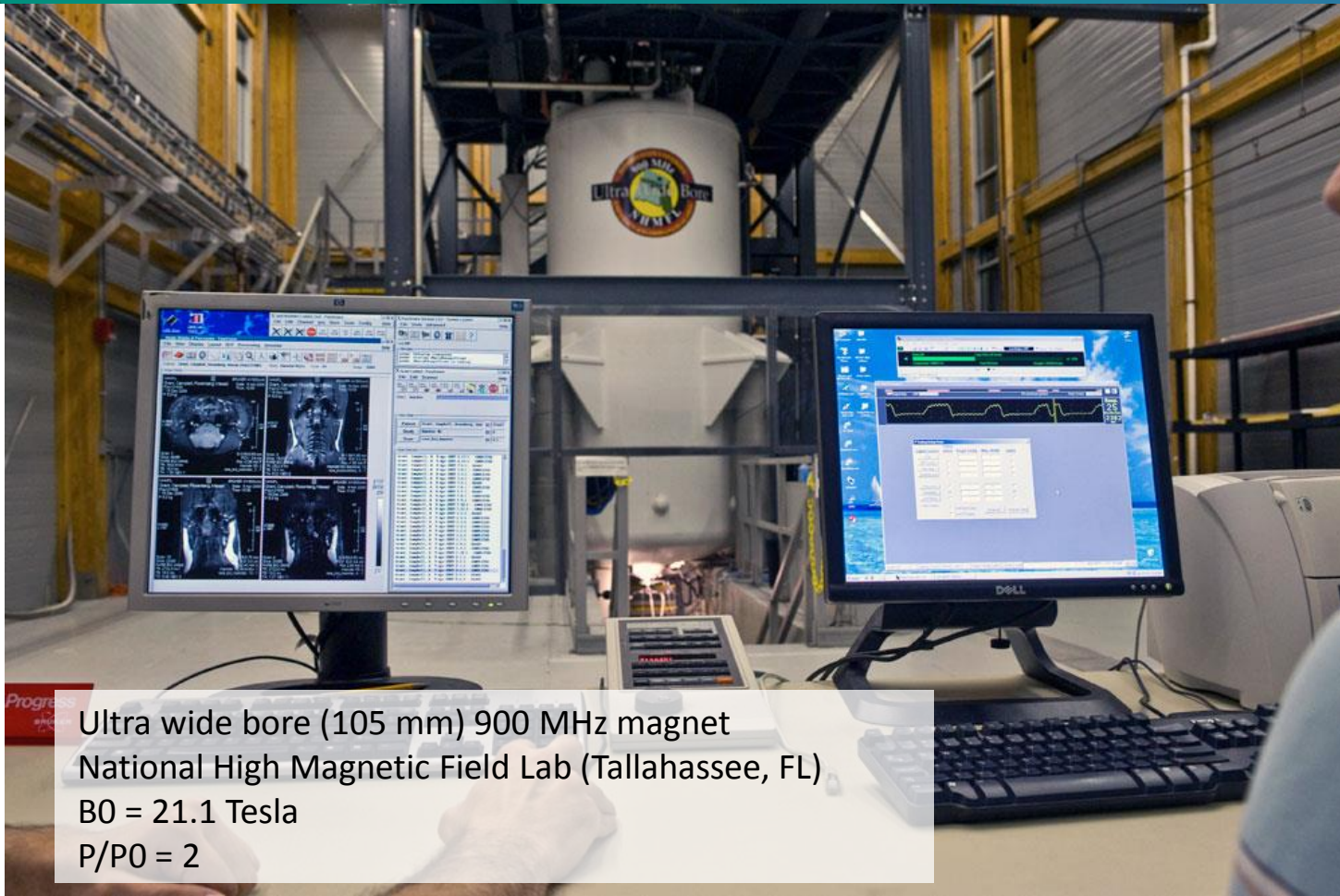
### Temperature and Field Dependence

$$P = \frac{n^+ - n^-}{n^+ + n^-} = \tanh\left(\frac{\gamma\hbar B_0}{2k_B T}\right) \approx \frac{\gamma\hbar B_0}{2k_B T}$$





Bruker Aeon 1 GHz magnet  
Installed at the University of Bayreuth  
 $B_0 = 23.5$  Tesla  
 $P/P_0 = 2$



Ultra wide bore (105 mm) 900 MHz magnet  
National High Magnetic Field Lab (Tallahassee, FL)  
 $B_0 = 21.1$  Tesla  
 $P/P_0 = 2$

|                   | <b>Series<br/>Connected<br/>Hybrid Magnet</b> | <b>45 Tesla Hybrid<br/>Magnet</b> | <b>900 MHz Ultra-<br/>Wide Bore NMR<br/>Magnet</b> |
|-------------------|-----------------------------------------------|-----------------------------------|----------------------------------------------------|
| Frequency         | 1.5 GHz                                       | n/a                               | 900 MHz                                            |
| Field strength    | 36 tesla                                      | 45 tesla                          | 21.1 tesla                                         |
| Bore size         | 40 mm                                         | 32 mm                             | 105 mm                                             |
| Field Homogeneity | 1 part per million                            | 150 part per million              | 0.001 part per million                             |
| Year completed    | 2016                                          | 1999                              | 2004                                               |
| Cost              | \$18.7 million                                | \$14.4 million                    | \$16 million                                       |



|                               | <b>Series Connected Hybrid Magnet</b>                        | <b>45 Tesla Hybrid Magnet</b>                                  | <b>900 MHz Ultra-Wide Bore NMR Magnet</b> |
|-------------------------------|--------------------------------------------------------------|----------------------------------------------------------------|-------------------------------------------|
| Height                        | 3.6 meters (11.8 feet)                                       | 6.7 meters (22 feet)                                           | About 5 meters (16 feet)                  |
| Weight                        | 30,000 kg (33 tons)                                          | 31,752 kg (35 tons)                                            | 13,600 kg (15 tons)                       |
| Type of magnet                | Resistive-superconducting hybrid (coils connected in series) | Resistive-superconducting hybrid (coils connected in parallel) | Superconducting                           |
| Cost to operate at full field | \$678 /hour                                                  | \$1,452 /hour                                                  | \$33 /hour                                |
| Power required to operate     | 14 Megawatts                                                 | 30 Megawatts                                                   | 0 Megawatts                               |

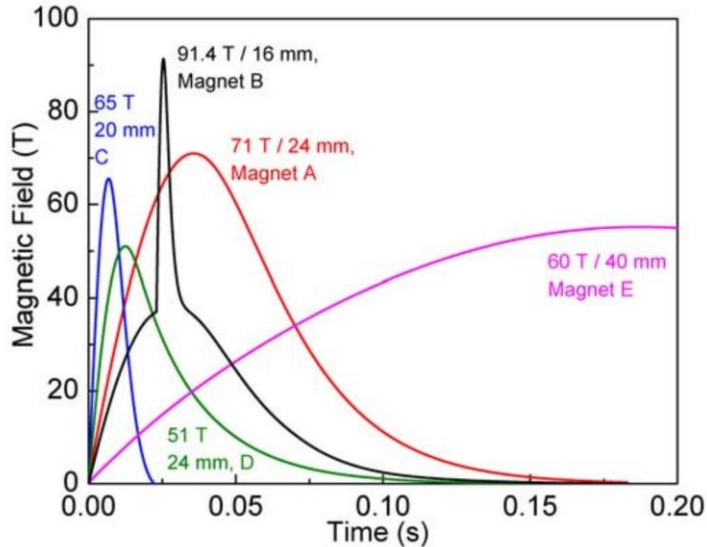
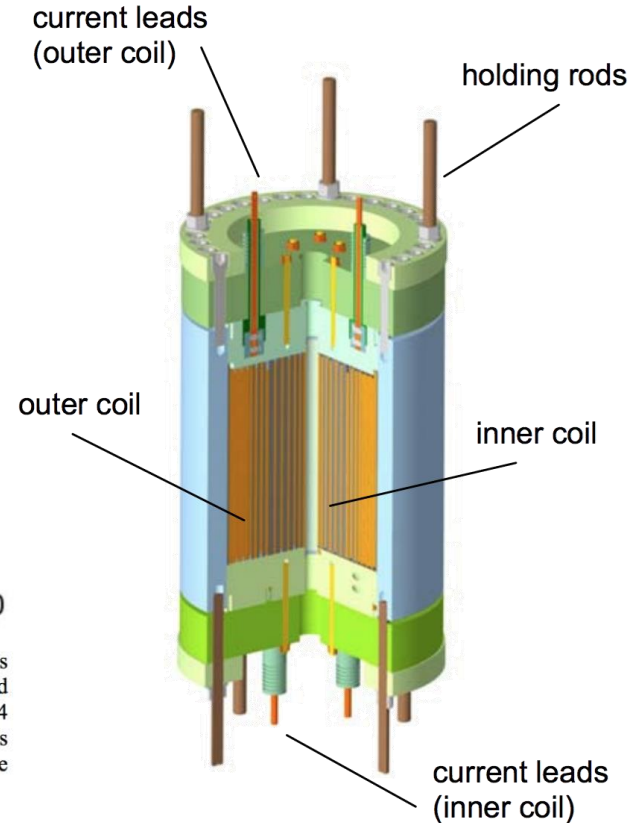
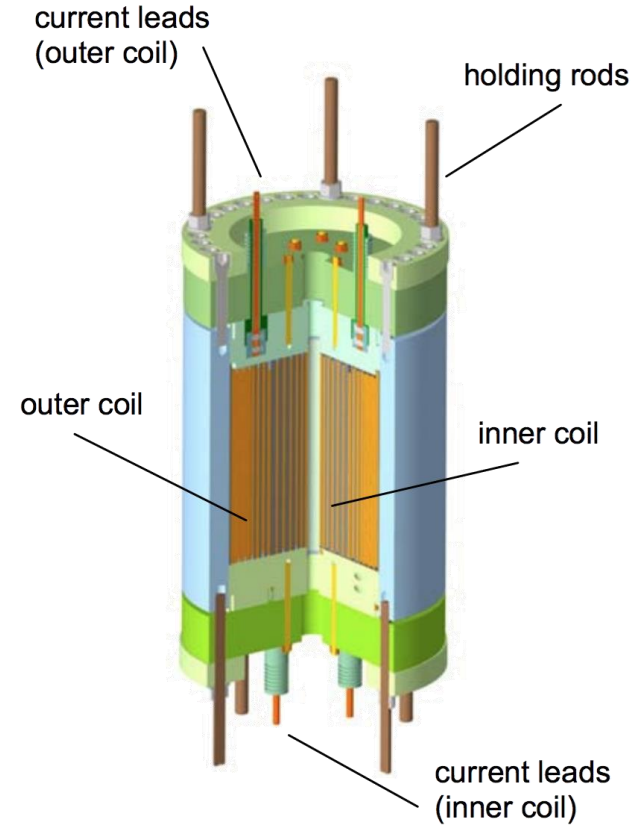
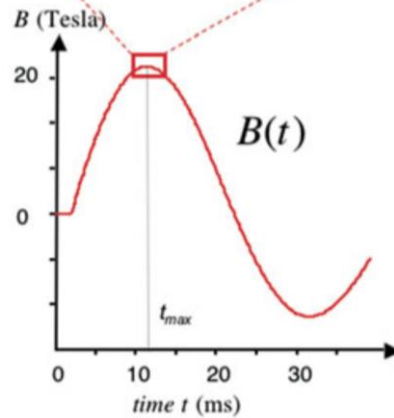
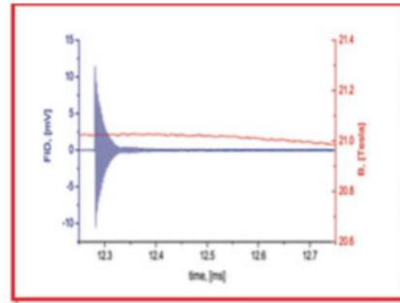


Fig. 2. (Color online) Measured time dependences of the magnetic fields obtained with various pulsed magnets operational at the HLD. Magnets A and D are 8.5 and 1.5 MJ magnets, respectively. Both of them have bores of 24 mm. Magnet B is a dual-coil 9.5 MJ magnet with a 16 mm bore. Magnet C is a 1 MJ pulsed magnet with a 20 mm bore. Magnet E is a 43 MJ long-pulse magnet with a 40 mm bore.





