

# Searching for **New Physics** at the intensity frontier

(at low energies with lots of particles)  
Part II

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Heidelberg Graduate Days  
Heidelberg, April 2018



The weak mixing angle

(also: Weinberg-angle)



# The weak mixing angle

- One of the fundamental parameters of the standard model
- Electroweak symmetry breaking creates photon and  $Z^0$
- Angle shows up both in masses and couplings (charges)

$$\begin{pmatrix} \gamma \\ Z^0 \end{pmatrix} = \begin{pmatrix} \cos \theta_W & \sin \theta_W \\ -\sin \theta_W & \cos \theta_W \end{pmatrix} \begin{pmatrix} B^0 \\ W^0 \end{pmatrix}$$

$$\cos \theta_W = \frac{m_W}{m_Z}$$

$$\sin^2 \theta_W = \frac{g'^2}{g^2 + g'^2}$$

# Which weak mixing angle?

- The last slide is true at tree level
- But there are quantum corrections...

Two options:

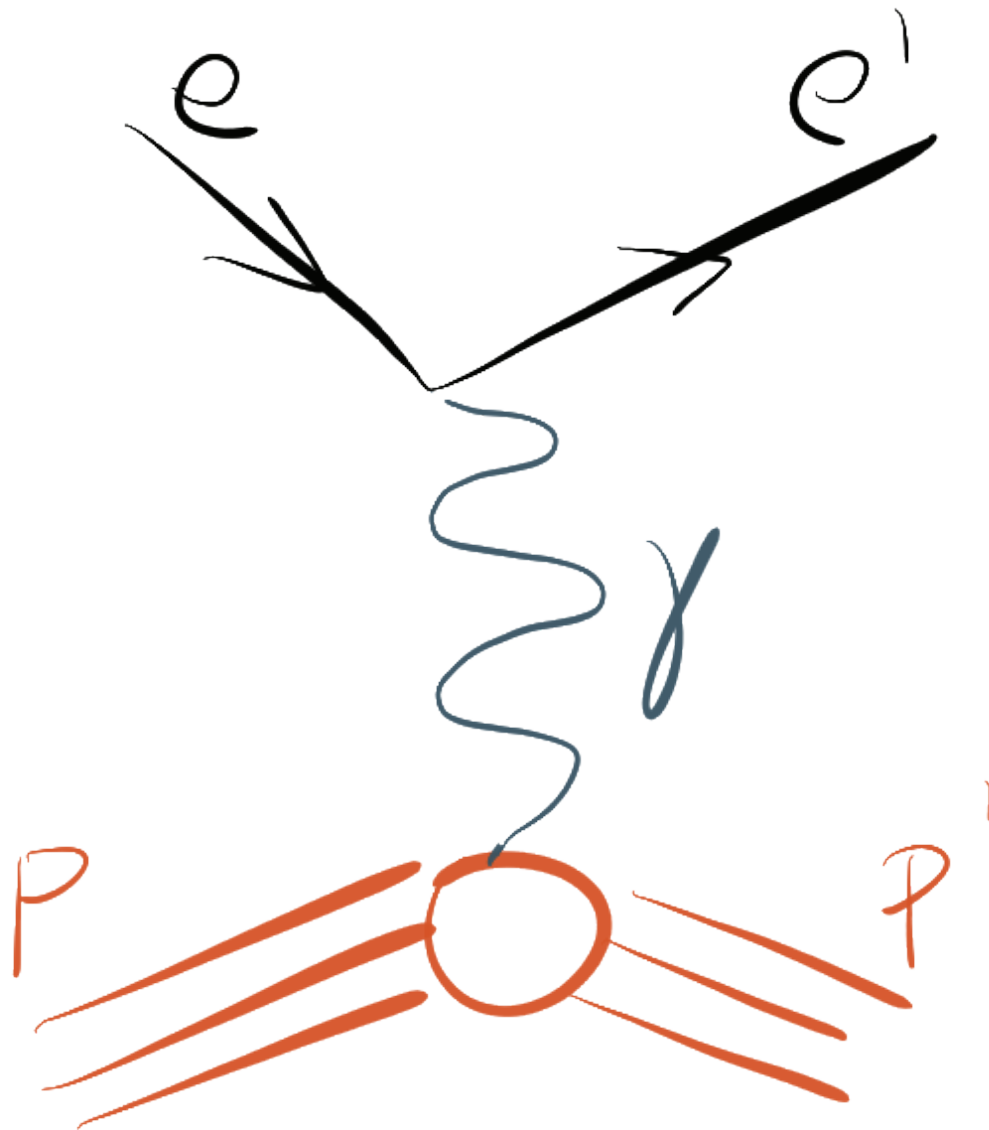
- Use the masses for the definition:  
(at all orders of perturbation theory)  
"On-shell scheme"
- Or use the couplings:  
(which change with energy, and so does  
the angle)  
" $\overline{\text{MS}}$ -scheme"
- Use second option from here on