## Is Brute Force Convenient?

Average enhancements (ref: H = 11.7 Tesla, T = 300 K) :

|                                       |                       |
|---------------------------------------|-----------------------|
| High fields (100 T), room temp:       | $P/P_0 = 10$          |
| Standard fields and low temp (10 mK): | $P/P_0 = 10^4 - 10^5$ |
| High fields, low temp:                | $P/P_0 = 10^5$        |

### Pros:

- High enhancements at extremely low temperatures
- In situ polarization, independent from sample preparation
- Sample is not contaminated by foreign agents

### Cons:

- Technically demanding
- Slow relaxation times at low temperatures / high fields
- Long measurement times
- Not necessarily safe for the sample



## NMR Hyperpolarisation

- Brute Force NMR
  - Very low temperatures
  - Very high fields
- **Dynamic nuclear polarisation (DNP)**
- Spin-exchange optical pumping (SEOP)
- Parahydrogen induced polarisation (PHIP)

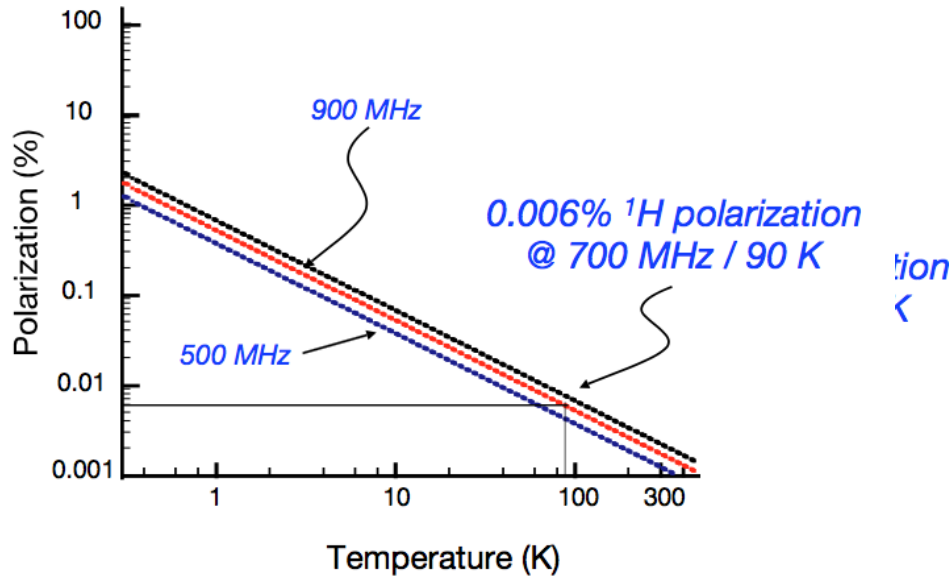


## Electron and Nuclear Spin Polarization

### Temperature and Field Dependence

$$P = \frac{n^+ - n^-}{n^+ + n^-} = \tanh\left(\frac{\gamma\hbar B_0}{2k_B T}\right) \approx \frac{\gamma\hbar B_0}{2k_B T}$$

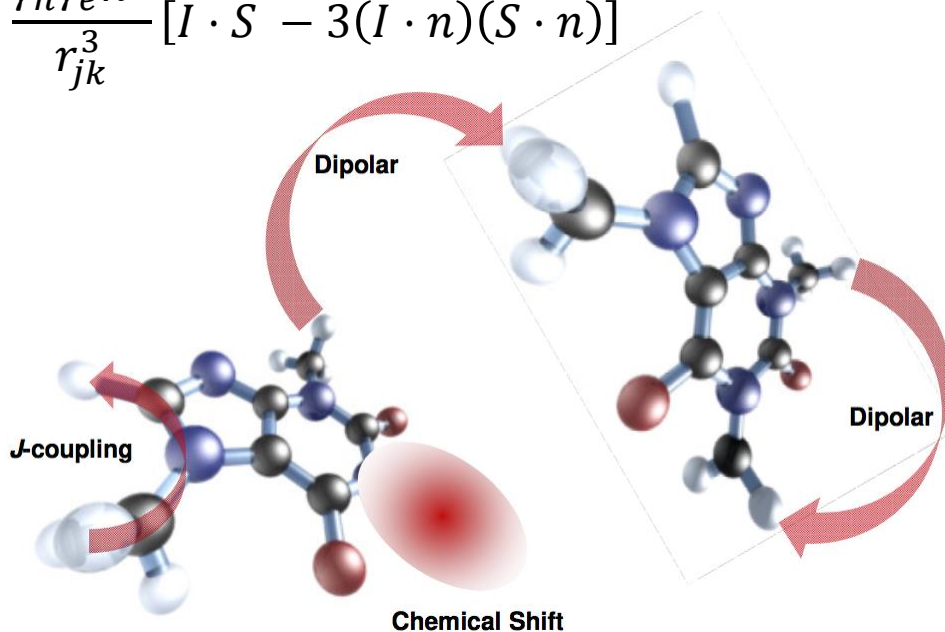
$$\frac{\gamma_e}{\gamma^1_H} \sim 660$$



## Dipolar couplings

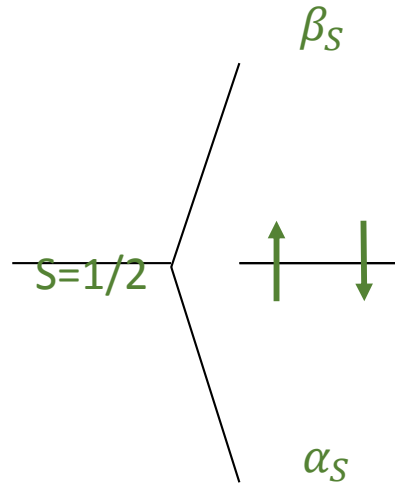
$$H = \omega_e S_Z - \omega_n I_Z + H_{ee} + H_{en} + H_{nn}$$

$$H_{en}^D = \frac{\gamma_n \gamma_e \hbar^2}{r_{jk}^3} [I \cdot S - 3(I \cdot n)(S \cdot n)]$$

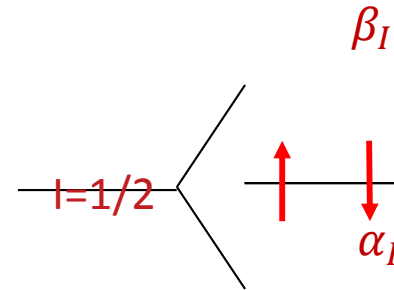


In liquids we don't see an effect on the spectrum due to tumbling motions  
**BUT it affects relaxation**

Electron



Proton



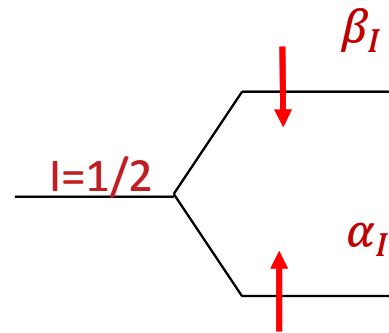
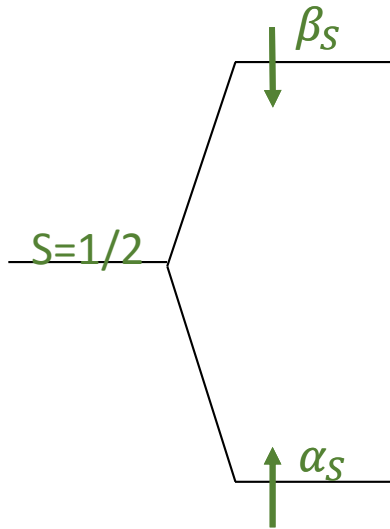
$H \gg 0$



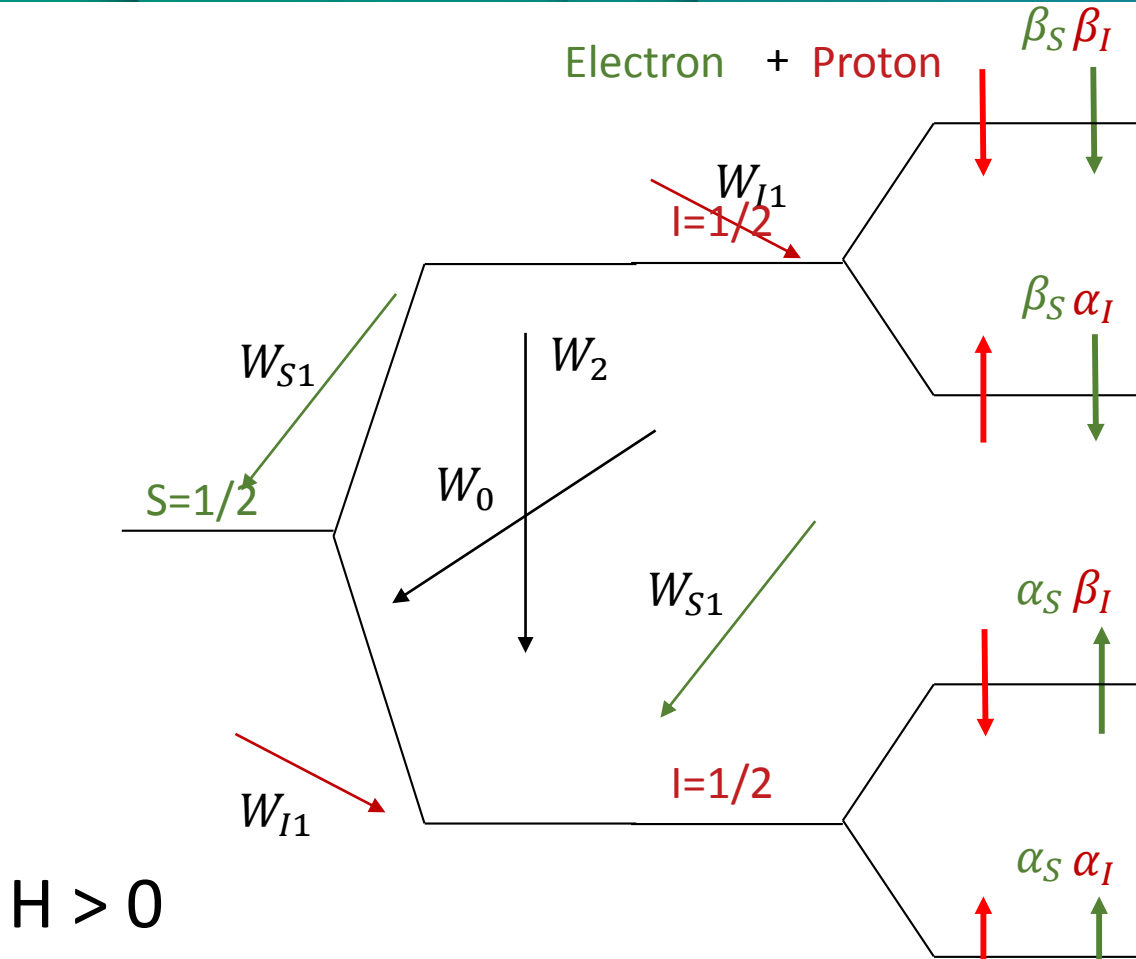
Electron

+

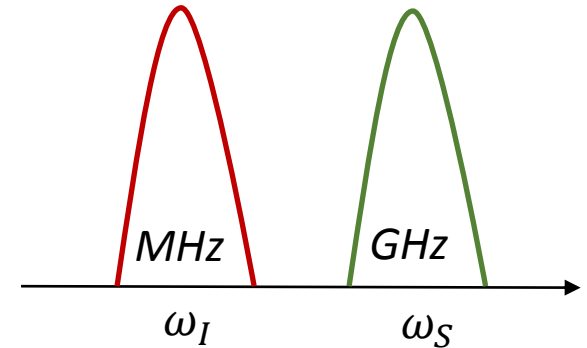
Proton

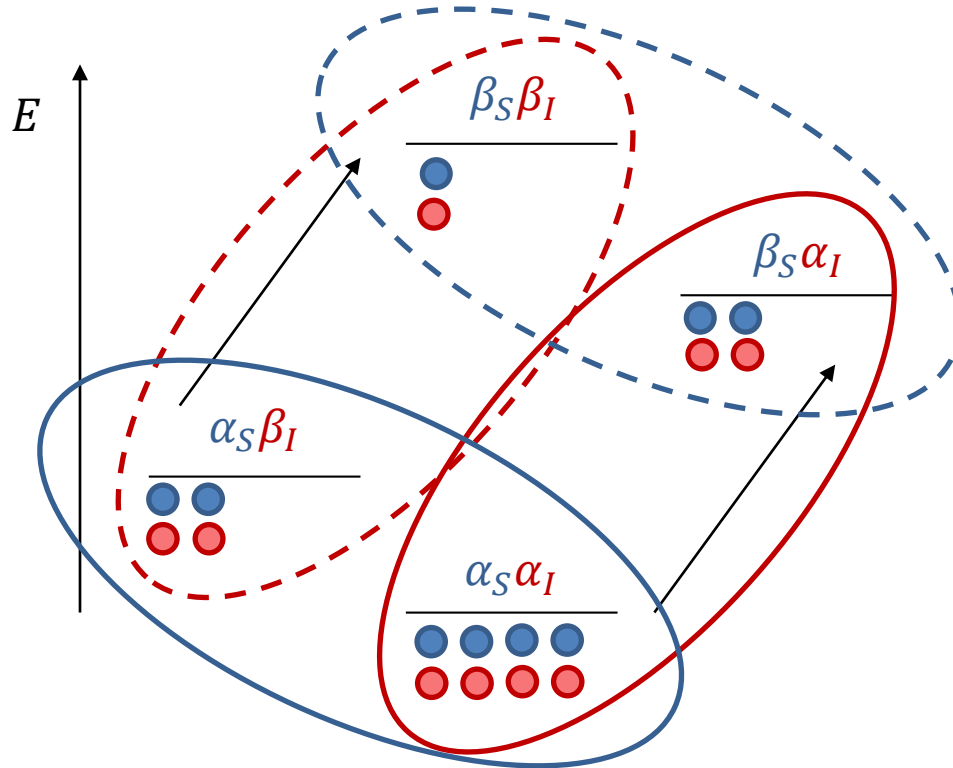


$H > 0$



$W_0$  and  $W_2$  are allowed by the presence of e-H dipolar couplings but they don't show up in the NMR spectrum