$$\cos \theta_W = \frac{m_W}{m_Z}$$

$$\sin^2 \theta_W = \frac{g'^2}{g^2 + g'^2}$$

$$\sin^2 \theta_W(q^2)$$

# Weak mixing angle and charges



Proton electric charge

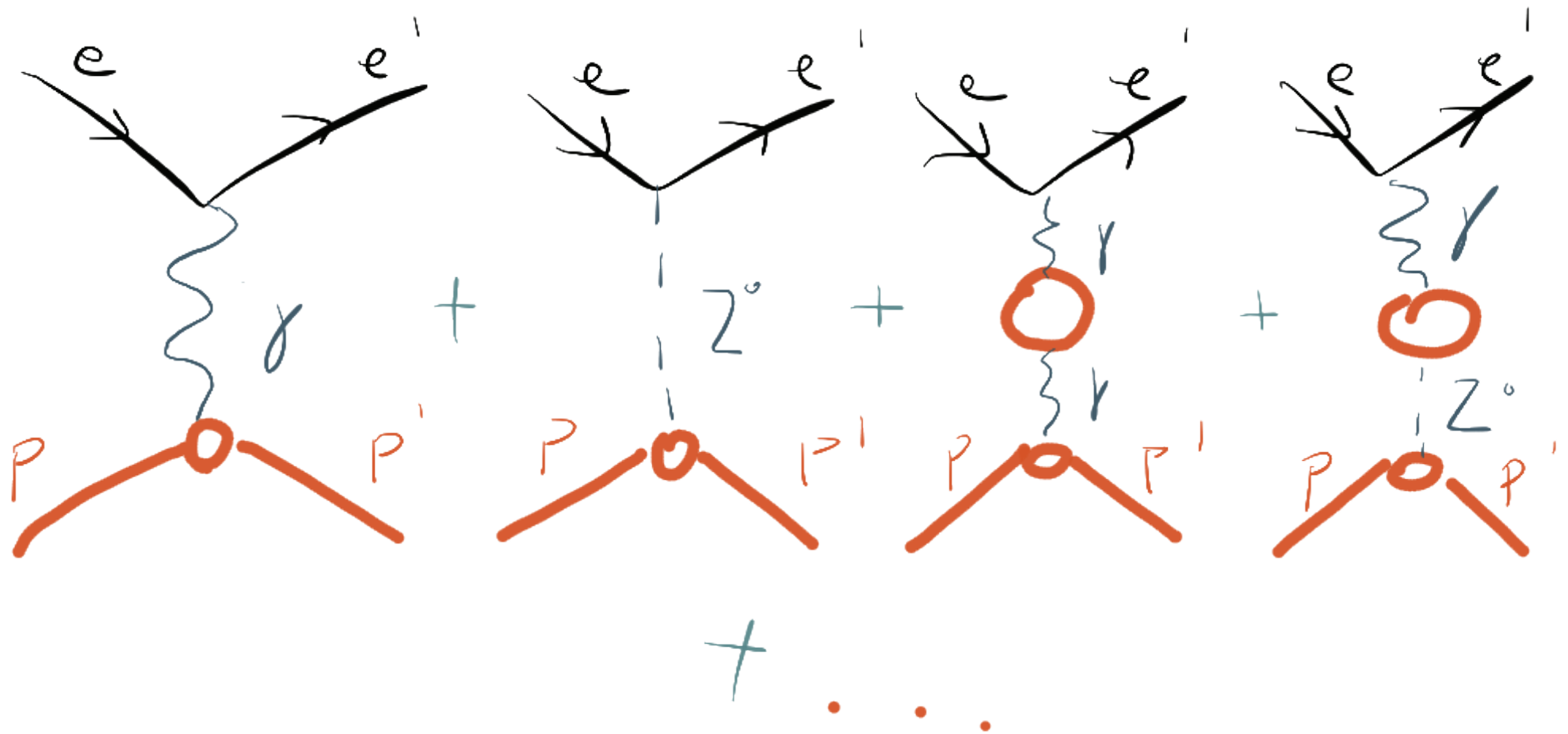
$+e$



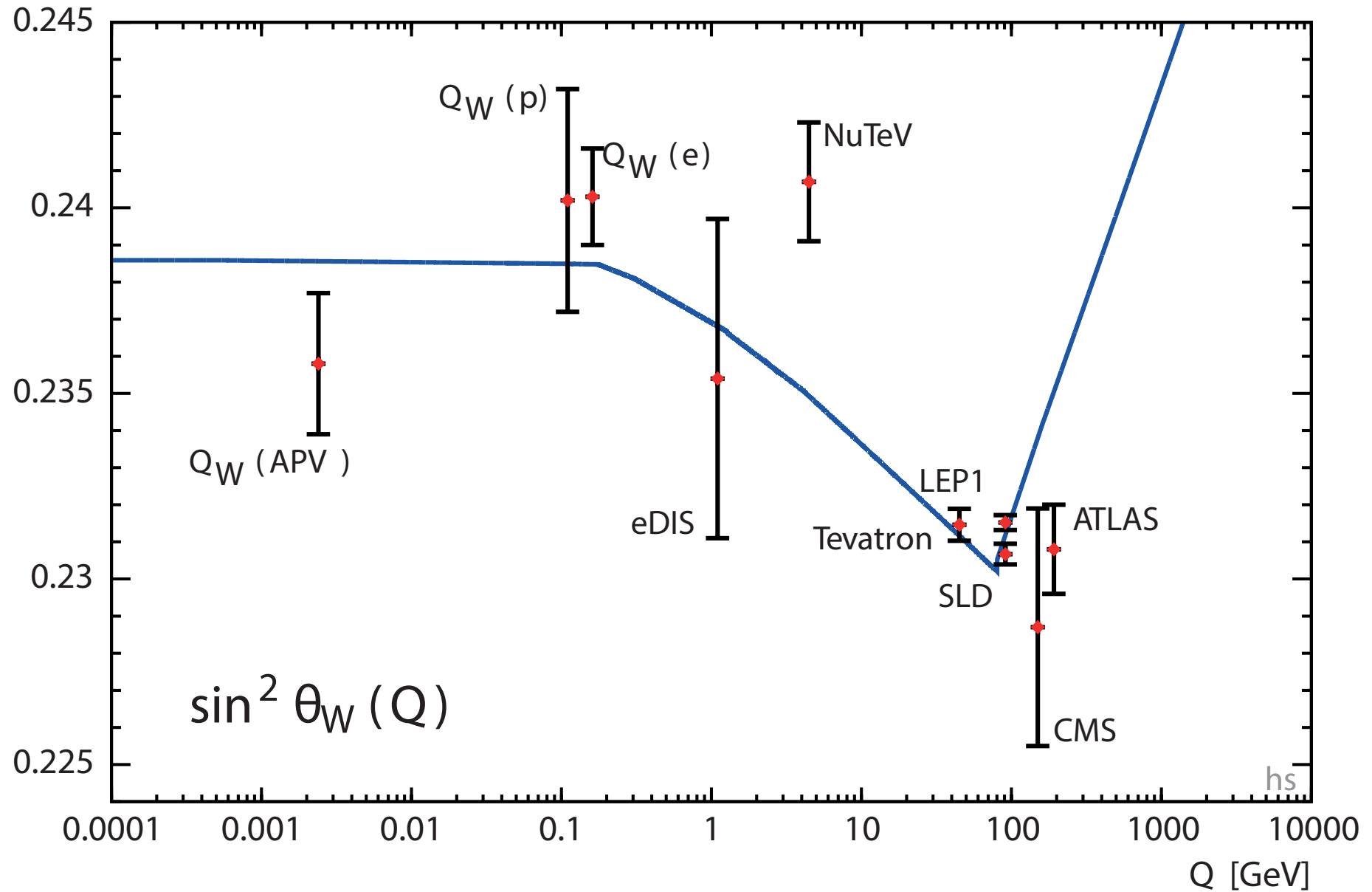
Proton weak charge

$1 - 4 \sin^2 \theta_w$

# Scale dependence (running) of $\sin^2\theta_w$

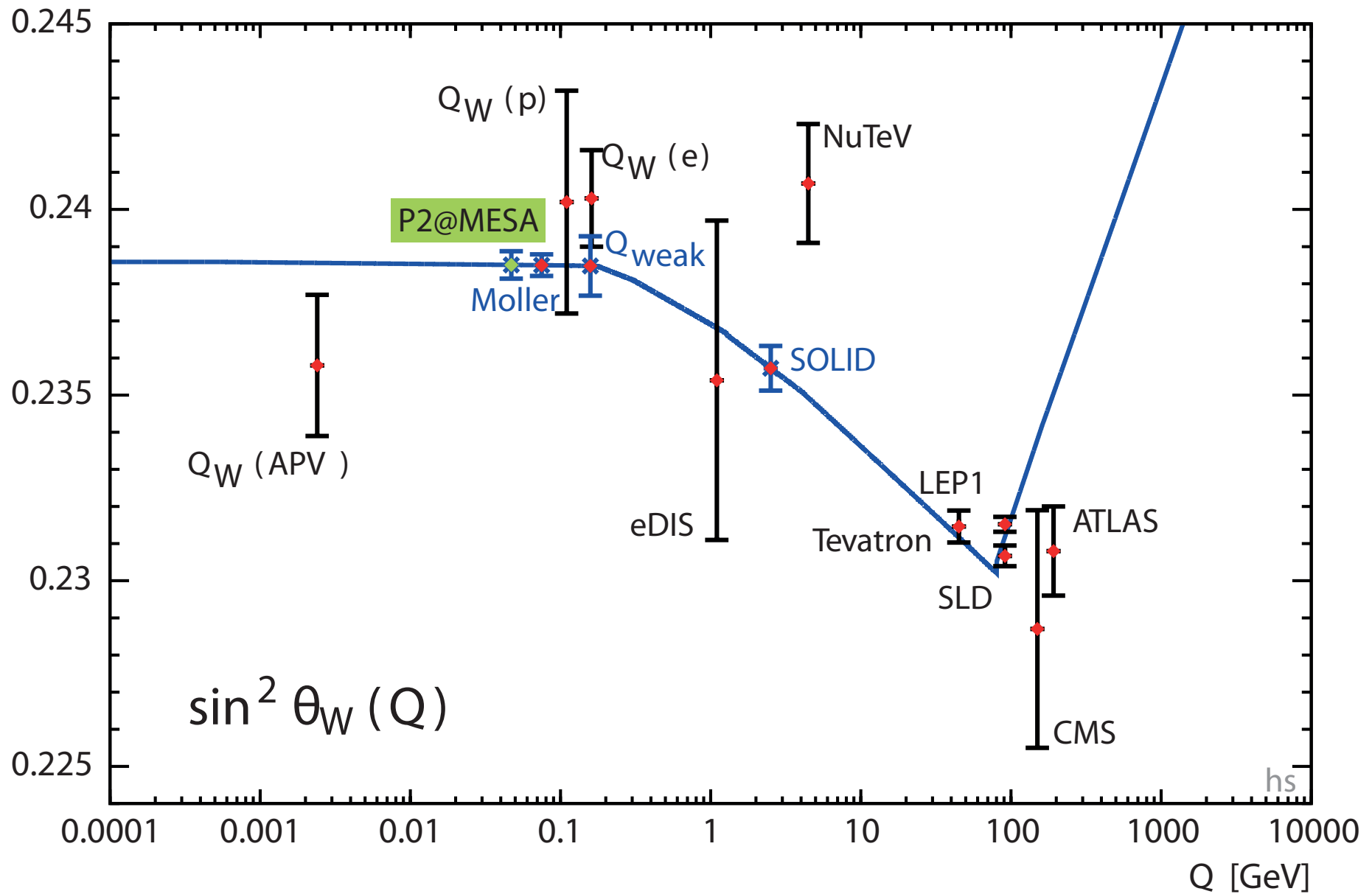


# Scale dependence (running) of $\sin^2\theta_W$

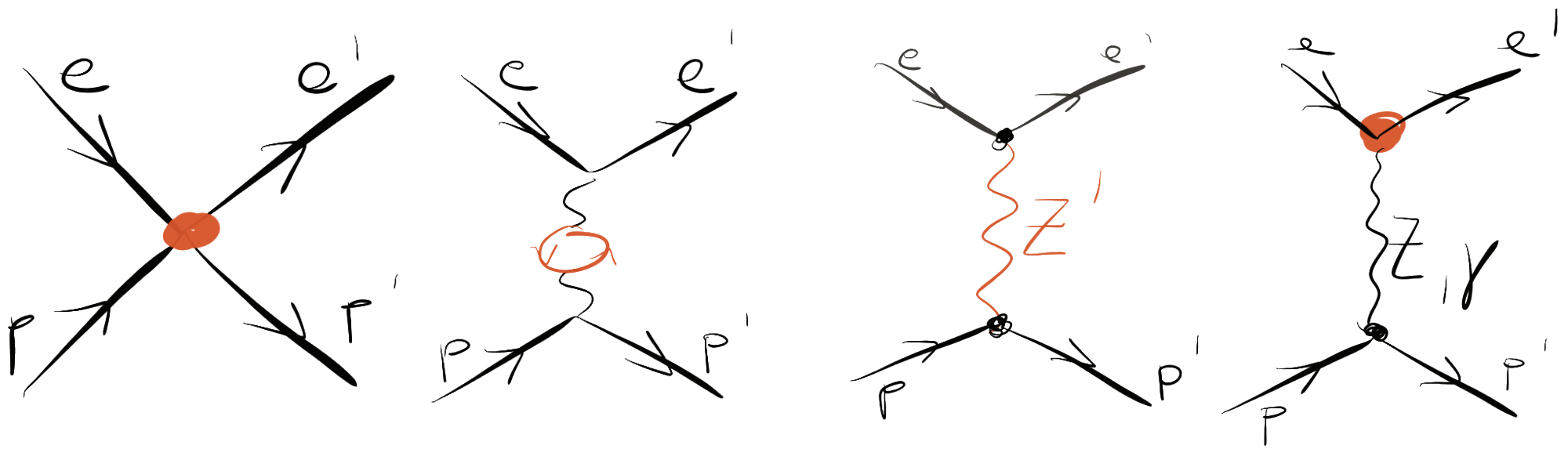
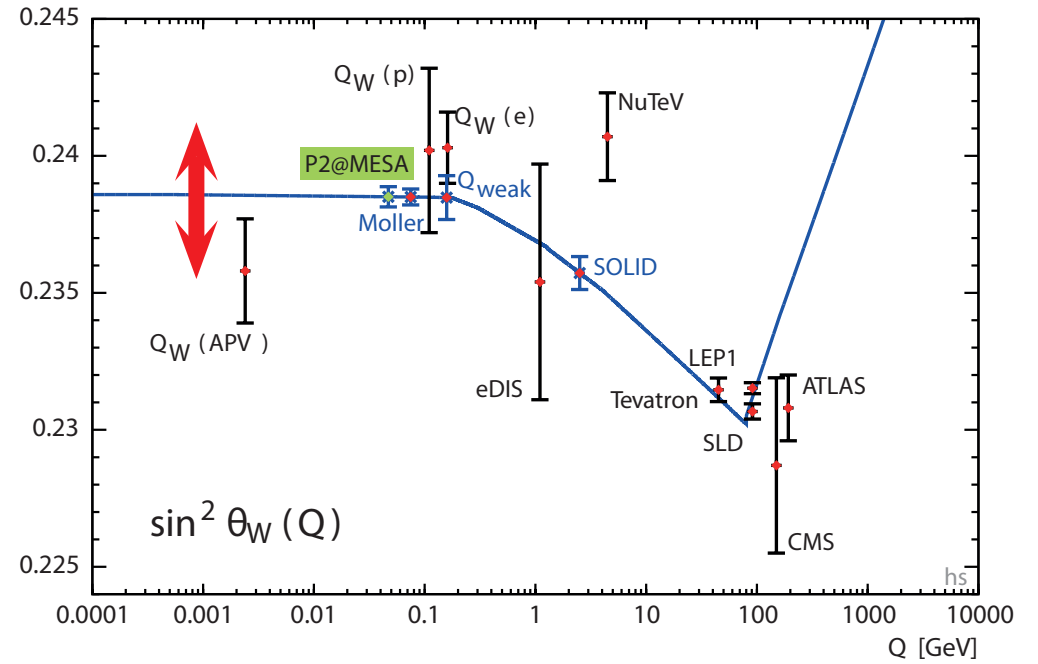
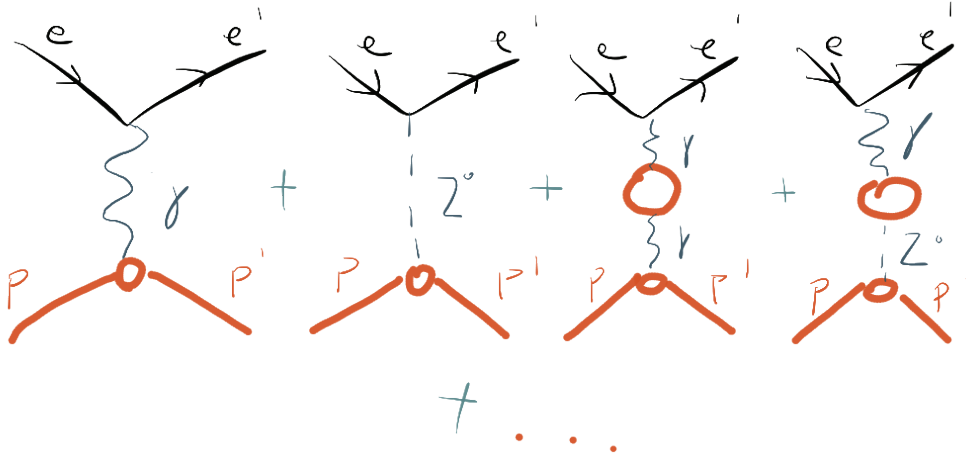




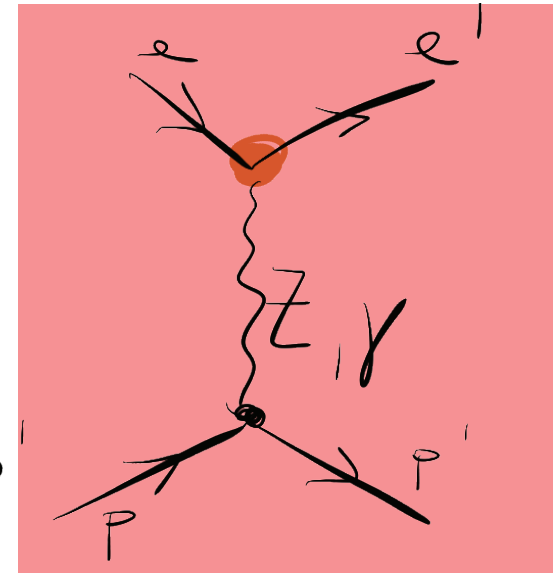
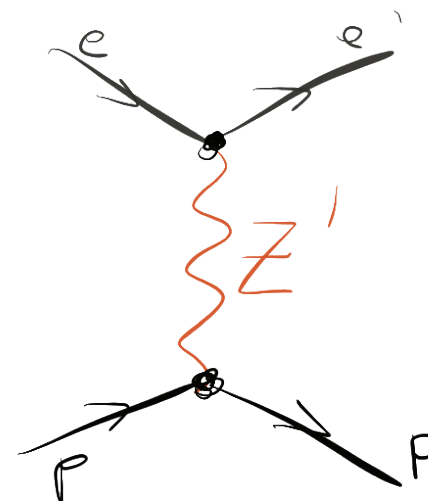
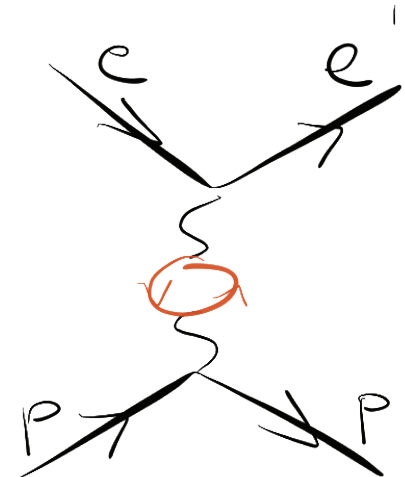
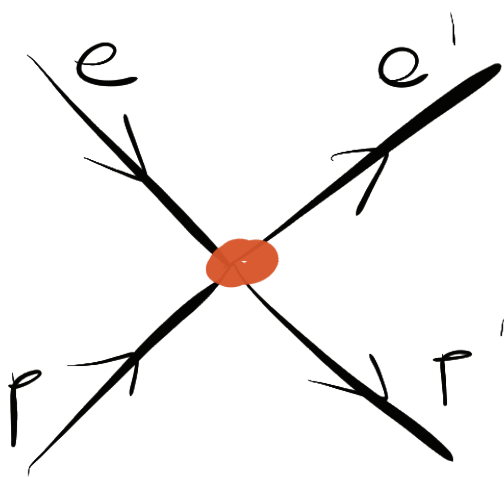
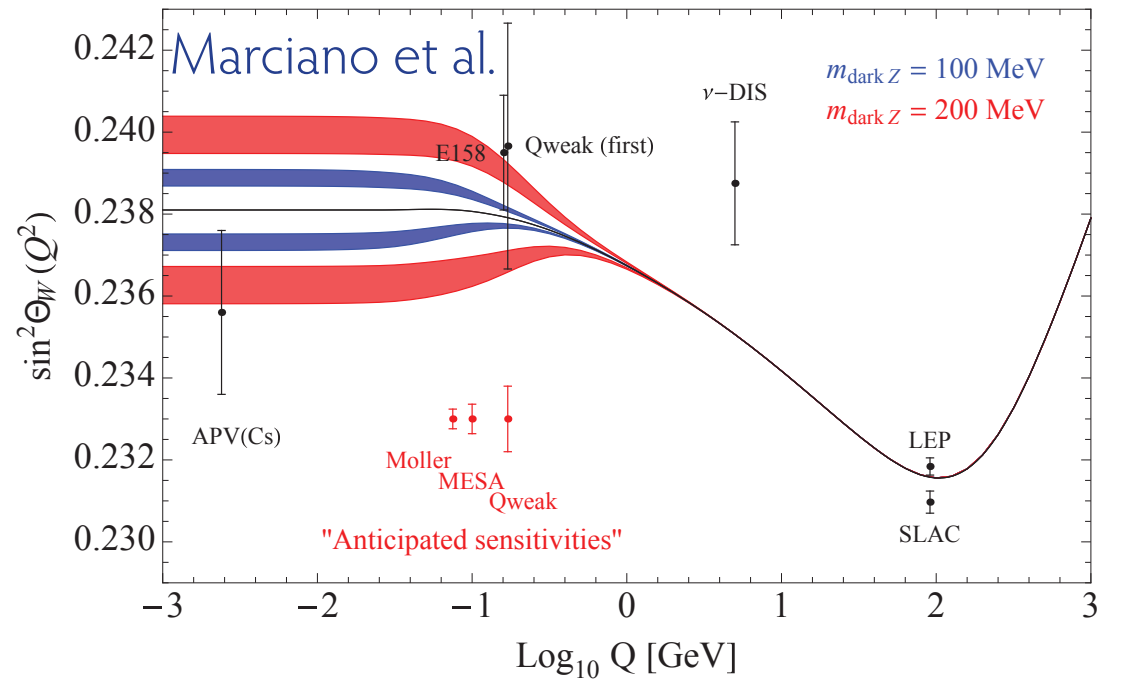
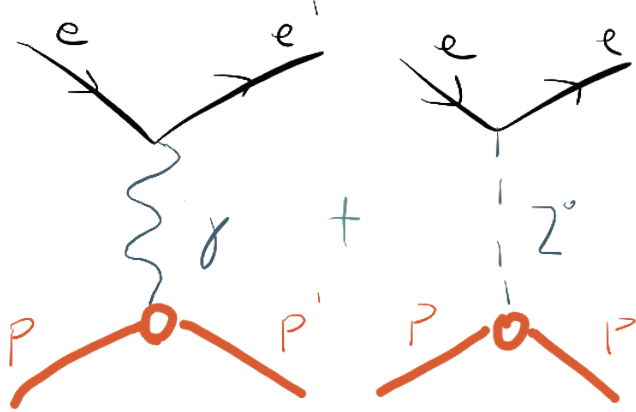
# Scale dependence (running) of $\sin^2\theta_W$