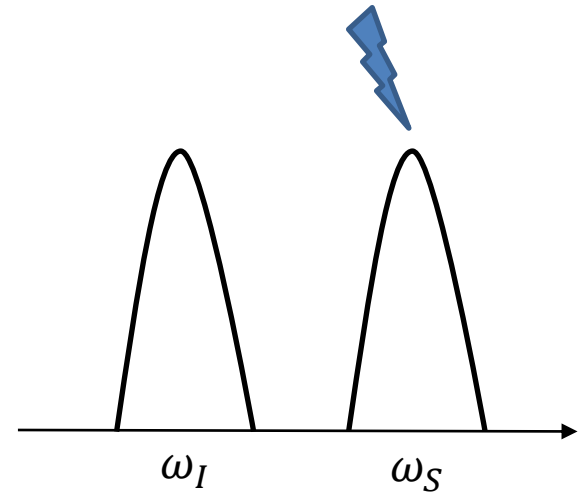


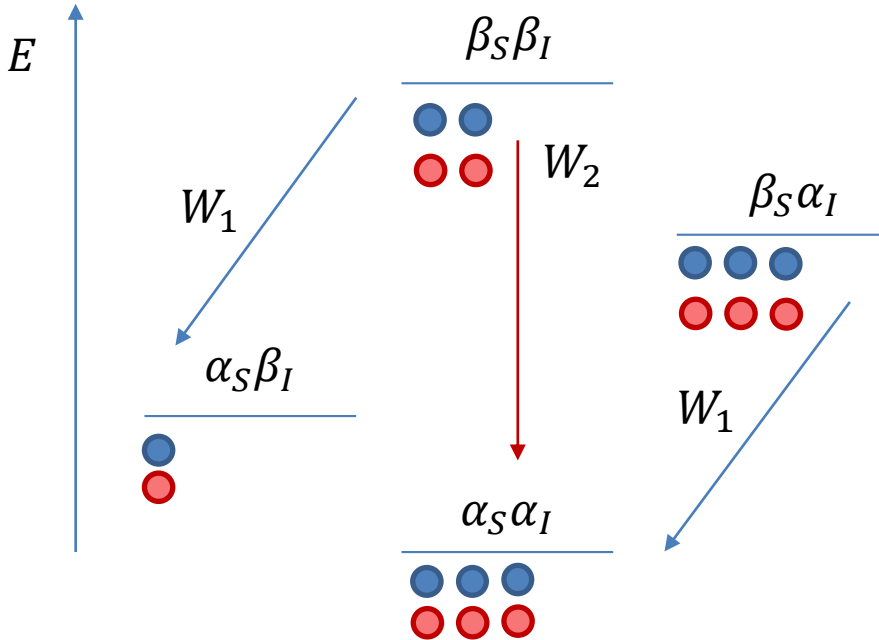
Populations

$$N_i = N_{i\alpha} - N_{i\beta}$$

$$N_S = 3 \cdot 1$$

$$N_I = 3$$



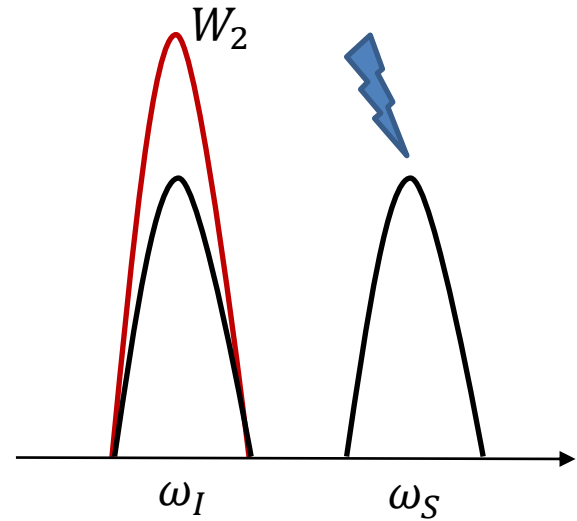


Populations

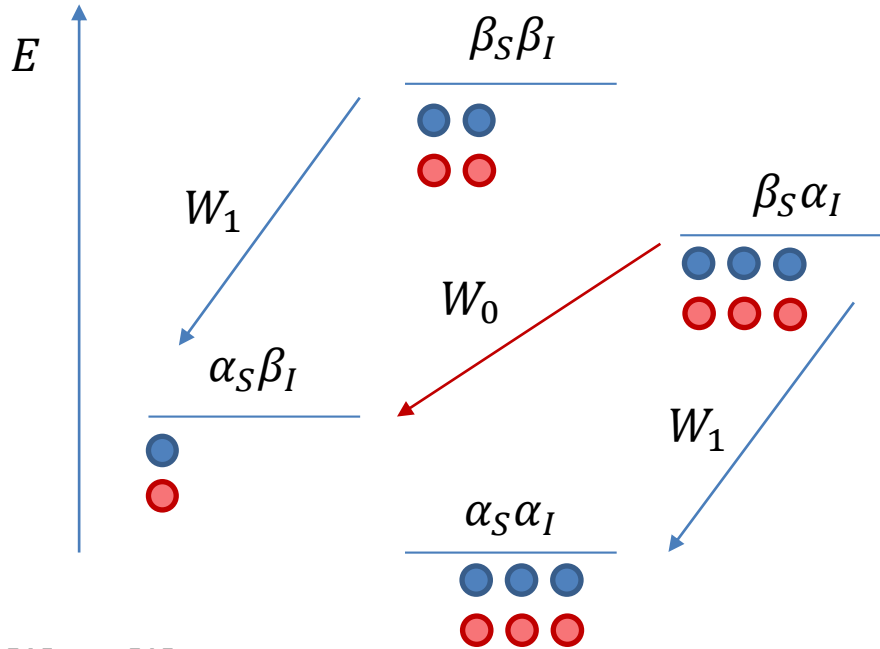
$$N_i = N_{i\alpha} - N_{i\beta}$$

$$N_S = \pm 1$$

$$N_I = 3$$



$W_2$  dominates for fast tumbling molecules

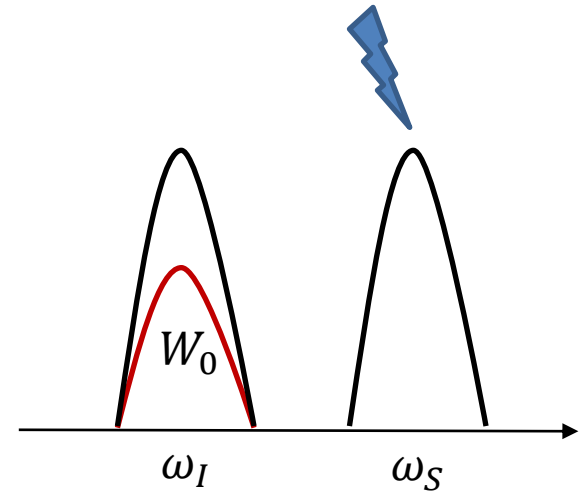


Populations

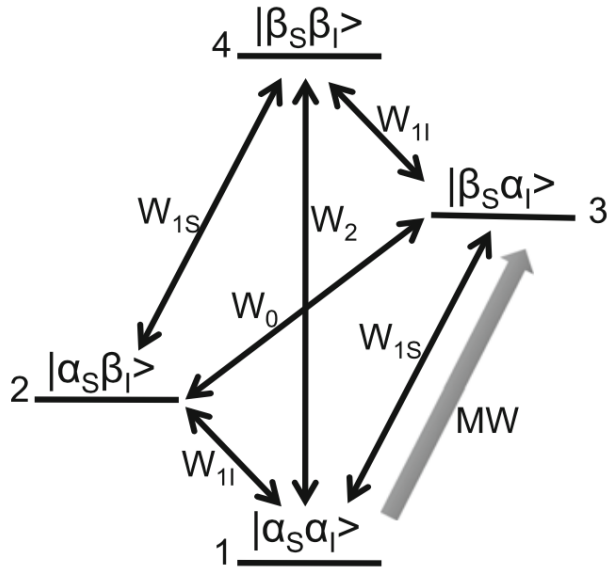
$$N_i = N_{i\alpha} - N_{i\beta}$$

$$N_S = \pm 1$$

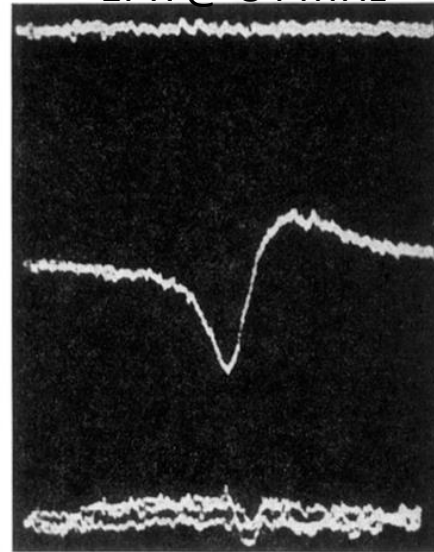
$$N_I = \mathbf{3}$$



$$\frac{P_I}{P_{I0}} = 1 + \frac{W_2 - W_0}{2W_1 + W_2 + W_0} \cdot \frac{\gamma_S}{\gamma_I} \cdot \chi_{DD}$$



$^7\text{Li}$  NMR @ 50kHz  
EPR @ 84 MHz



$^7\text{Li}$  w/o DNP

$^7\text{Li}$  with DNP  
@ 84 MHz

$^1\text{H}$  glycerol

Carver and Slichter, *Phys. Rev.* 92, 212-213 (1953)  
*Phys. Rev.* 102, 975-980 (1956)

$\epsilon \sim 100$

DNP can be classified by the experimental arrangement that is used to carry out the experiment.

## **Overhauser DNP (O-DNP)**

Room temperature

Driven by the Overhauser effect (OE), where the entire process is carried out in the liquid state.

## **Solid-state DNP (SS-DNP)**

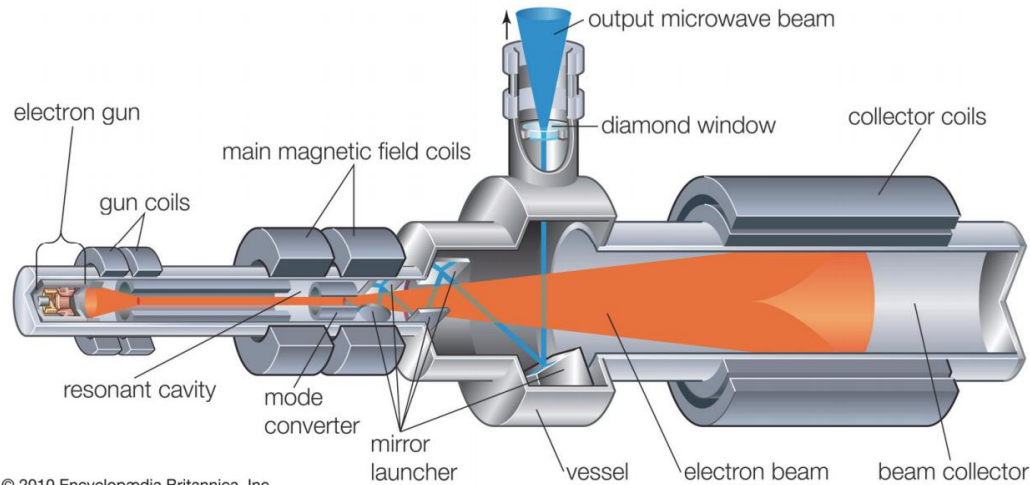
Low temperatures  
(typically 90 K)

Driven by the cross effect (CE), using high-power microwave sources in conjunction with MAS for solid-state NMR.

## **Dissolution DNP (D-DNP)**

Low temperatures <1.5 K

Driven by thermal mixing (TM) or the solid effect (SE).

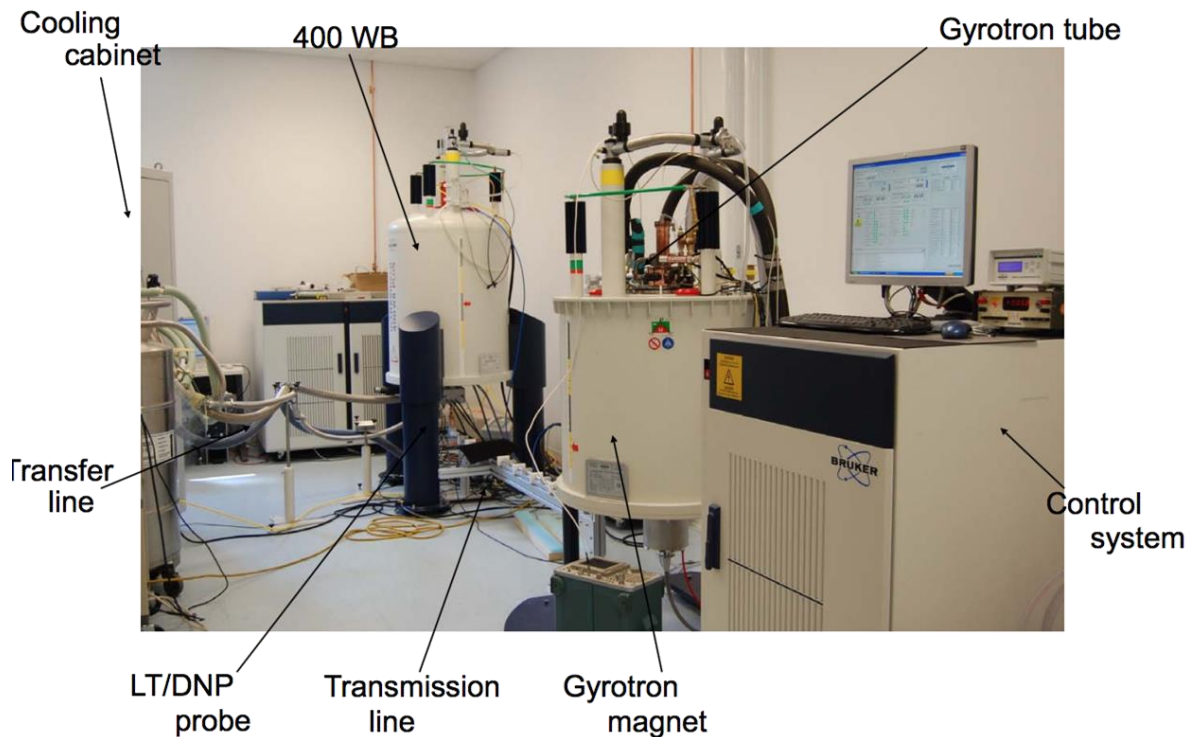


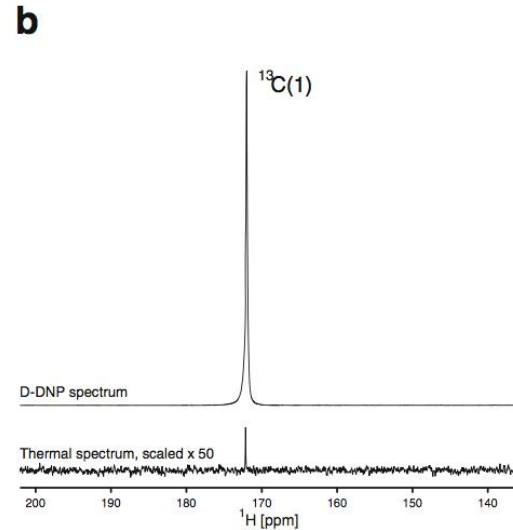
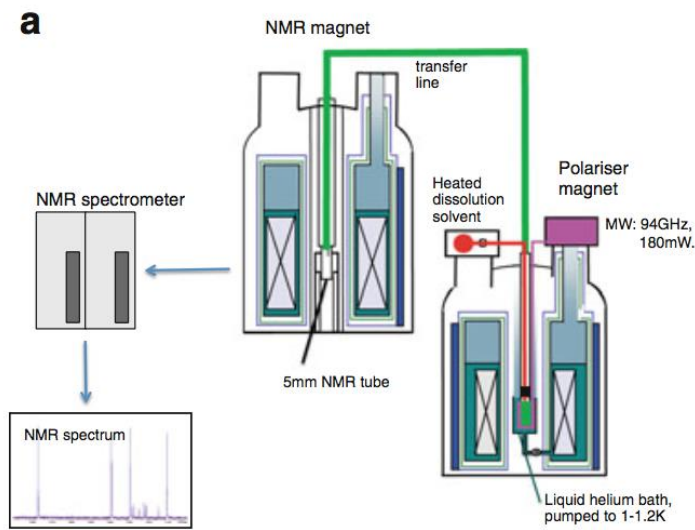
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A *gyrotron*, also known as *cyclotron resonance maser*, is a type of free electron maser. The device creates a strong beam of electrons with spiral trajectories moving in a strong magnetic field and passing through a resonant cavity. The bunching in this case occurs along the trajectory spiral. Gyrotron output power can reach many megawatts. The output frequency is easily tuned by changing the magnet field and the resonance frequency of the cavity.



## 263 GHz gyrotron in Bruker Billerica (MA) DNP Lab

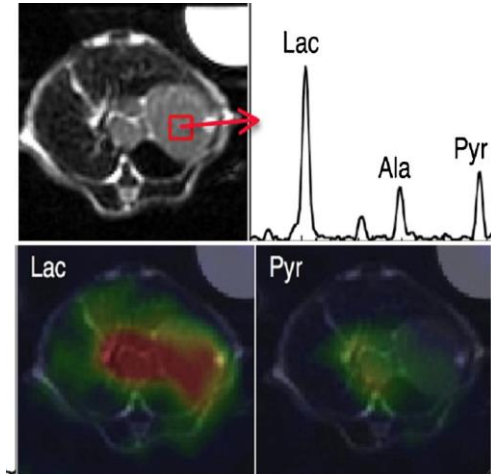




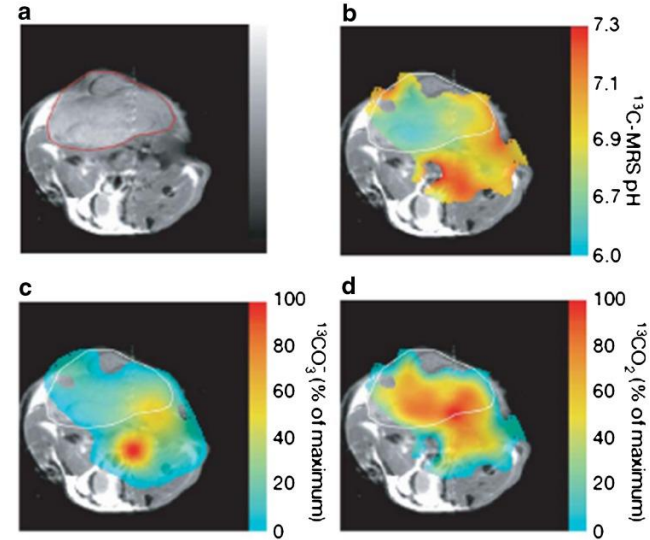
*Dissolution DNP (also ex-situ DNP), pioneered by Ardenkjær-Larsen, yields by far the largest overall enhancements of all implementations of DNP.*

$$\epsilon' = \epsilon \frac{B_{DNP}}{B_{NMR}} \cdot \frac{T_{NMR}}{T_{DNP}}$$

$$\sim 10^5$$



<sup>13</sup>C MRS of hyperpolarized pyruvate  
To observe tumor metabolites

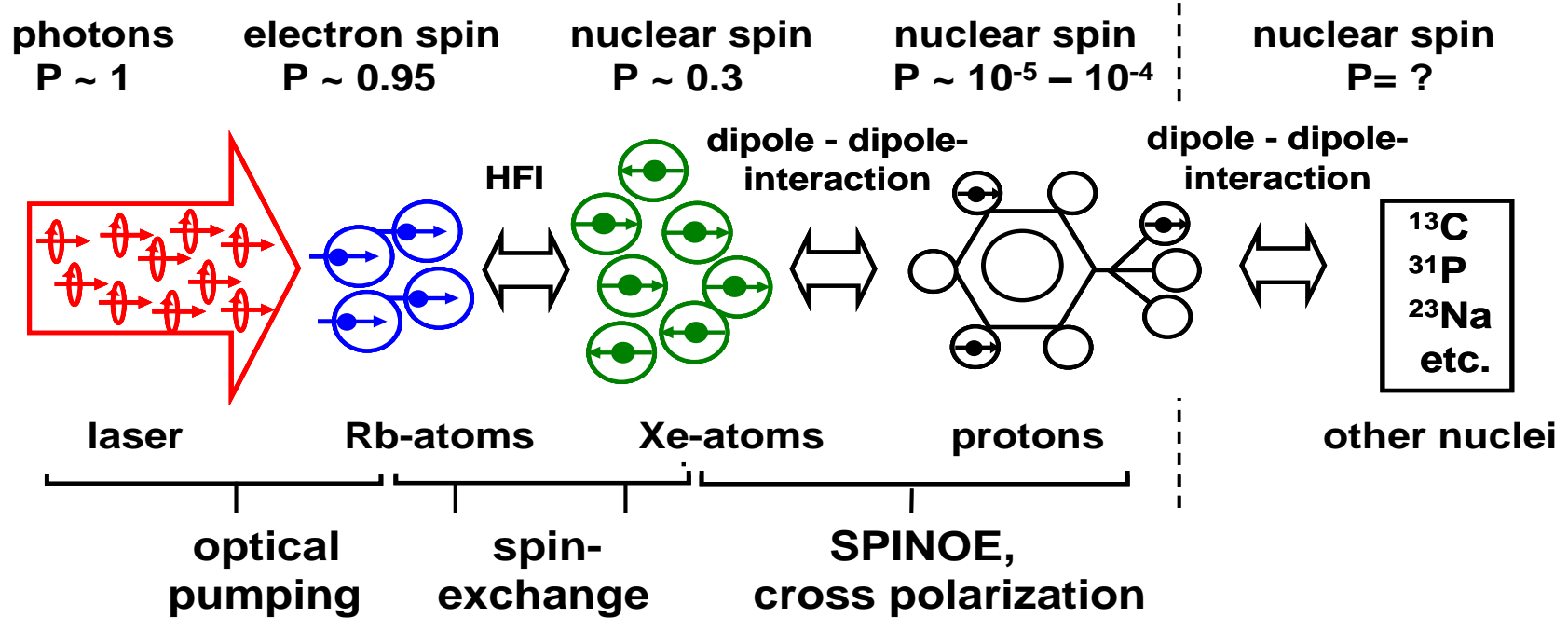


Imaging of tumor pH in-vivo  
with hyperpolarized <sup>13</sup>C bicarbonate

## NMR Hyperpolarisation

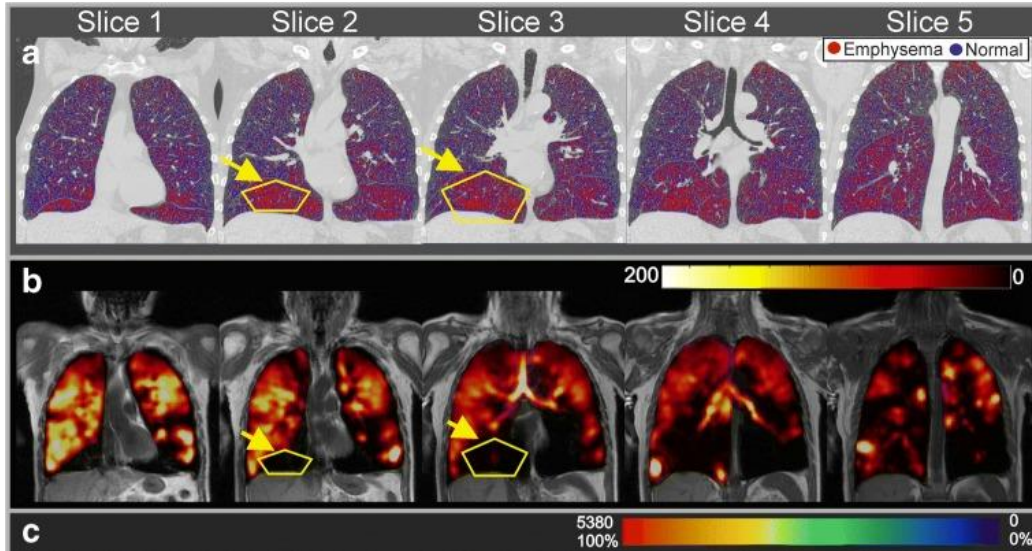
- Brute Force NMR
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## Why hyperpolarized $^{129}\text{Xe}$ ?

- $^{129}\text{Xe}$  is an excellent probe for non-invasive NMR measurements on materials
  - Zero chemical reactivity
  - Favourable absorption properties
  - Relatively high isotopic abundance (~26%)
  - Spin  $\frac{1}{2}$
  - High electronic polarizability, which means high chemical shifts dependent on shape, size and nature of occupied cavities.
  - High enhanced polarization (~85-90%).
  - Lack of NMR background



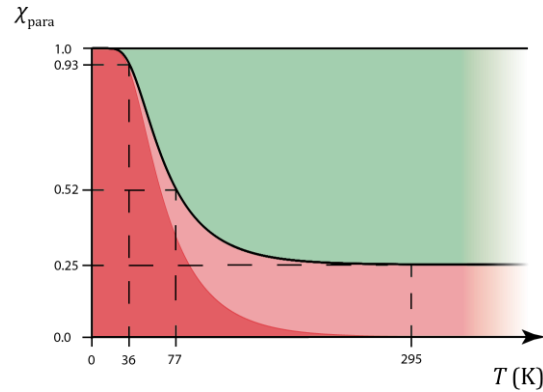
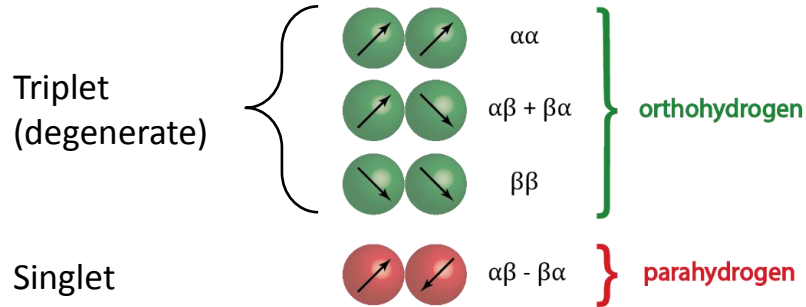
- a) CT and
  - b) combined  $^1\text{H}$  and hyperpolarized  $^{129}\text{Xe}$  MRI
- in a stage IV COPD patient

## NMR Hyperpolarisation

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- **Parahydrogen induced polarisation (PHIP)**



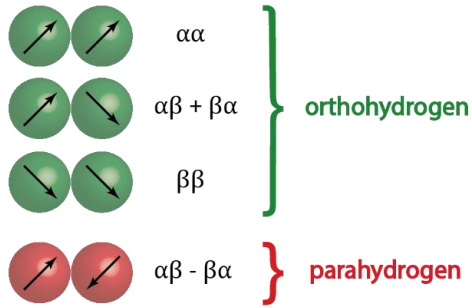




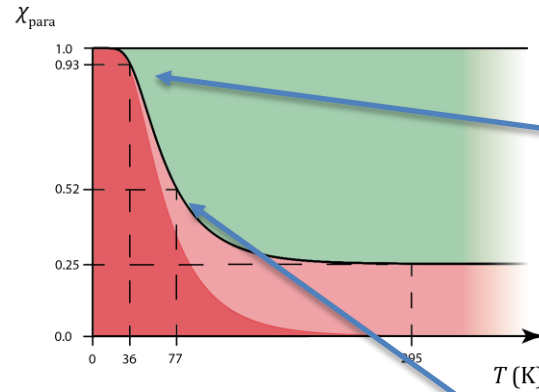
Energy difference: 1.5 kJ/mol (175 K or  $\sim 10^{12}$  Hz)

$$\Psi = \Psi_{\text{elec}} \Psi_{\text{vib}} \Psi_{\text{rot}} \Psi_{\text{trans}} \Psi_{\text{nuclearspin}}$$

But: Symmetry has to be broken to yield NMR signals



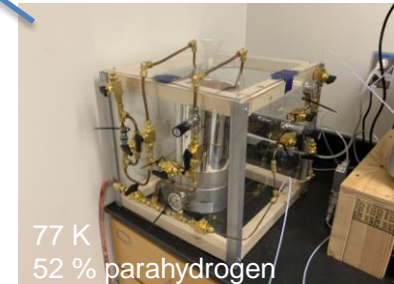
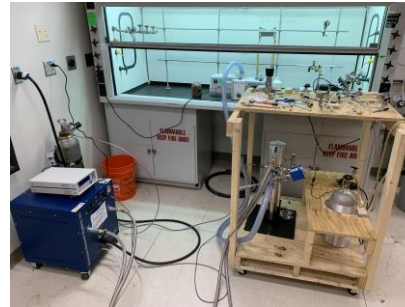
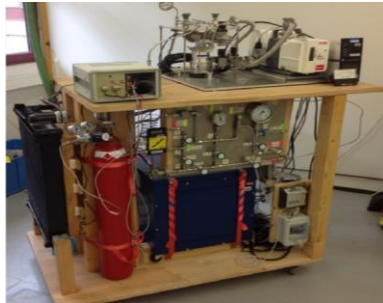
$$\Psi = \Psi_{\text{elec}} \Psi_{\text{vib}} \Psi_{\text{rot}} \Psi_{\text{trans}} \Psi_{\text{nuclearspin}}$$