# New Physics in the running



# Dark Z in mixing

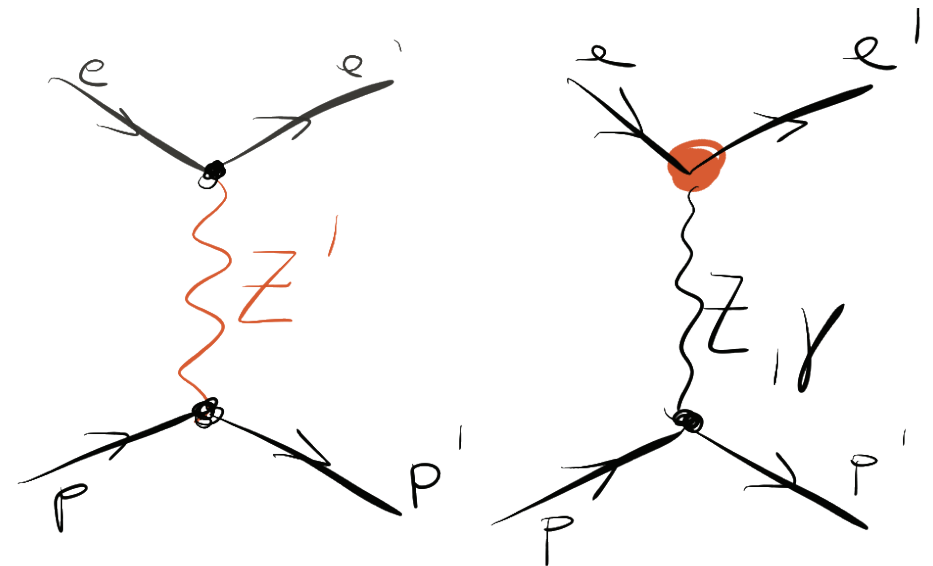
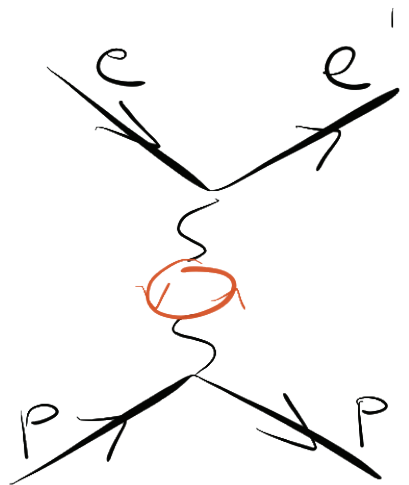
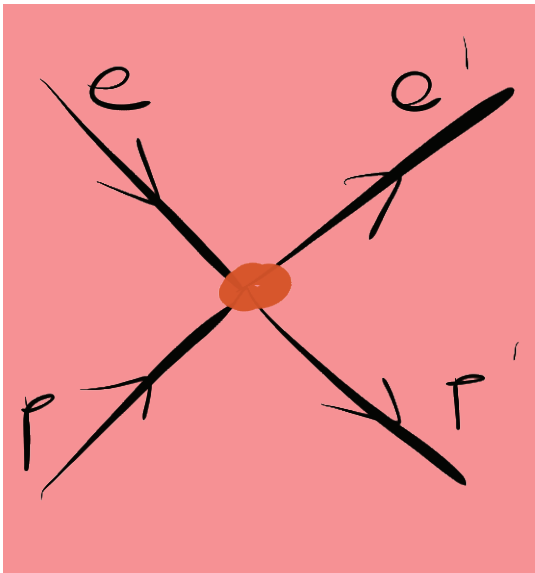
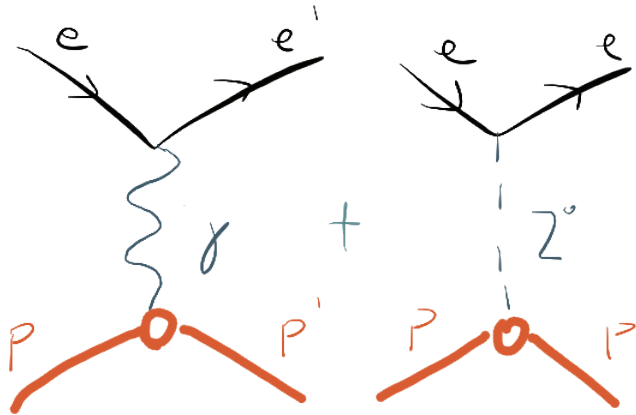


# Contact Interactions

Contact interactions up to

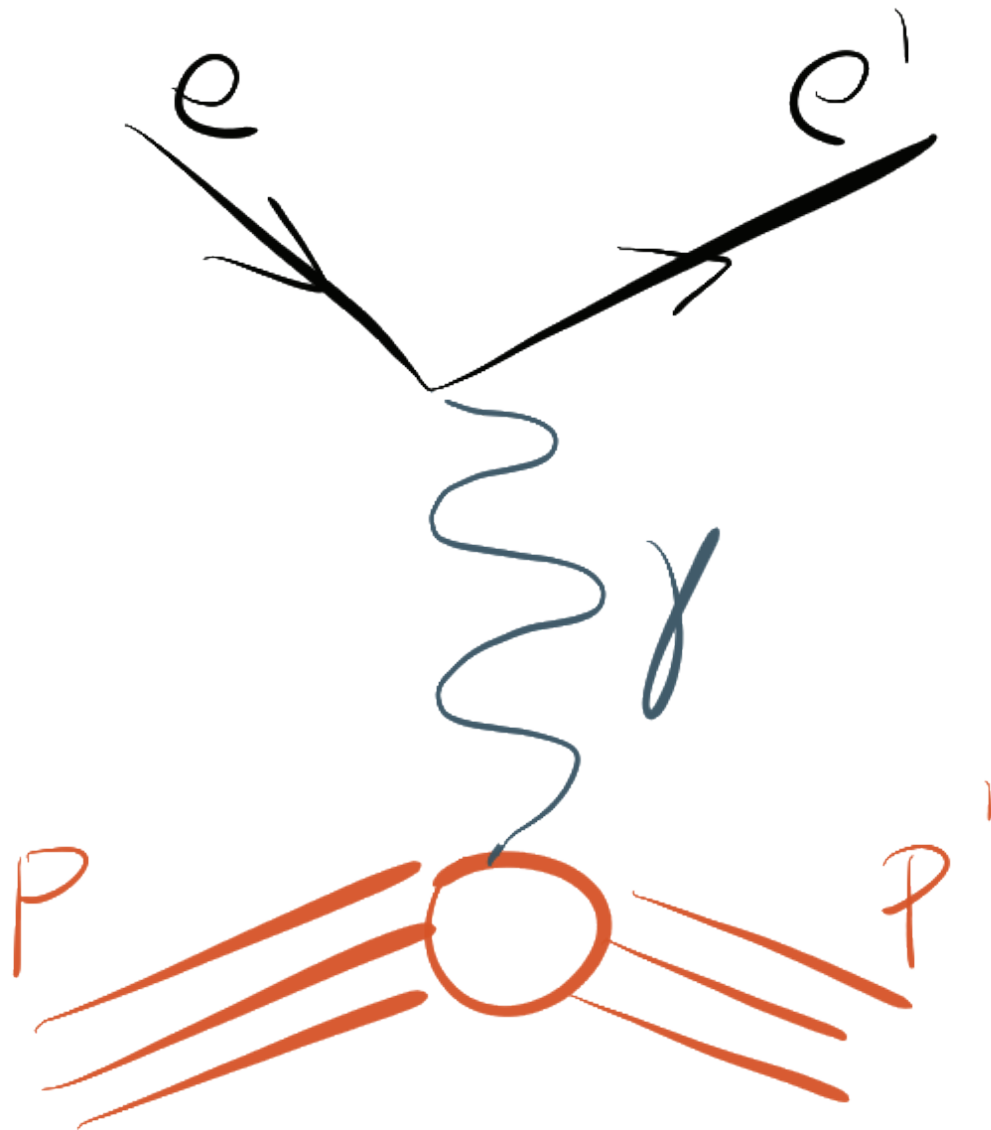
**49 TeV**

(comparable to LHC at  $300 \text{ fb}^{-1}$ )



How to measure the weak charge?

# Weak mixing angle and charges



Proton electric charge

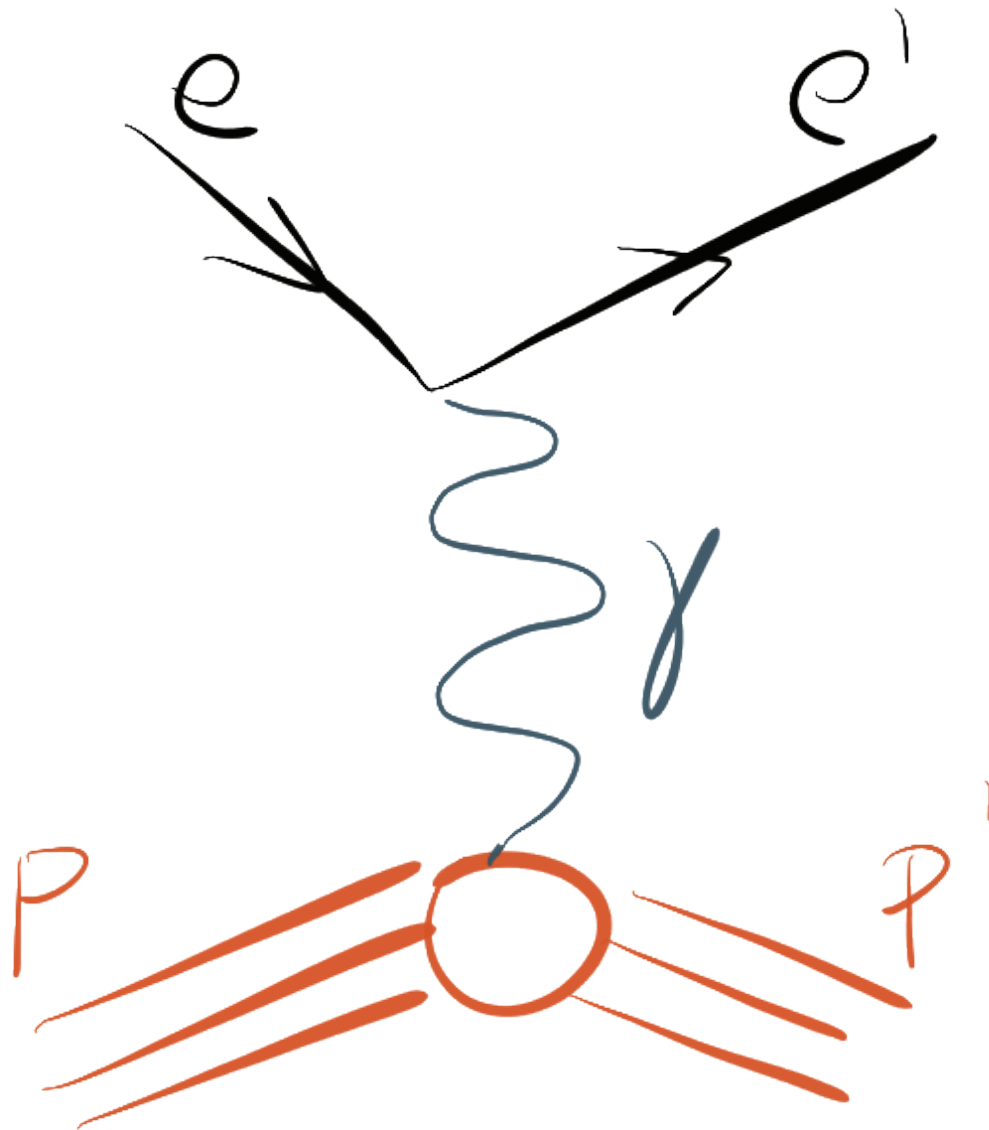
$+e$



Proton weak charge

$1 - 4 \sin^2 \theta_w$

# Weak mixing angle and charges



Proton electric charge

$+e$

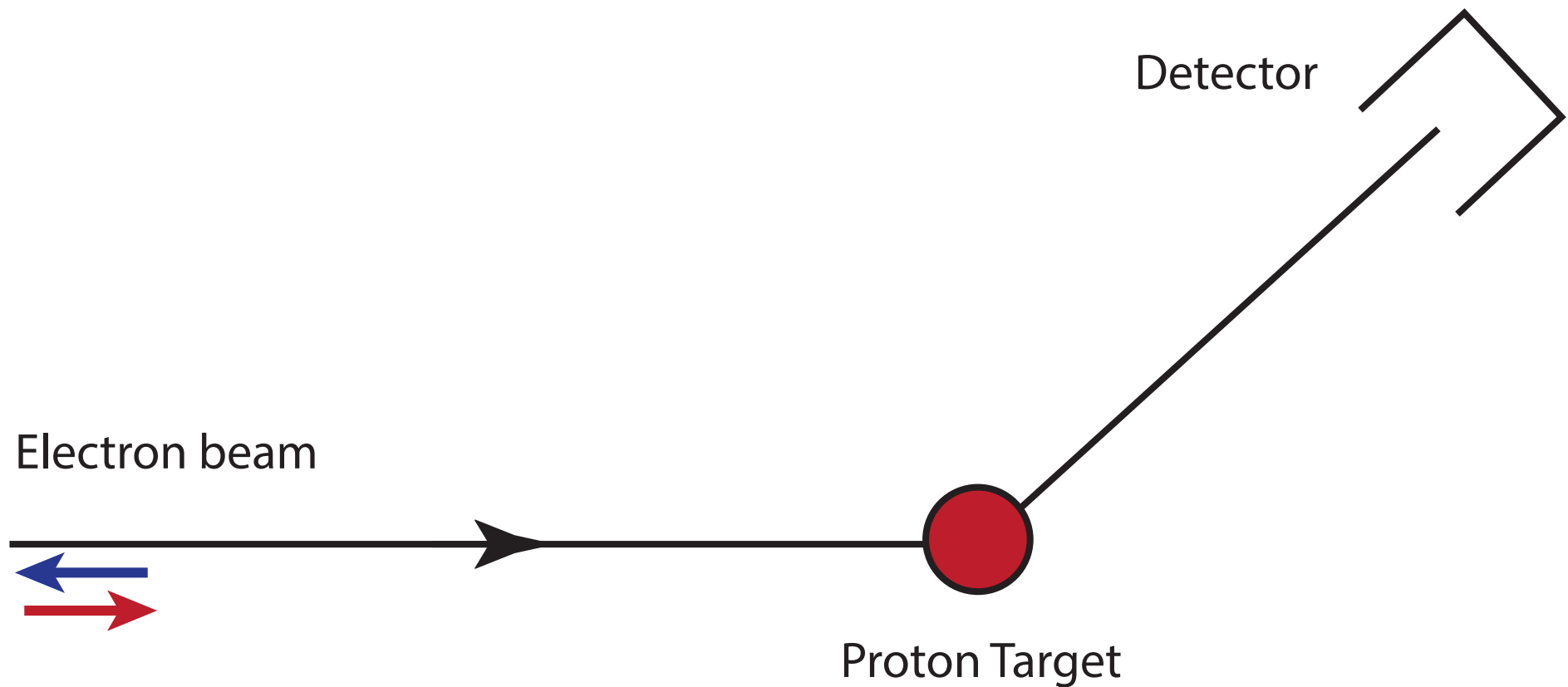


Violates parity!