Commercial Bruker generator

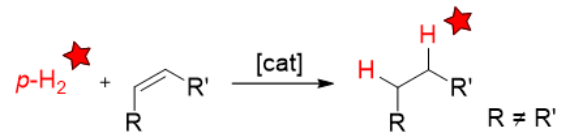


IMT @ KIT

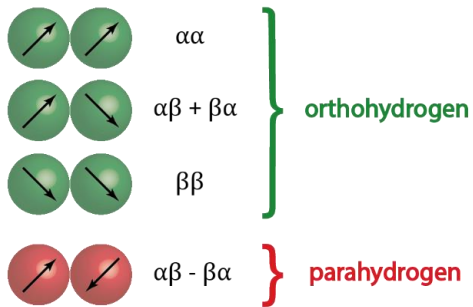
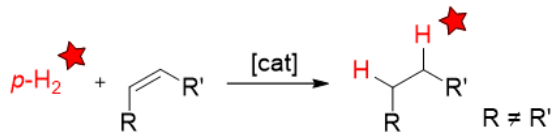


Custom made generators to achieve over 95% parahydrogen

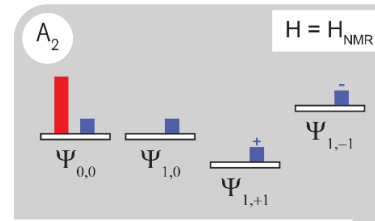
*a) Hydrogenation with  $p\text{-H}_2$*



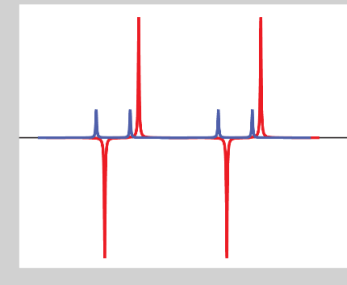
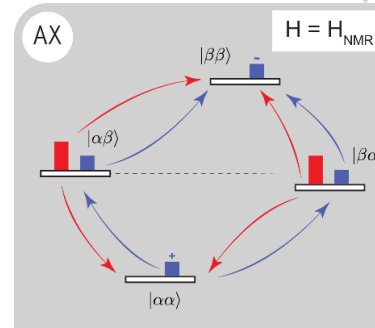
## a) Hydrogenation with $p\text{-H}_2$



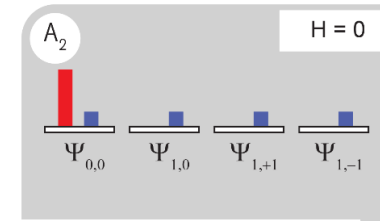
### PASADENA



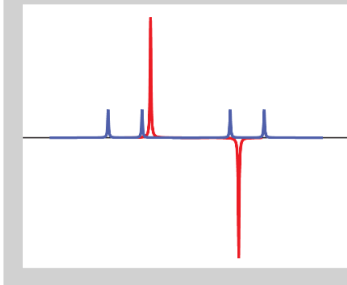
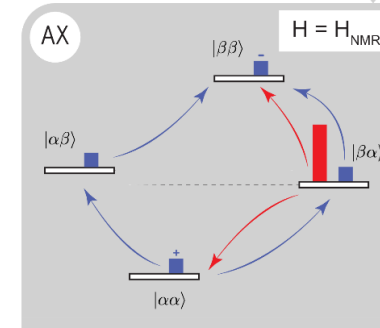
Hydrogenation



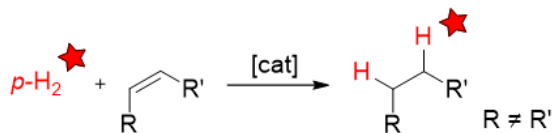
### ALTADENA



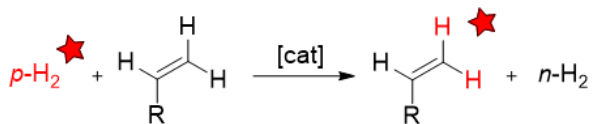
Hydrogenation and transport



## a) Hydrogenation with $p\text{-H}_2$



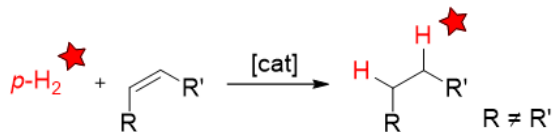
## b) Pairwise replacement with $p\text{-H}_2$



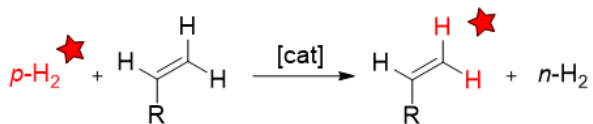
[a] C. R. Bowers and D. P. Weitekamp, *J. Am. Chem. Soc.* **1987**, 109, 5541-5542.

[b] A. Harthun, R. Giernoth, C. J. Elsevier and J. Bargon, *Chem. Commun.* **1996**, 2483.

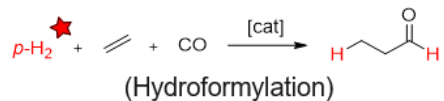
## a) Hydrogenation with $p\text{-H}_2$



## b) Pairwise replacement with $p\text{-H}_2$



## c) Single proton transfer

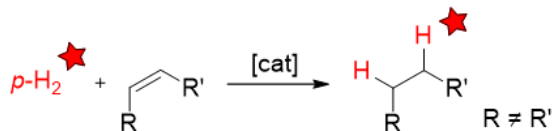


[a] C. R. Bowers and D. P. Weitekamp, *J. Am. Chem. Soc.* **1987**, 109, 5541-5542.

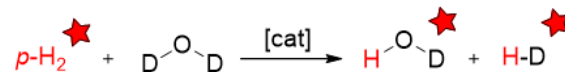
[b] A. Harthun, R. Giernoth, C. J. Elsevier and J. Bargon, *Chem. Commun.* **1996**, 2483.

[c] A. B. Permin and R. Eisenberg, *J. Am. Chem. Soc.*, **2002**, 124, 12406-12407.

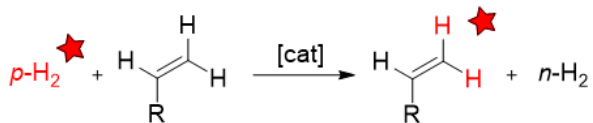
## a) Hydrogenation with $p\text{-H}_2$



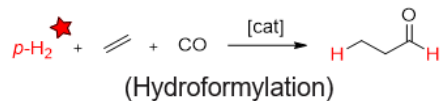
## d) Single proton exchange



## b) Pairwise replacement with $p\text{-H}_2$



## c) Single proton transfer



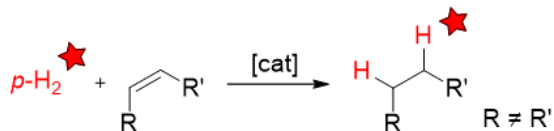
[a] C. R. Bowers and D. P. Weitekamp, *J. Am. Chem. Soc.* **1987**, 109, 5541-5542.

[b] A. Harthun, R. Giernoth, C. J. Elsevier and J. Bargon, *Chem. Commun.* **1996**, 2483.

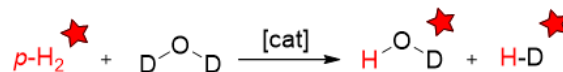
[c] A. B. Permin and R. Eisenberg, *J. Am. Chem. Soc.*, **2002**, 124, 12406-12407.

[d] S. Lehmkuhl, M. Emondts, L. Schubert, P. Spanning, J. Klankermayer, B. Blümich and P. P. M. Schleker, *ChemPhysChem* **2017**, 18, 2426-2429.

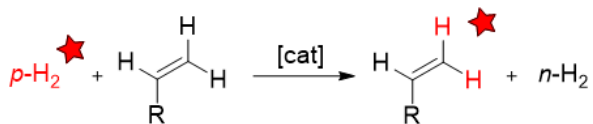
## a) Hydrogenation with $p\text{-H}_2$



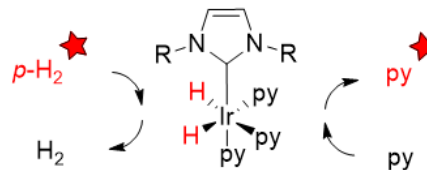
## d) Single proton exchange



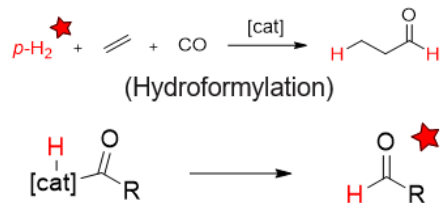
## b) Pairwise replacement with $p\text{-H}_2$



## e) SABRE: J-coupling transfer



## c) Single proton transfer



[a] C. R. Bowers and D. P. Weitekamp, *J. Am. Chem. Soc.* **1987**, 109, 5541-5542.

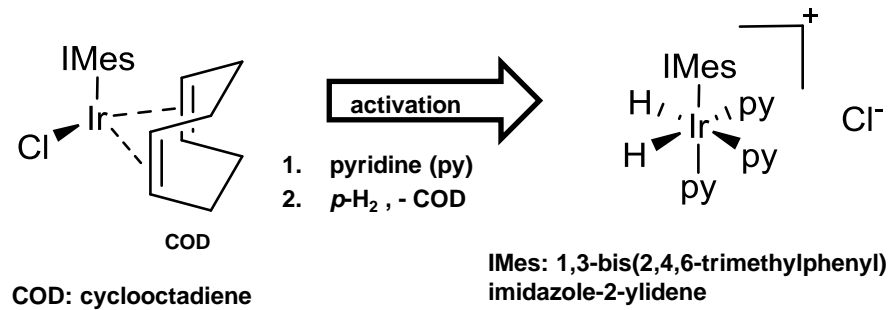
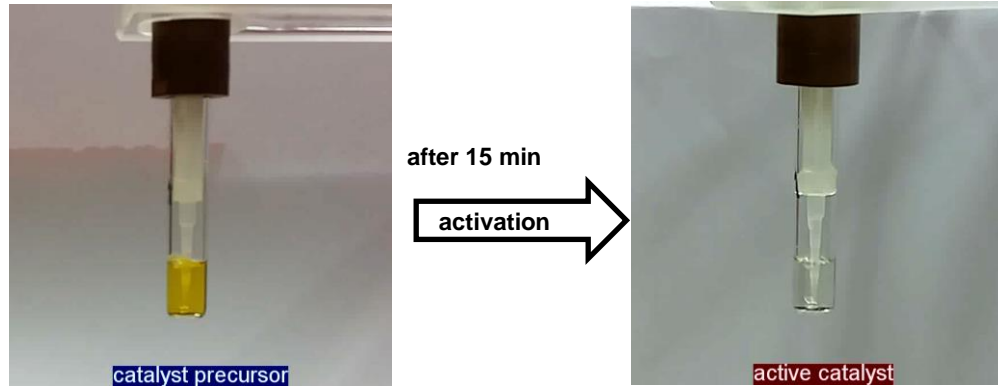
[b] A. Harthun, R. Giernoth, C. J. Elsevier and J. Bargon, *Chem. Commun.* **1996**, 2483.

[c] A. B. Permin and R. Eisenberg, *J. Am. Chem. Soc.*, **2002**, 124, 12406-12407.