Proton weak charge

$1 - 4 \sin^2 \theta_w$

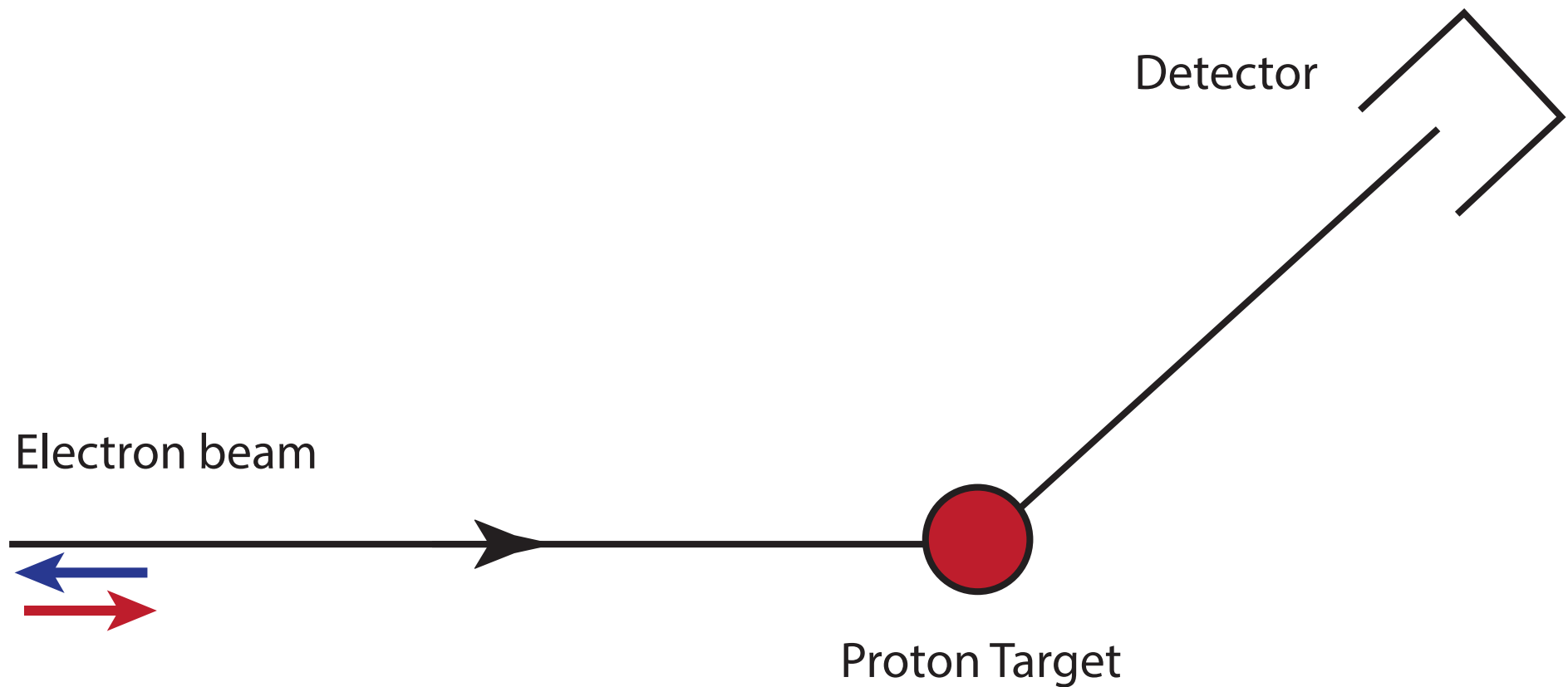
# Parity violating electron scattering





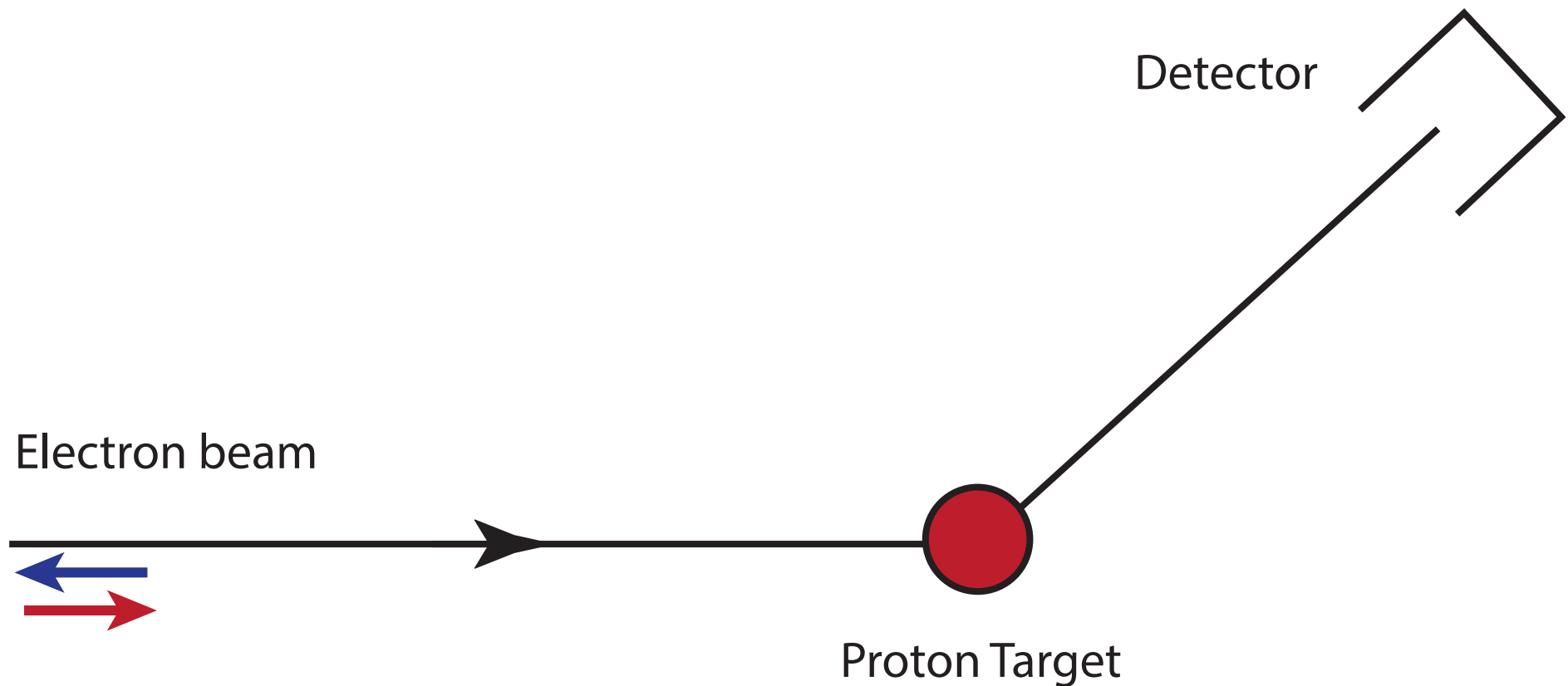
# Parity violating electron scattering

$$A_{PV} = \frac{N_R - N_L}{N_R + N_L}$$



# Parity violating electron scattering

$$A_{PV} = \frac{N_R - N_L}{N_R + N_L} = \frac{G_F Q^2}{4\sqrt{2}\pi\alpha} (Q_W - F(Q^2))$$



# Parity violating electron scattering

Momentum transfer  
sets scale

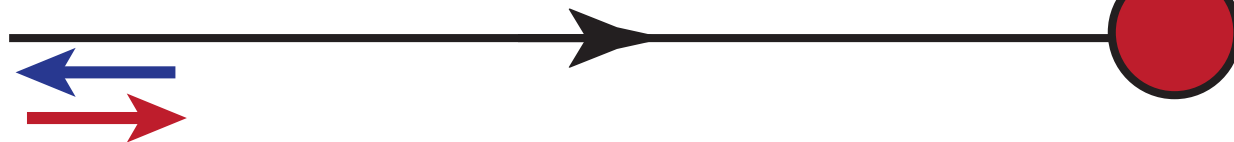
Proton structure -  
small nuisance if  $Q^2$  small

$$A_{PV} = \frac{N_R - N_L}{N_R + N_L} = \frac{G_F Q^2}{4\sqrt{2}\pi\alpha} (Q_W - F(Q^2))$$

Weak charge -  
what we want

Detector

Electron beam



Proton Target

# Parity violating electron scattering

Momentum transfer  
sets scale

Proton structure -  
small nuisance if  $Q^2$  small

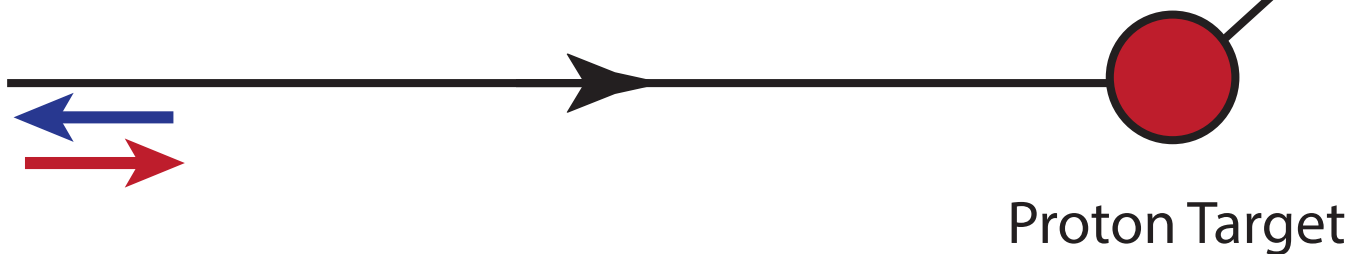
$$A_{PV} = \frac{N_R - N_L}{N_R + N_L} = \frac{G_F Q^2}{4\sqrt{2}\pi\alpha} (Q_W - F(Q^2))$$

Weak charge -  
what we want

Detector

$$\sin^2 \theta_W = \frac{1 - Q_W}{4}$$

Electron beam



# Why is this difficult?

- $\sin^2\theta_W \approx 0.25$ : Weak charge is tiny

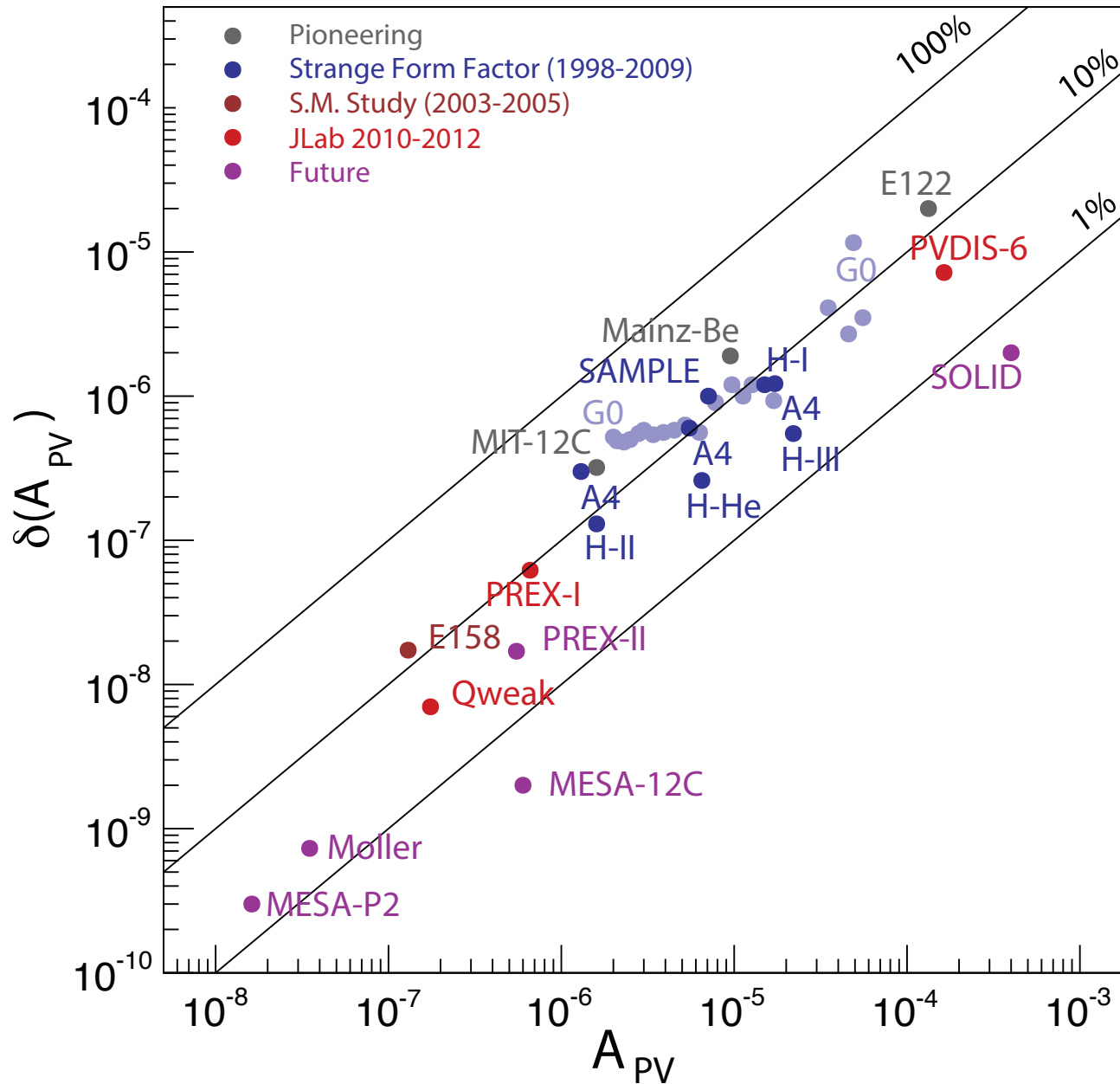
$$Q_W = 1 - 4 \sin^2 \theta_W$$

- At low  $Q^2$ : Asymmetry is tiny (40 parts per billion):  
need very large statistics

$$A_{PV} = \frac{N_R - N_L}{N_R + N_L} = \frac{G_F Q^2}{4\sqrt{2}\pi\alpha} (Q_W - F(Q^2))$$

- We are subtracting two huge numbers from each other  
(not really - switching helicity with a few KHz)

# PVeS Experiment Summary



# How much statistics do we need?

- Want to measure  $\sin^2\theta_W$  to 0.13%

- Need  $Q_W$  at 1.5%

$$\frac{\Delta \sin^2 \theta_W}{\sin^2 \theta_W} = \frac{1 - 4 \sin^2 \theta_W}{4 \sin^2 \theta_W} \frac{\Delta Q_W}{Q_W}$$

- Essentially means 1.5% on  $A_{PV}$

- $A_{PV}$  is 40 parts per billion

- $\delta(A_{PV})$  is 0.6 parts per billion

$$\delta(A_{PV}) \propto \frac{1}{\sqrt{N}}$$

- $N$  a few  $10^{18}$

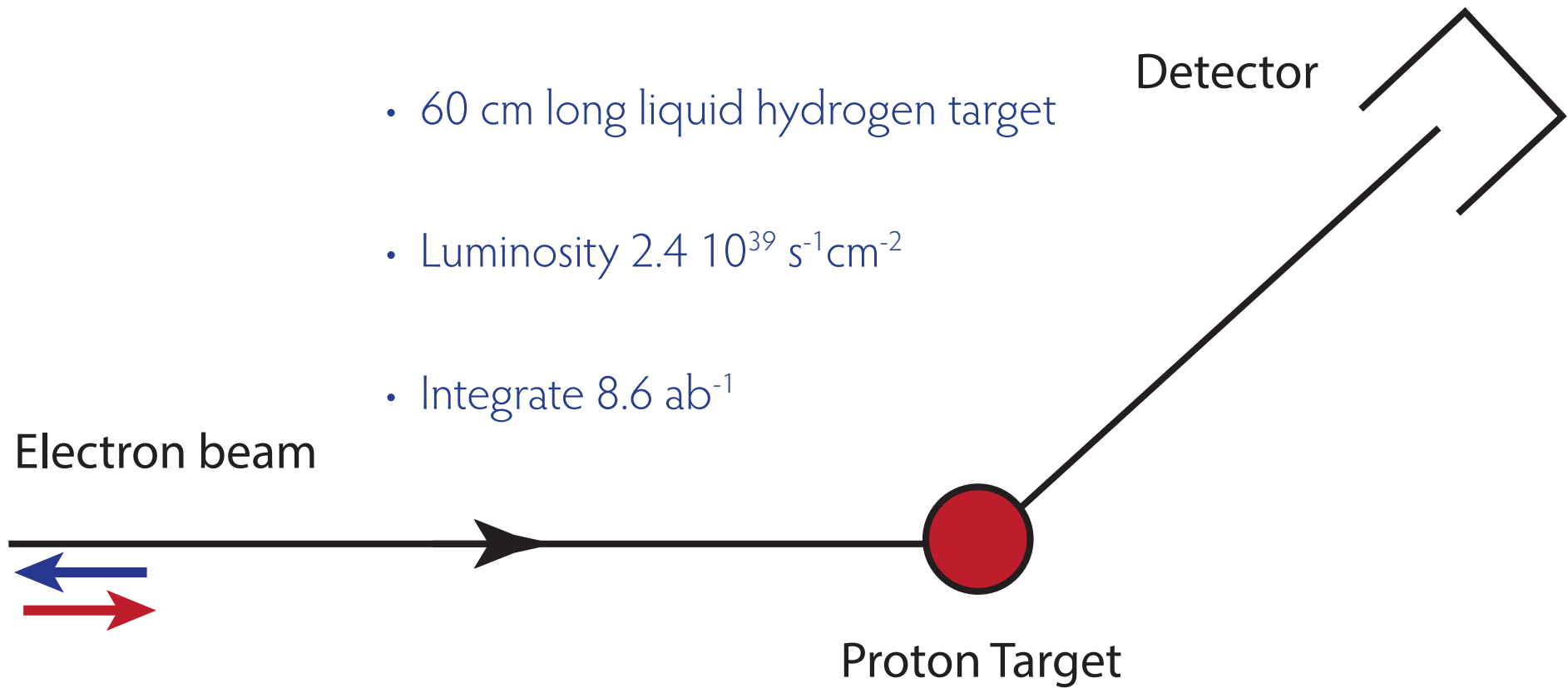
- Measure 10'000 hours (absolute maximum anyone thinks shifts are organisable)

- Need close to  $10^{11}$  electrons/s - 100 GHz

# Can we get that rate?

Yes!

- 150  $\mu\text{A}$  of electron beam current
- 60 cm long liquid hydrogen target
- Luminosity  $2.4 \cdot 10^{39} \text{ s}^{-1} \text{ cm}^{-2}$
- Integrate  $8.6 \text{ ab}^{-1}$





10'000 hours is 417 days 24/7 of measurements

Hard to get that amount of time at a shared  
accelerator facility...

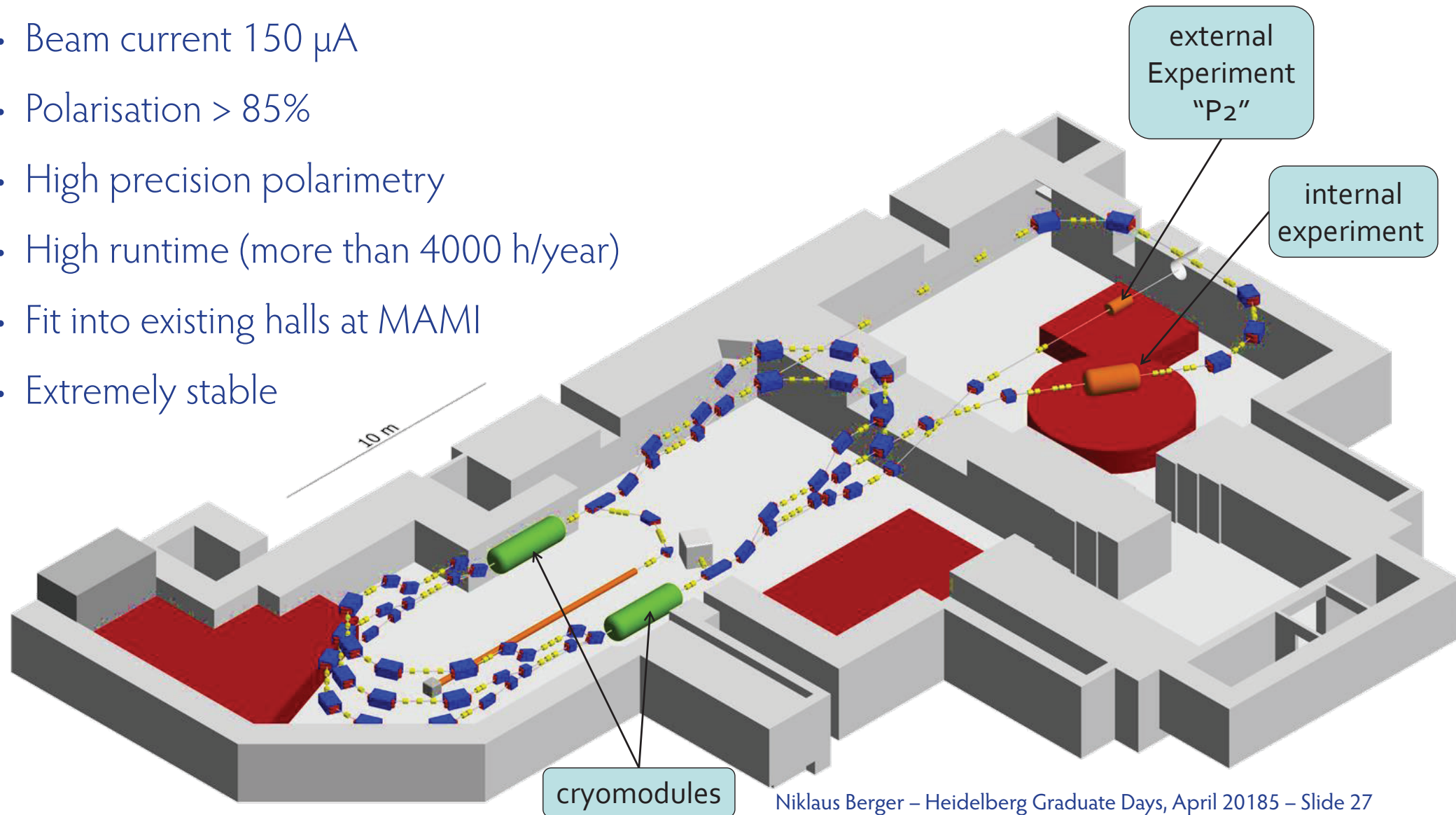
If you cannot rent it, build it:

# The MESA accelerator

Mainz Energy-recovery Superconducting Accelerator

# Requirements

- Beam current  $150 \mu\text{A}$
- Polarisation  $> 85\%$
- High precision polarimetry
- High runtime (more than 4000 h/year)
- Fit into existing halls at MAMI
- Extremely stable

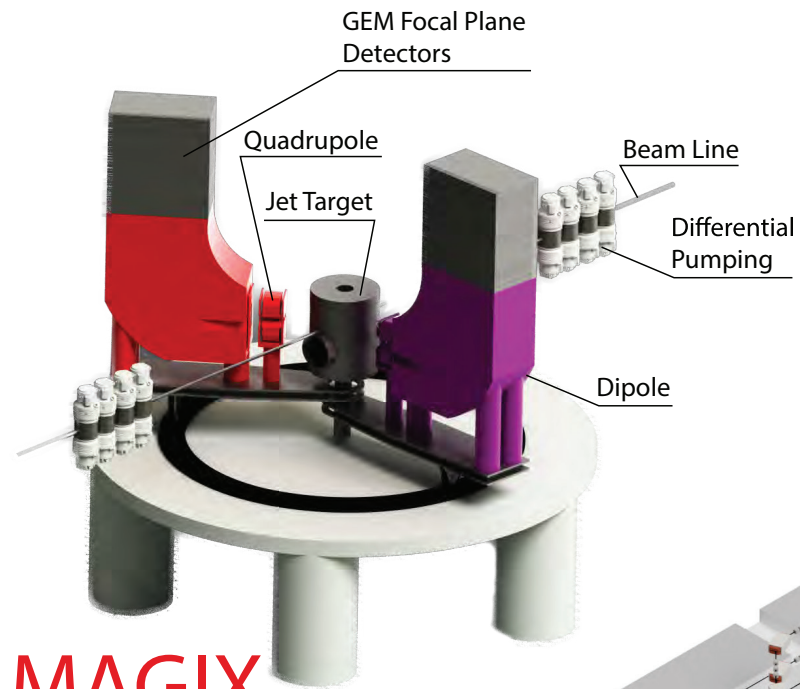


# Stability Requirements

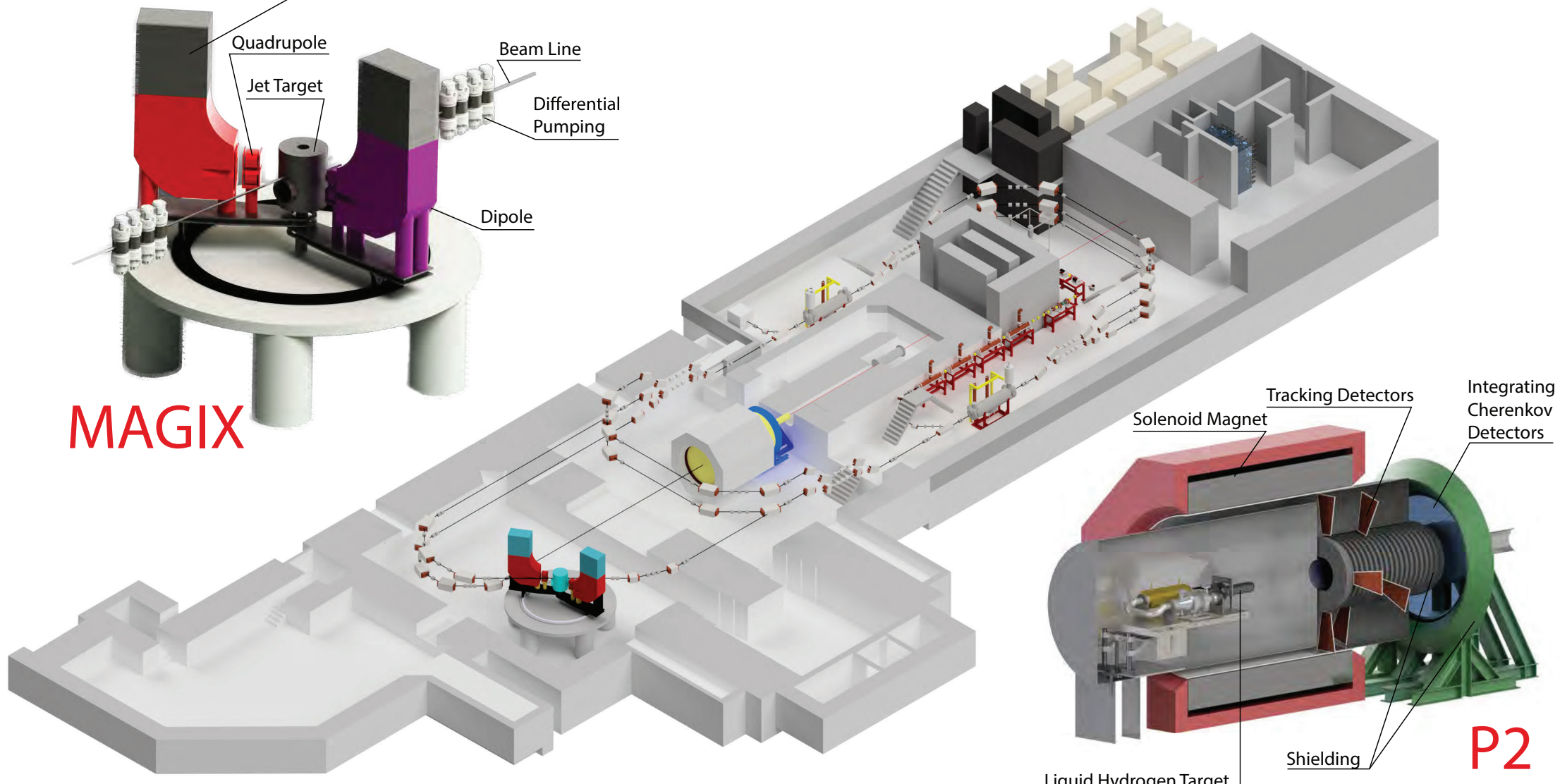
The main worry are beam fluctuations correlated with the helicity:

	Achieved at MAMI	$\sin^2\theta_w$ uncertainty	requirement
• Energy fluctuations:	0.04 eV	< 0.1 ppb	ok!
• Position fluctuations	3 nm	5 ppb	0.13 nm
• Angle fluctuations	0.5 nrad	3 ppb	0.06 nrad
• Intensity fluctuations	14 ppb	4 ppb	0.36 ppb

# MESA



**MAGIX**



Liquid Hydrogen Target

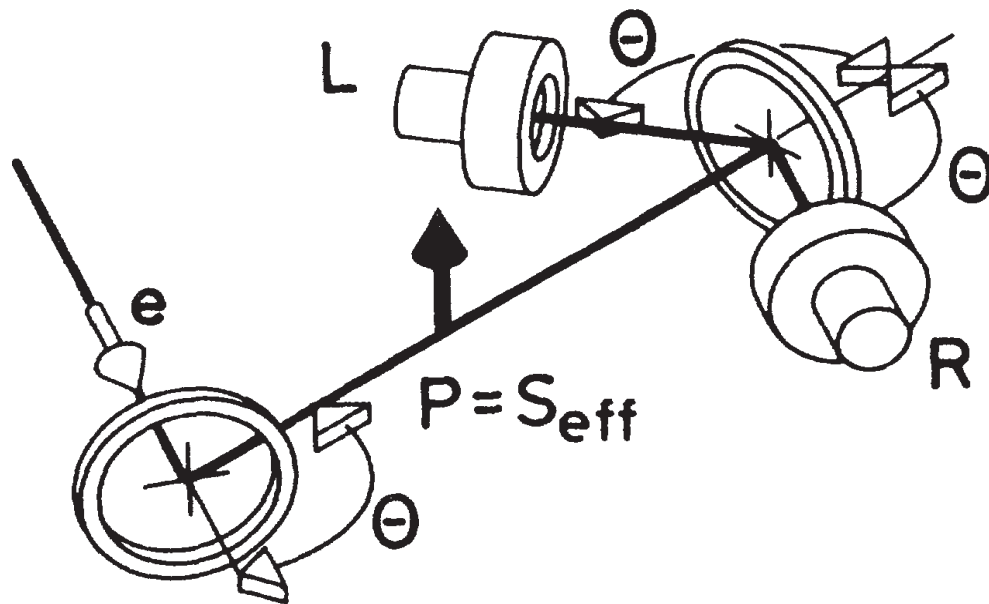
# Polarimetry: Double Mott Polarimeter

Mott Polarimetry:

- Measure left/right asymmetry to obtain spin polarisation
- Analysing power of foils needs to be extrapolated