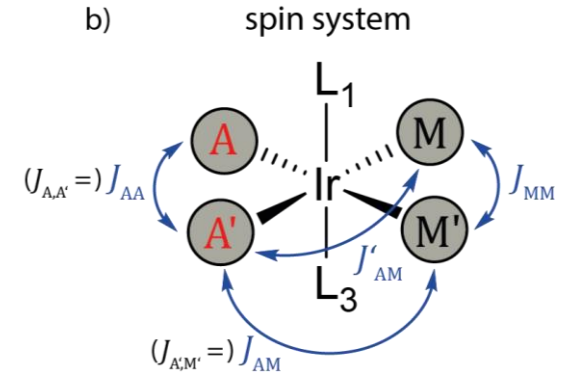
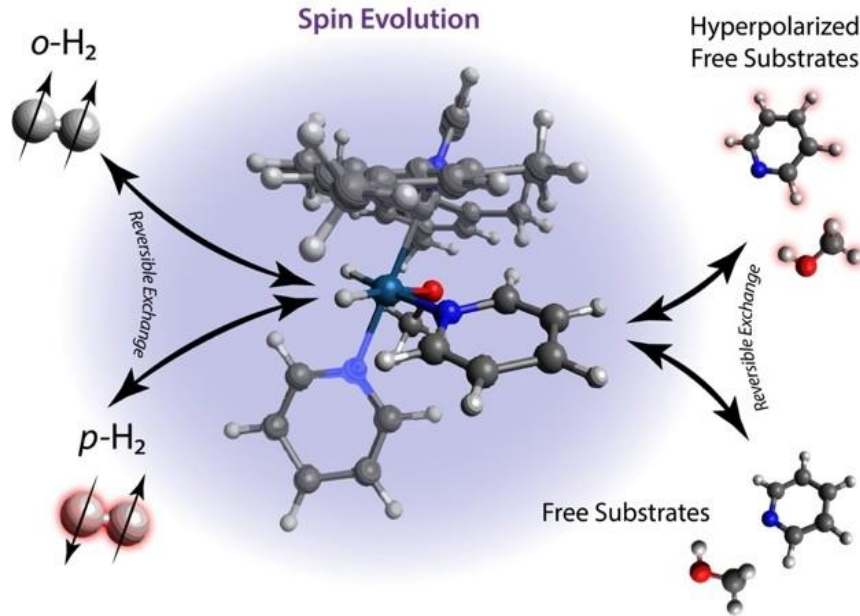
[d] S. Lehmkuhl, M. Emondts, L. Schubert, P. Spannring, J. Klankermayer, B. Blümich and P. P. M. Schleker, *ChemPhysChem* **2017**, 18, 2426-2429.

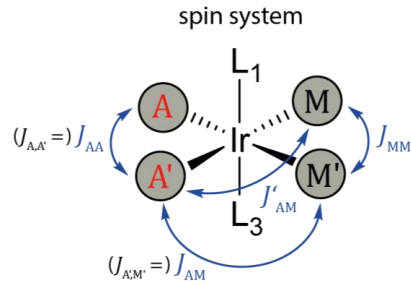
[e] R. Adams, J. Aguilar, K. Atkinson, M. Cowley, P. Elliott, S. Duckett, G. Green, I. Khazal, J. Lopez-Serrano and D. Williamson, *Science* **2009**, 323, 1708.



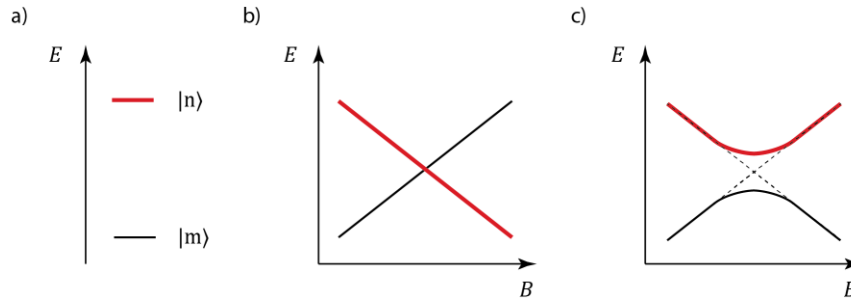


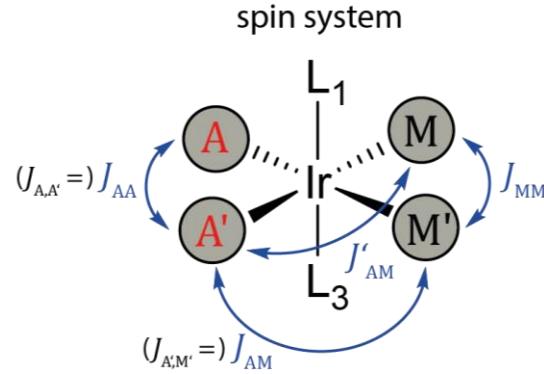
## Signal Amplification by Reversible Exchange





$$x = \frac{J}{\Delta\nu}$$



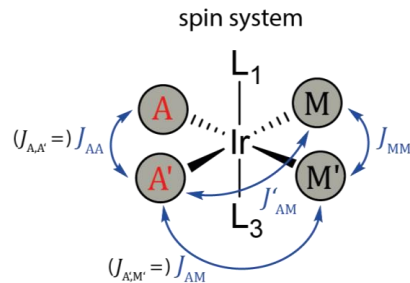


$$x = \frac{J}{\Delta\nu}$$

For pyridine SABRE:

~ 6 mT (60 G) for  $^1\text{H}$

~ 0.6  $\mu\text{T}$  for  $^{15}\text{N}$



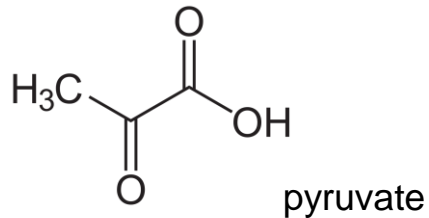
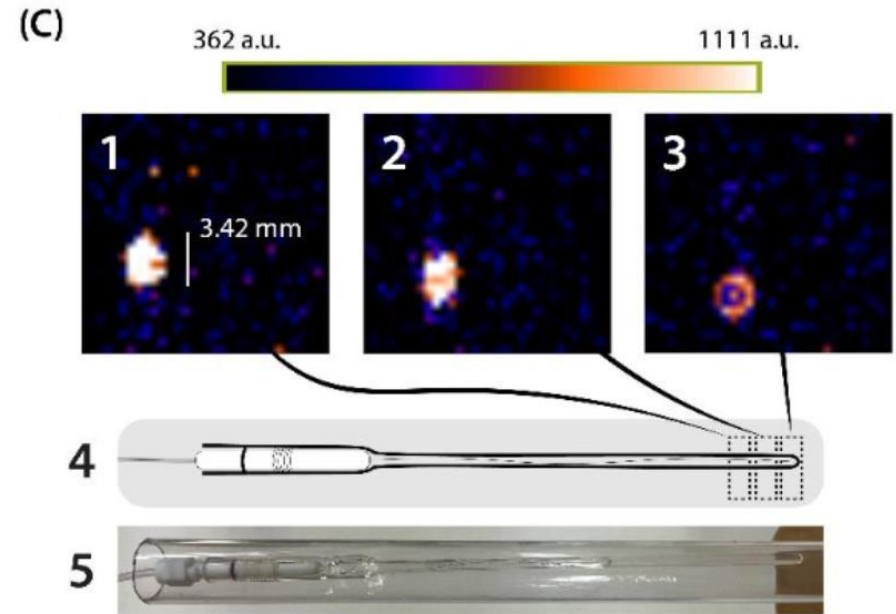
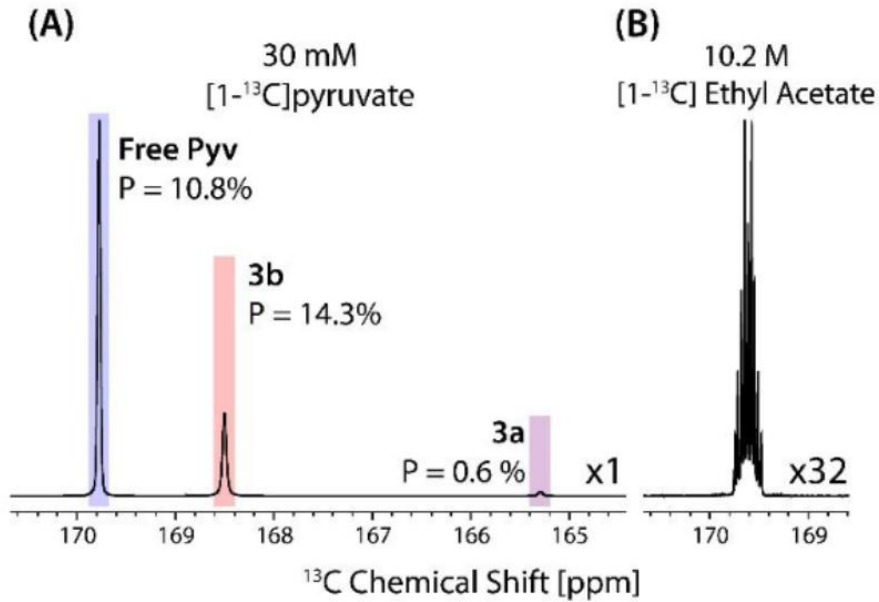
$$B_{\text{evo}} = \frac{\pm(J_{AA} - J_{MM})}{\gamma_A(1 - \delta_A) - \gamma_M(1 - \delta_M)}$$

$$B_{\text{evo}} = \frac{\pm(J_{AA} + J_{MM}) \mp \frac{1}{2}(J_{AM} + J'_{AM})}{\gamma_A(1 - \delta_A) - \gamma_M(1 - \delta_M)}$$

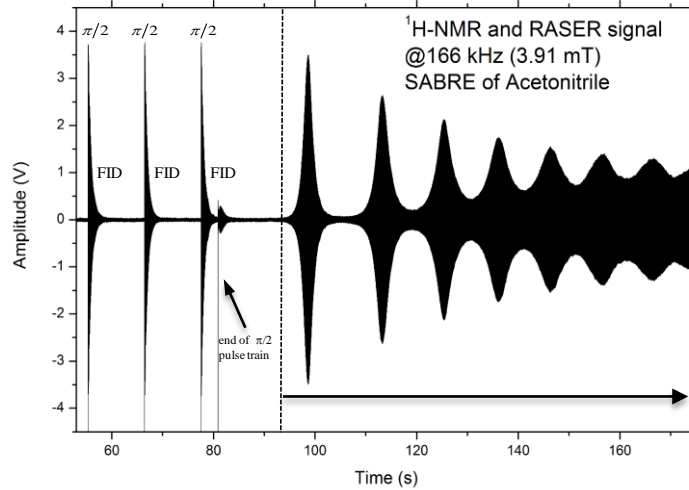


**SABRE-SHEATH**

(Shield Enables Alignment Transfer to Heteronuclei)



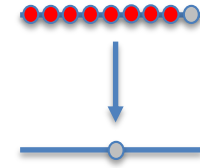
P. TomHon\*, M. Abdulmojeed, I. Adelabu, S. Nantogma, M. S. Hafez Kabir, S. Lehmkuhl, E. Y. Chekmenev, and T. Theis, Temperature Cycling Enables Efficient <sup>13</sup>C SABRESHEATH Hyperpolarization and Imaging of [1-<sup>13</sup>C]Pyruvate, *accepted JACS* 2021.