Double Mott Polarimeter:

- Obtain analysing power from measurement
- Precise measurement of spin polarisation
- Invasive measurement at source

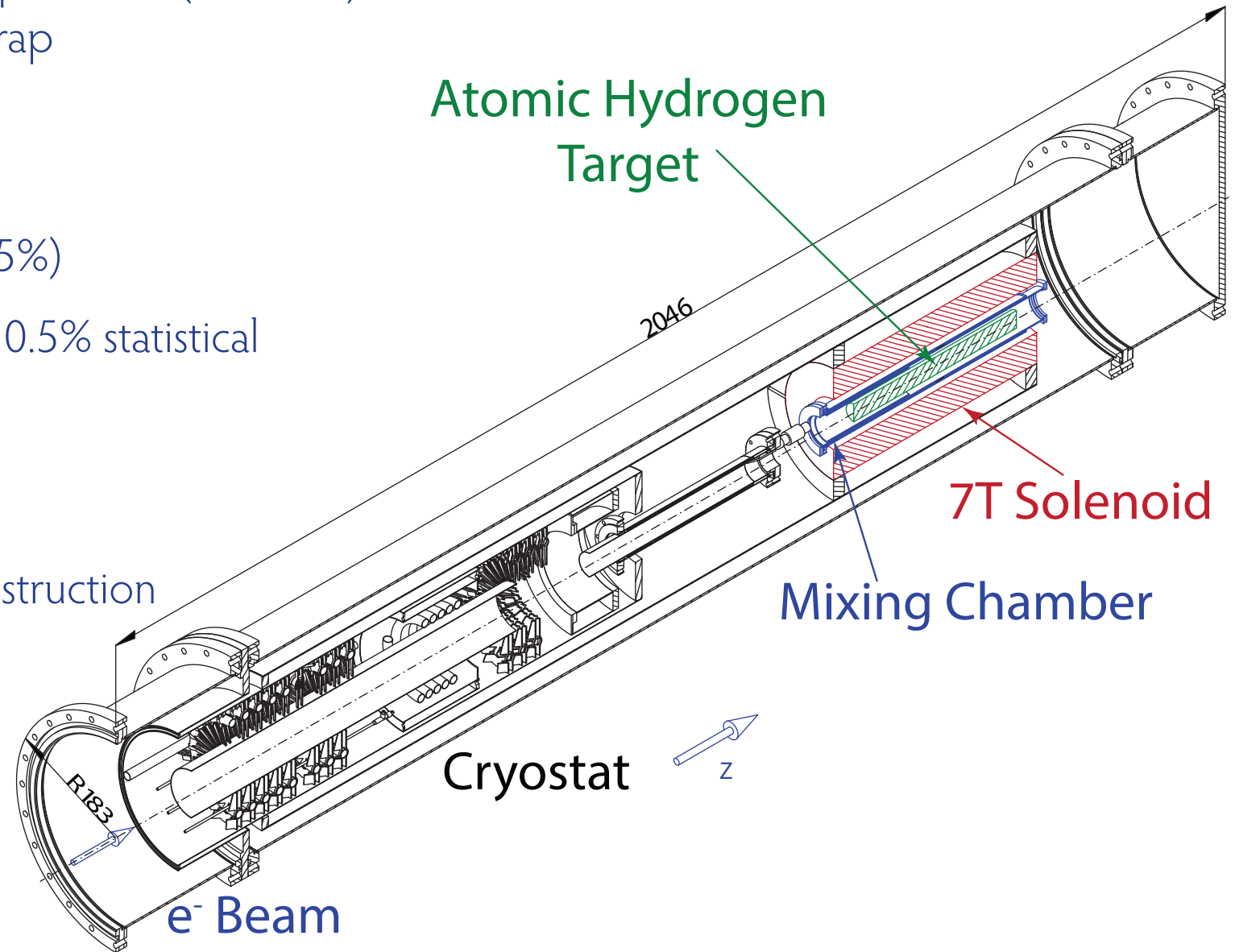


[Gellrich and Kessler, Phys.Rev.A. 43, 204 (1991)]

# Polarimetry: Hydro-Møller Polarimeter

Møller scattering from polarized (8 T field)  
atomic hydrogen in a trap

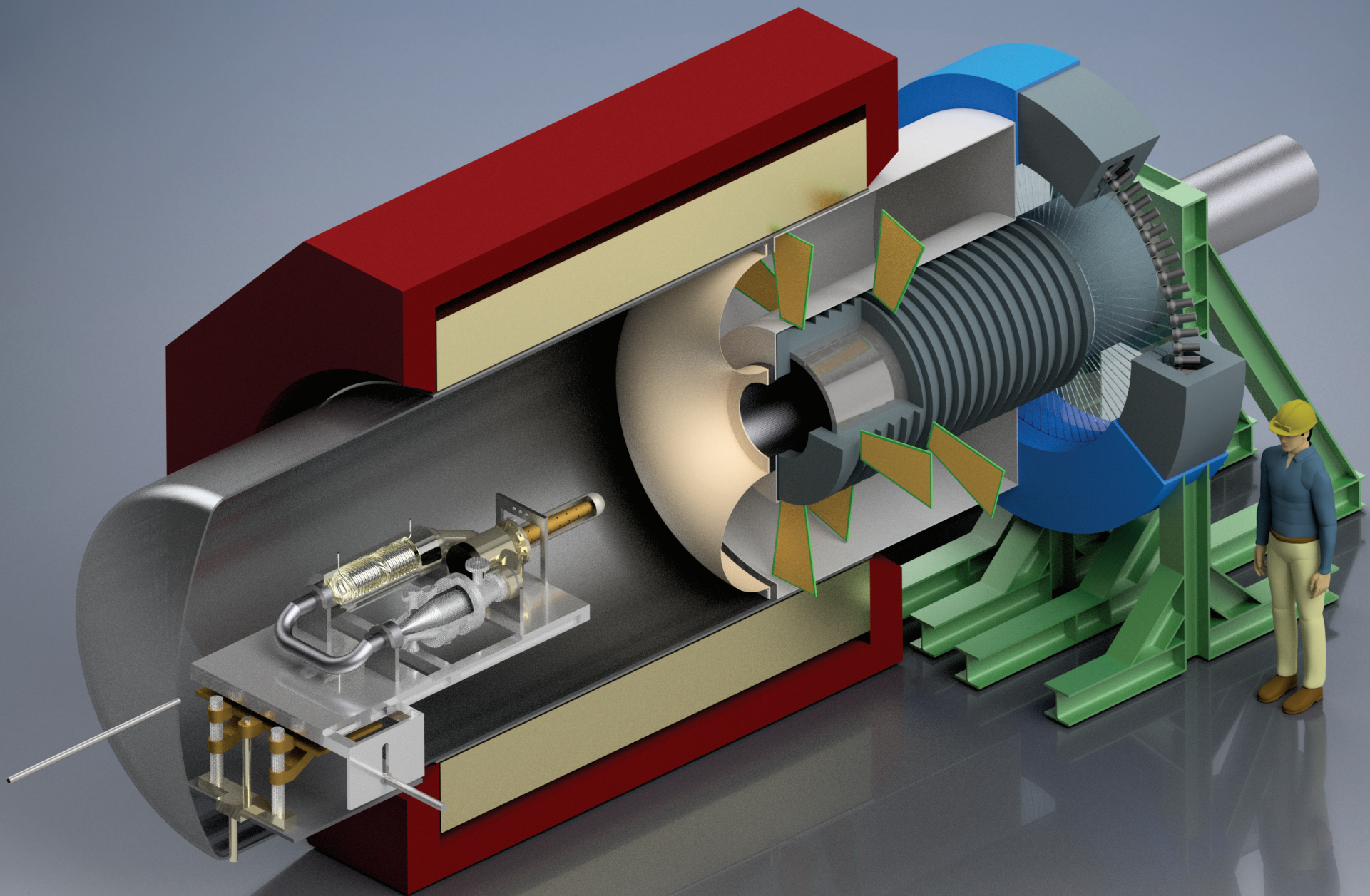
- Online capability
- High accuracy ( $< 0.5\%$ )
- About 2 h to reach 0.5% statistical accuracy
- Cryostat under construction in Mainz



P2:

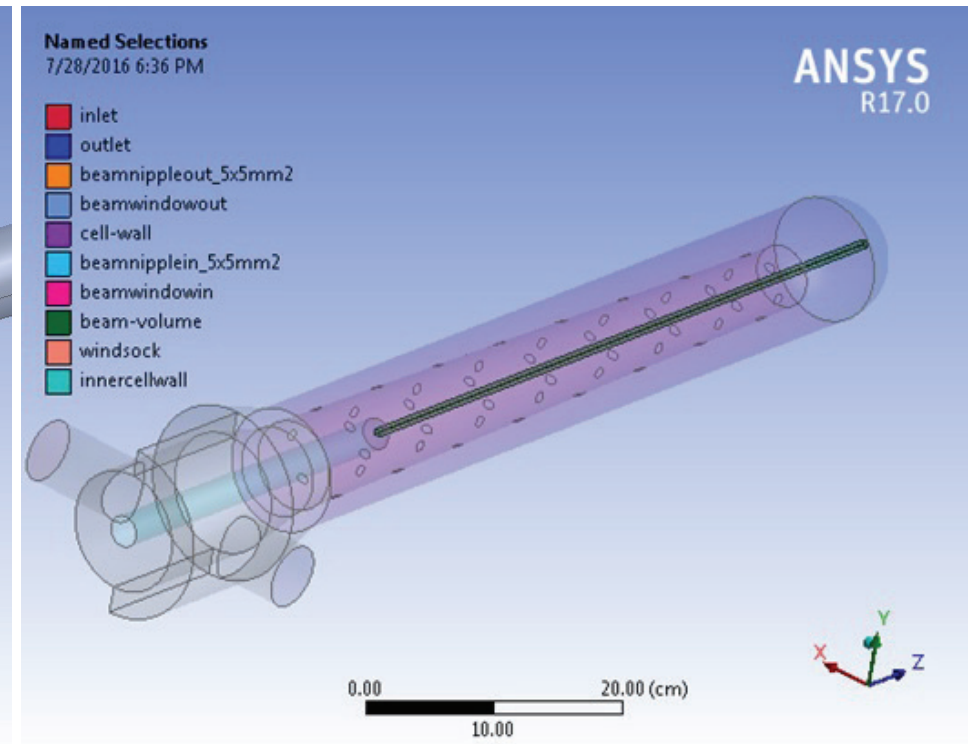
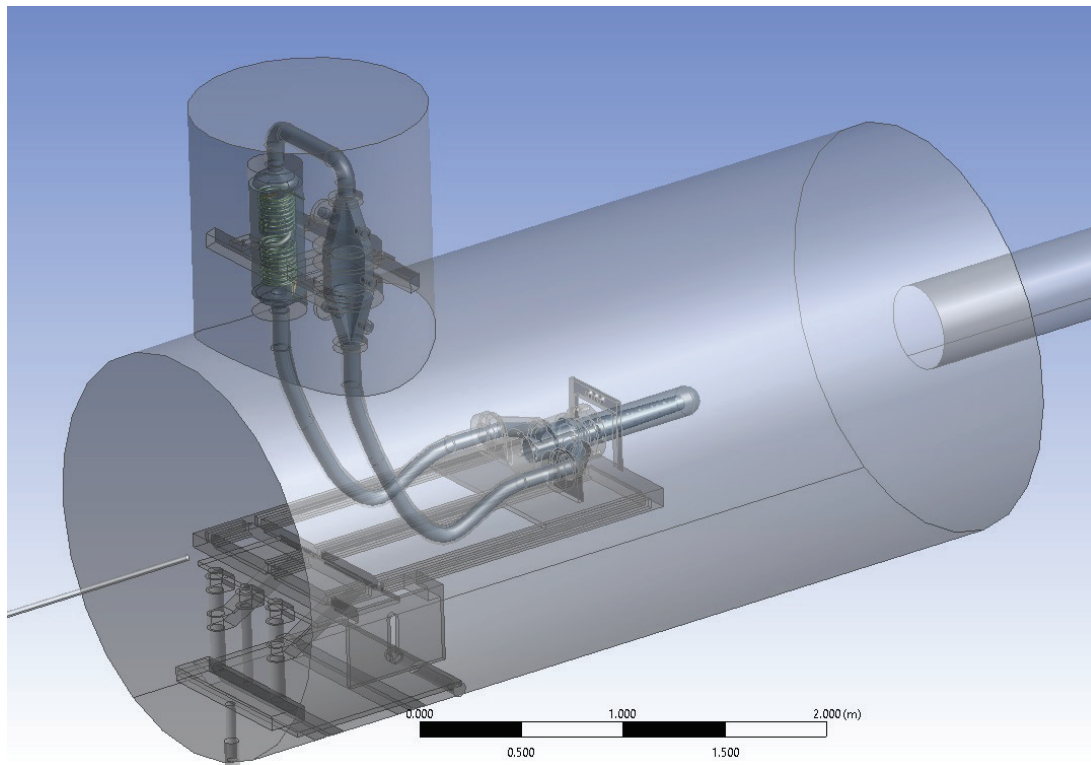
How to detect 100 GHz of (the right) electrons...