Unpublished data

**Need enough Polarization to surpass the RASER threshold**

Population inversion by hyperpolarization:



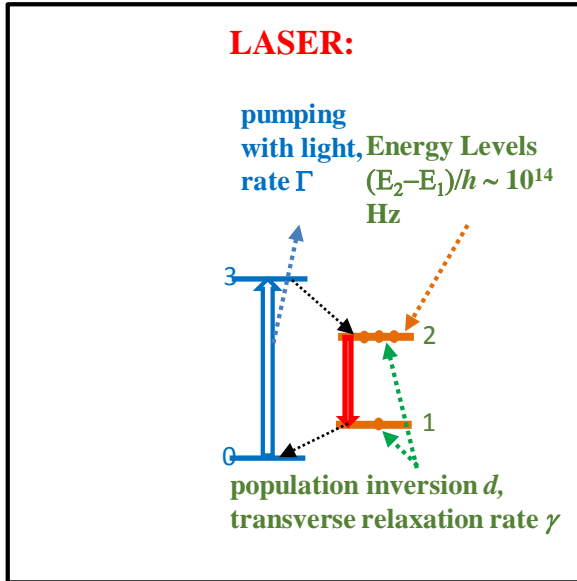
$$P_{threshold} = \frac{-4}{\mu_0 \eta \hbar \gamma^2 T_2^* Q n_s}$$

$\gamma$ : gyromagnetic ratio of the spin species

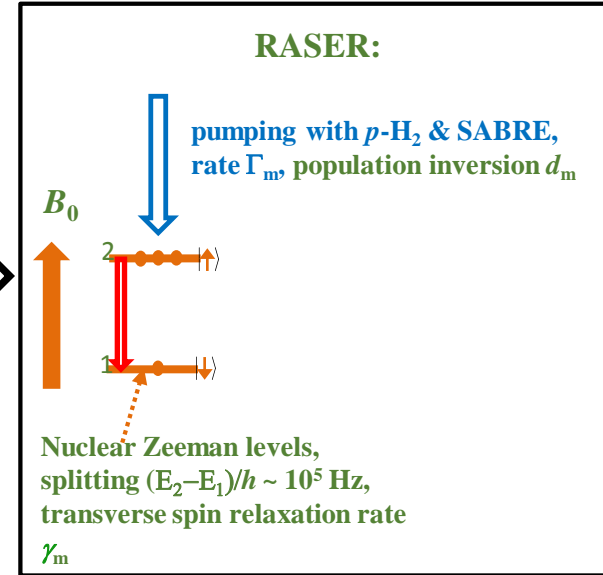
$T_2^*$ : effective transverse relaxation time

$Q$ : Quality factor of the resonant circuit

$n_s$ : Spin number density

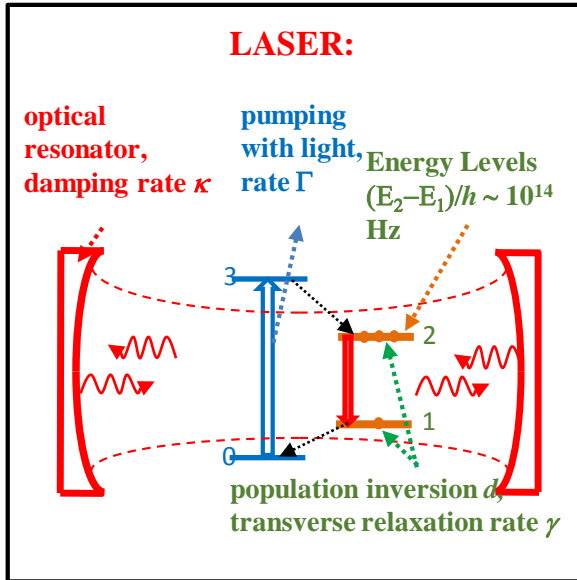


- Pumping with light or electrons with  $\Gamma$
- Atoms or molecules as LASER medium



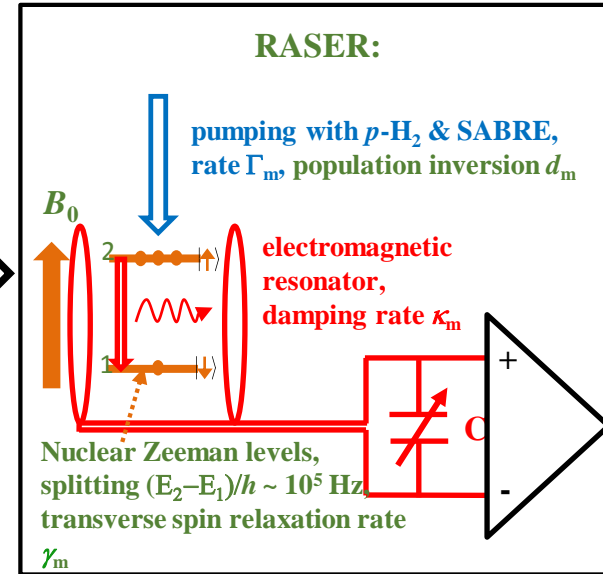
- Pumping with para-H<sub>2</sub> (spin zero) with  $\Gamma_m$
- Nuclear spins as RASER medium





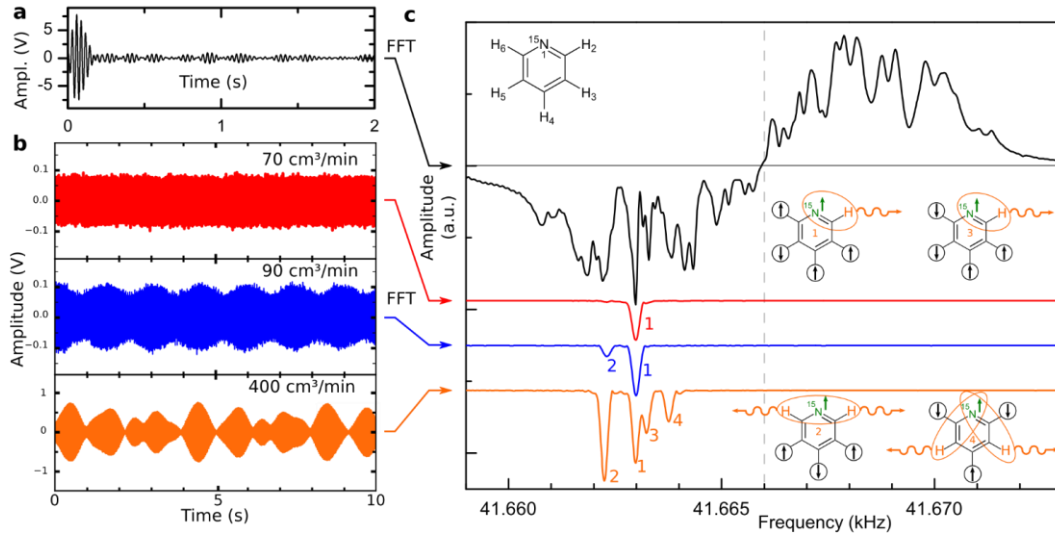
- Pumping with light or electrons with  $\Gamma$
- Atoms or molecules as LASER medium
- Optical resonator with damping rate  $\kappa$
- LASER emission starts at threshold  $d_0$

Frequency:  $\sim 10^{14}$  Hz,  $\Delta v_L/v_L \sim 10^{-14}$



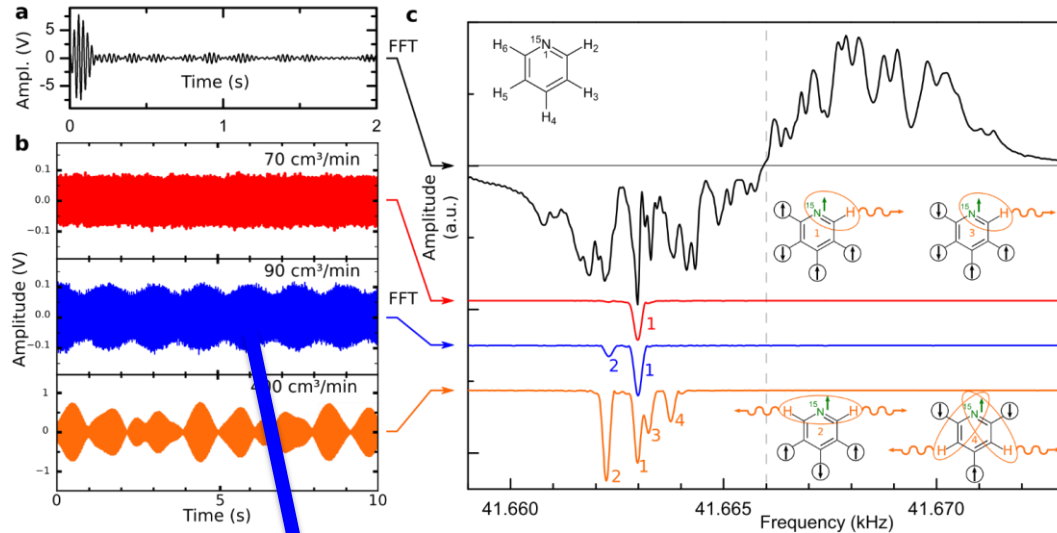
- Pumping with para-H<sub>2</sub> (spin zero) with  $\Gamma_m$
- Nuclear spins as RASER medium
- LC-resonator with damping rate  $\kappa_m$
- RASER emission starts at threshold  $d_{0,m}$

Frequency:  $\sim 10^5$  Hz,  $\Delta v/v < 10^{-10}$

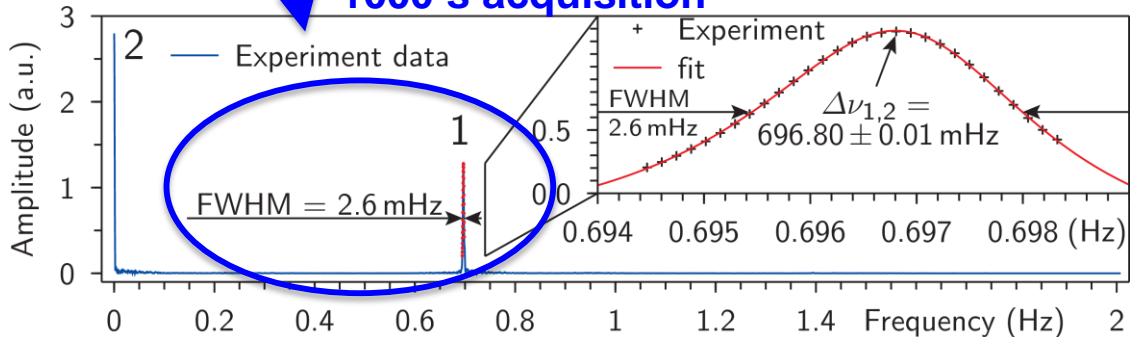


**MULTIMODE RASER  
OF MOLECULAR  
STATES!**

*M. Sűfke, S. Lehmkuhl,  
A. Liebisch, B. Blűmich &  
S Appelt: Nature Physics  
13, 568–572 (2017)*



1000 s acquisition



**MULTIMODE RASER  
OF MOLECULAR  
STATES!**

**UNPRECEDENTED  
RESOLUTION**

*M. Sűfke, S. Lehmkuhl,  
A. Liebisch, B. Blűmich &  
S Appelt: Nature Physics  
13, 568–572 (2017)*

## How to increase an NMR signal:

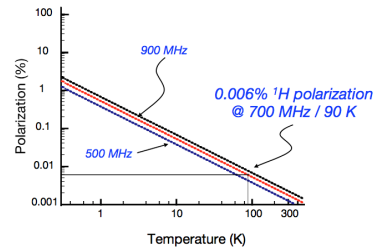
Spin density

Concentration  
Abundance

$$M = \frac{\hbar}{2} N \gamma P$$

Polarization

Temperature  
Magnetic field



Hyperpolarization

DNP  
SEOP  
PHIP