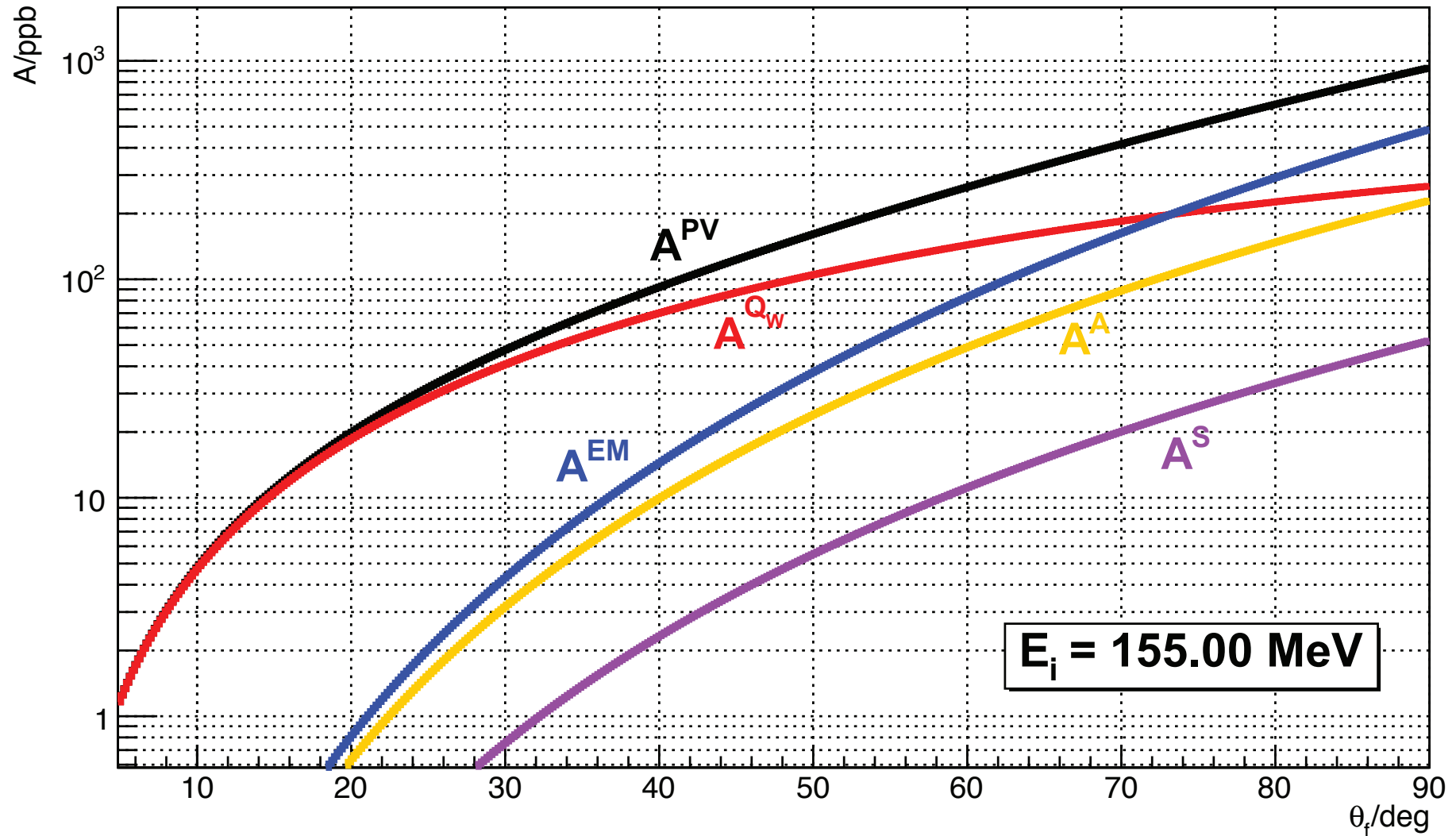


# Target

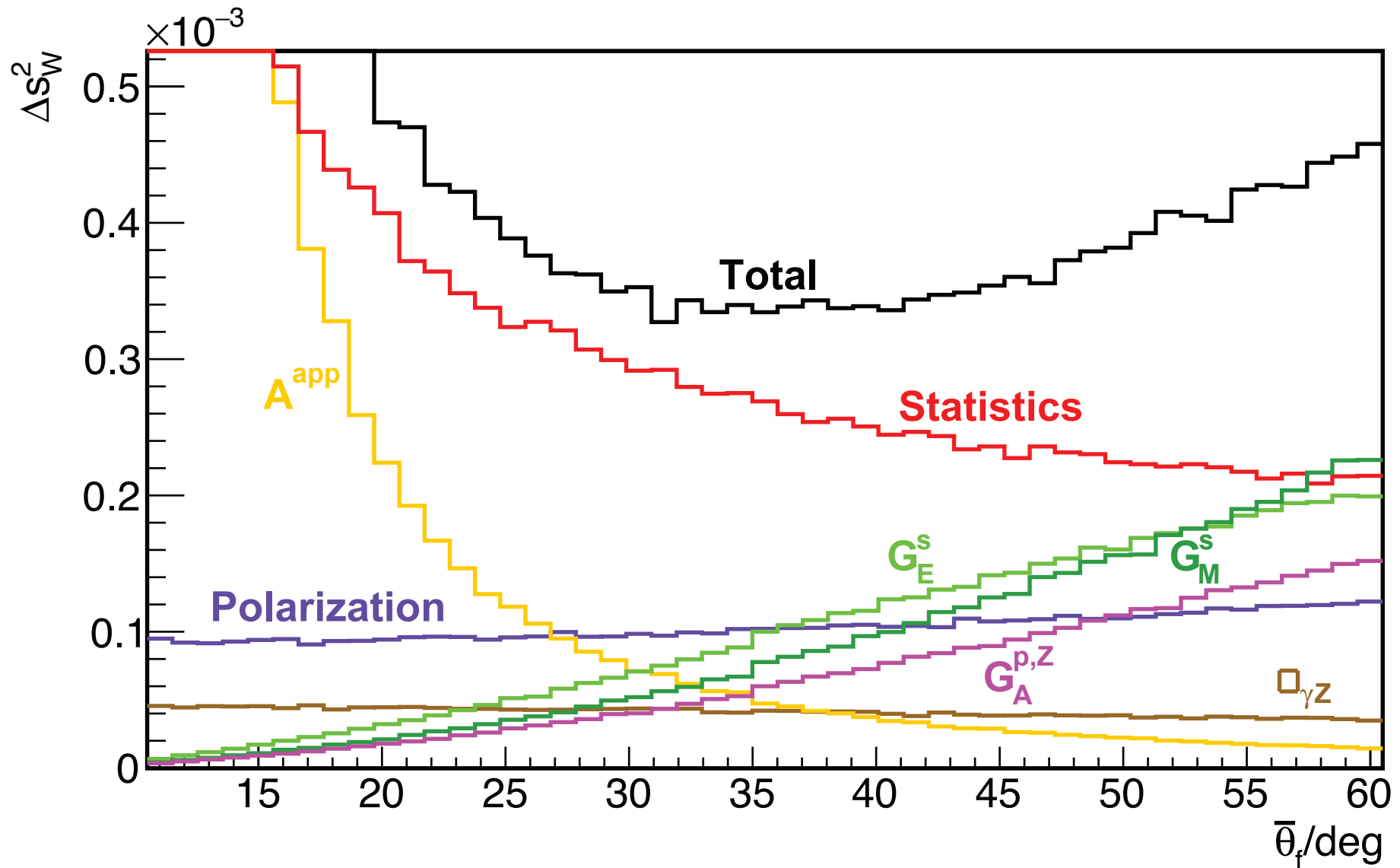
- 60 cm of liquid hydrogen
- 3.1 KW beam power deposited
- Should not boil...
- Challenging design using CFD tools (Silviu Covrig, JLab)



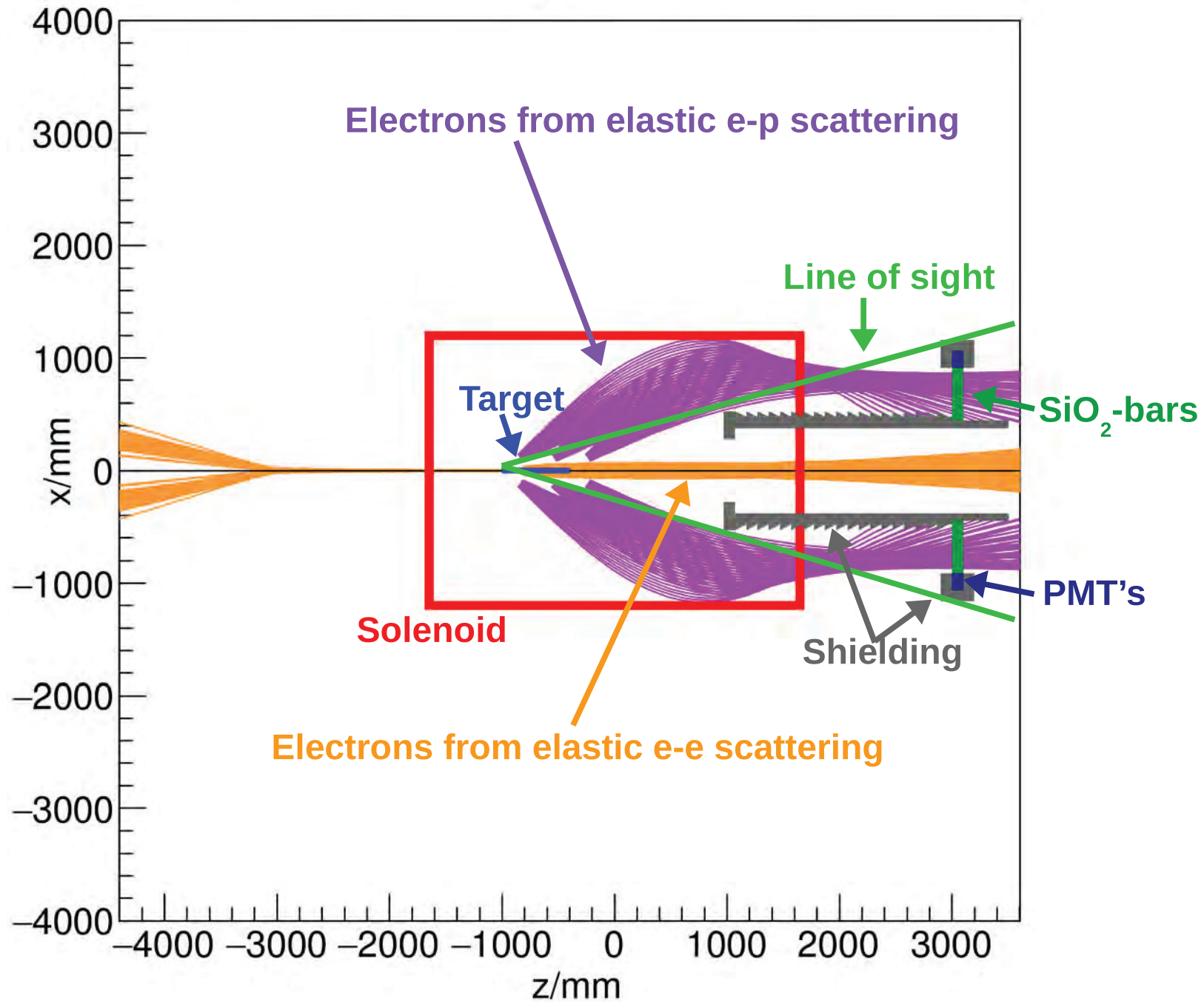
# Choice of scattering angle



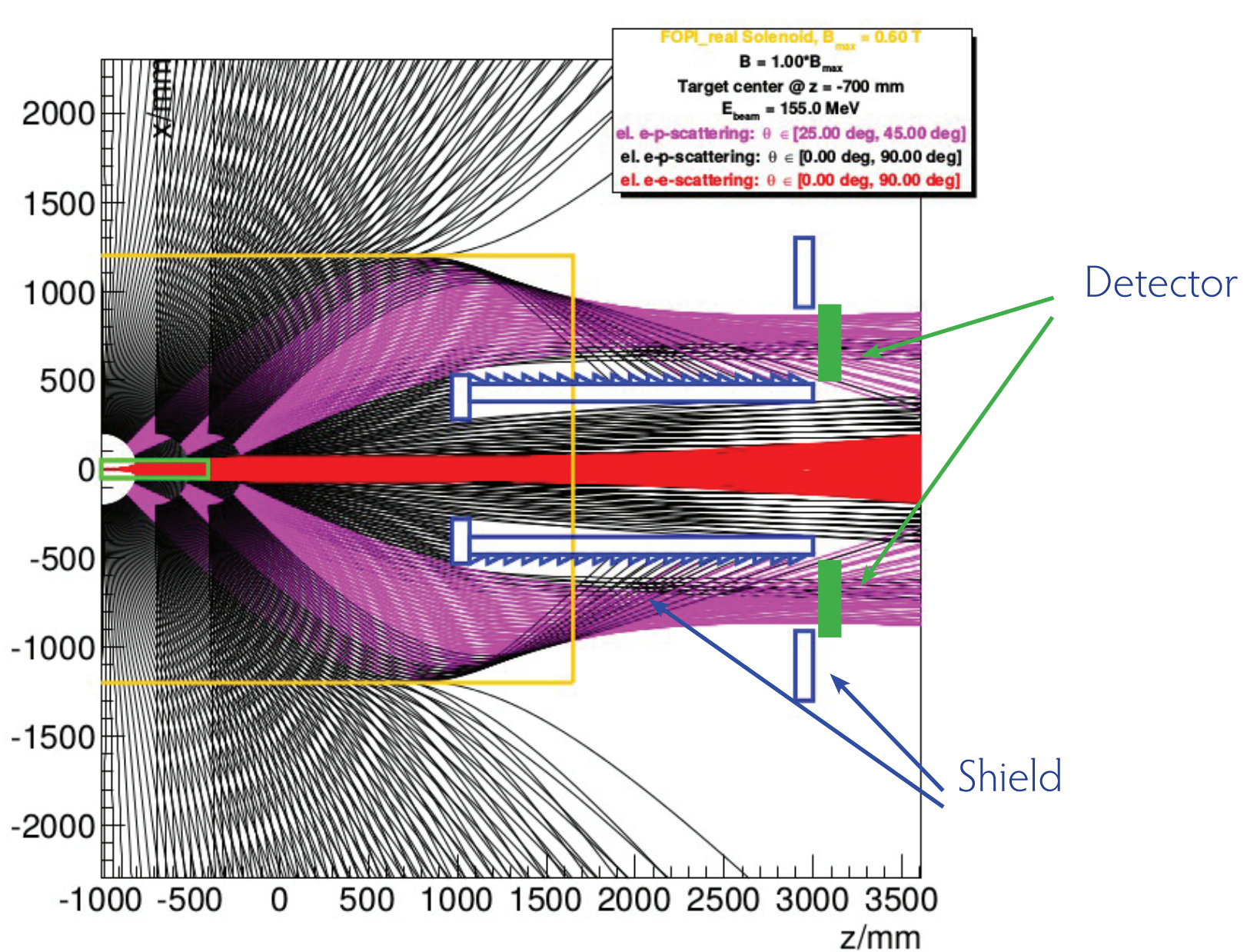
# Choice of scattering angle



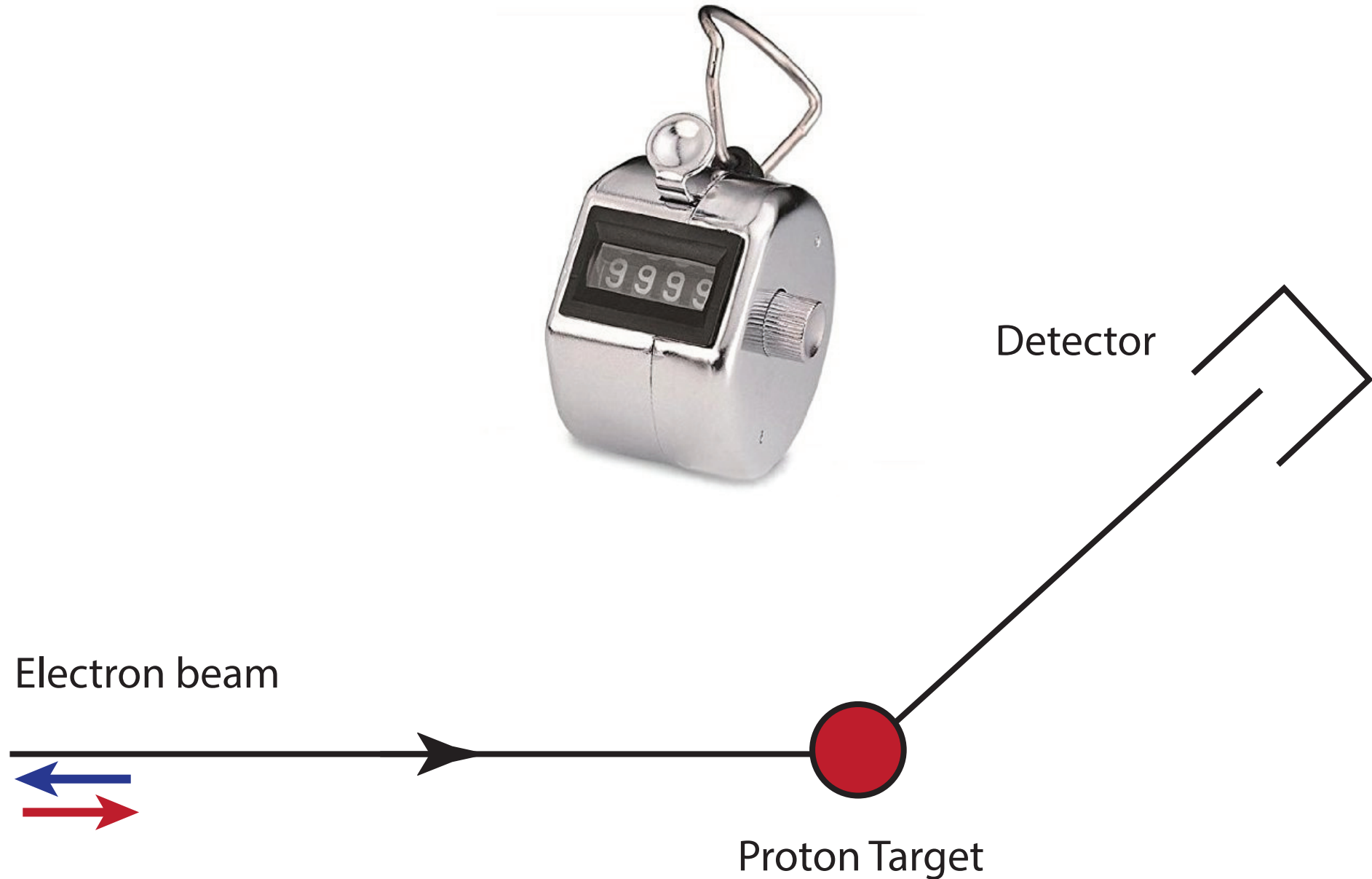
# Solenoid spectrometer



# Solenoid spectrometer



# Counting detectors



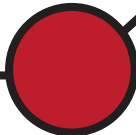
# Integrating detectors



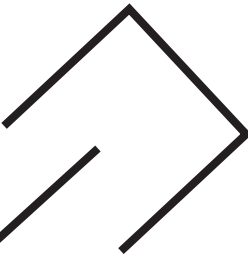
Electron beam



Proton Target

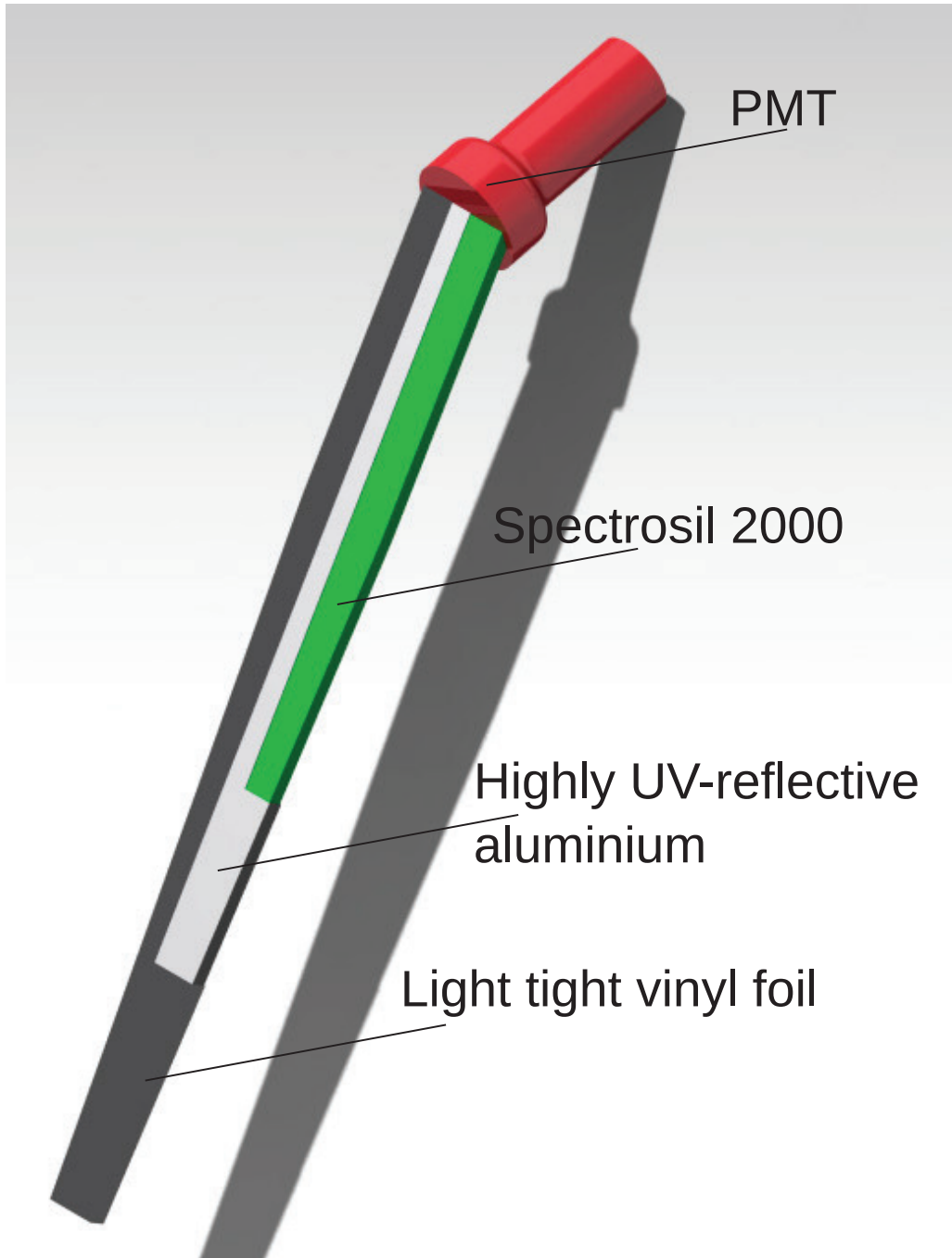


Detector



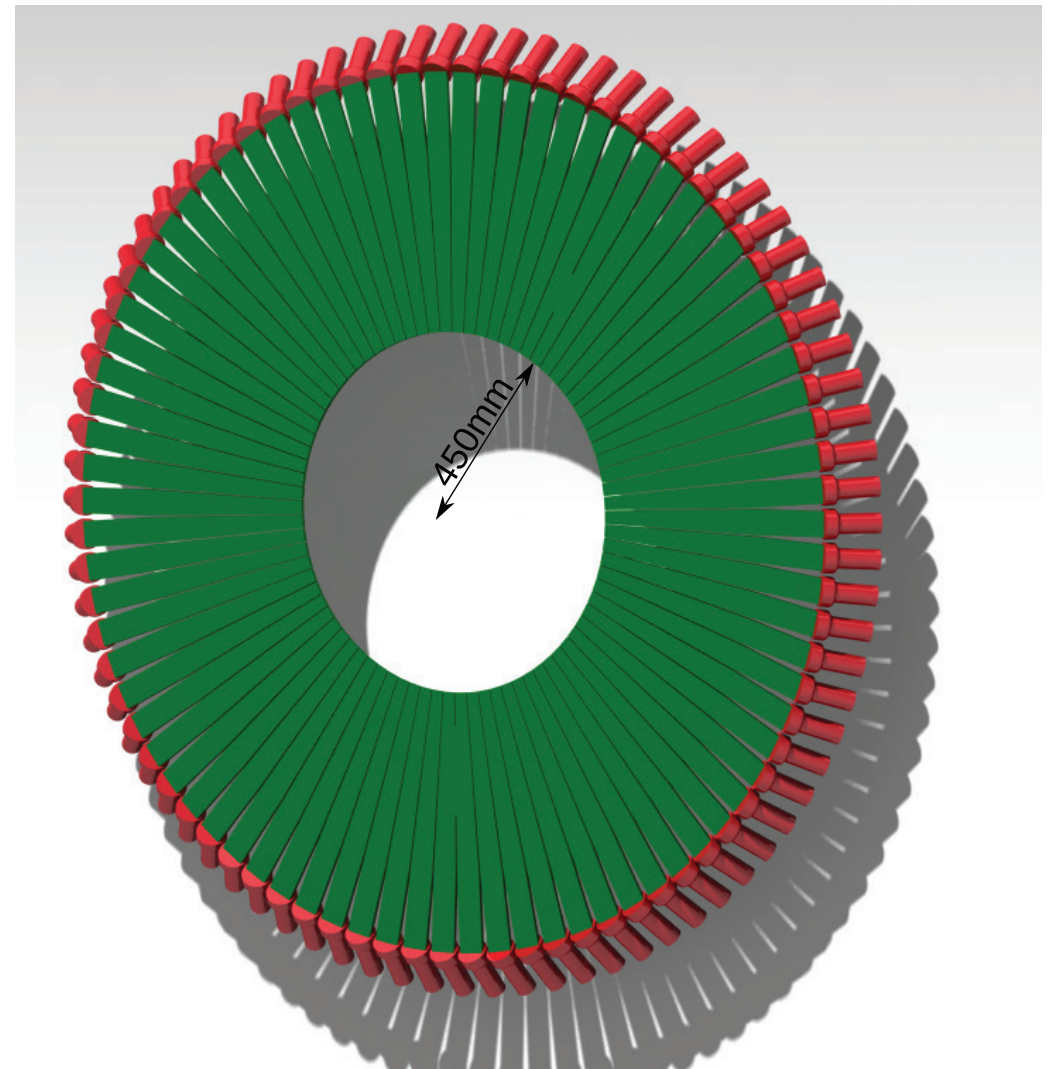


# Quartz-Bars & Photomultipliers

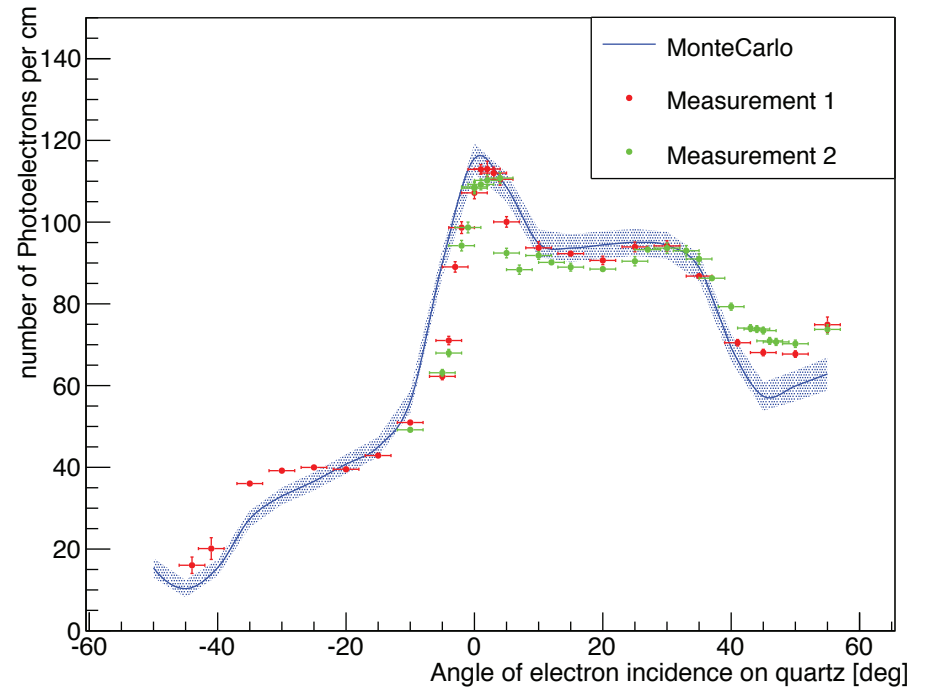
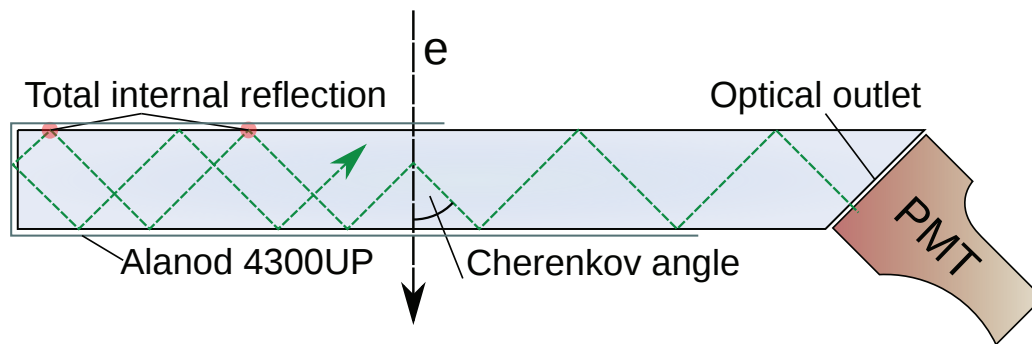
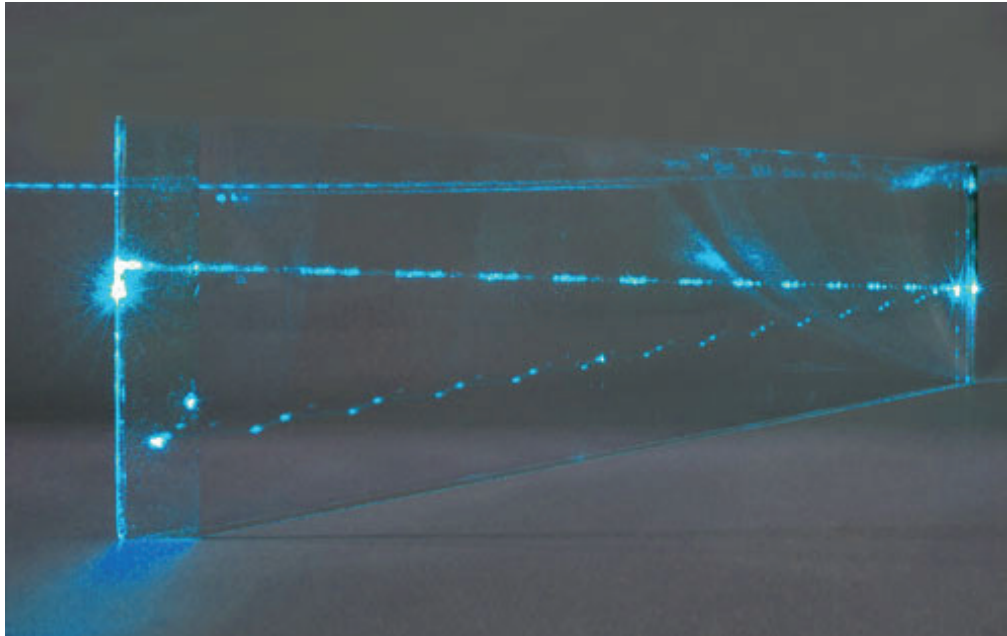


Detect Cherenkov-light created by electrons

Integrate photomultiplier current



# Quartz-Bars & Photomultipliers

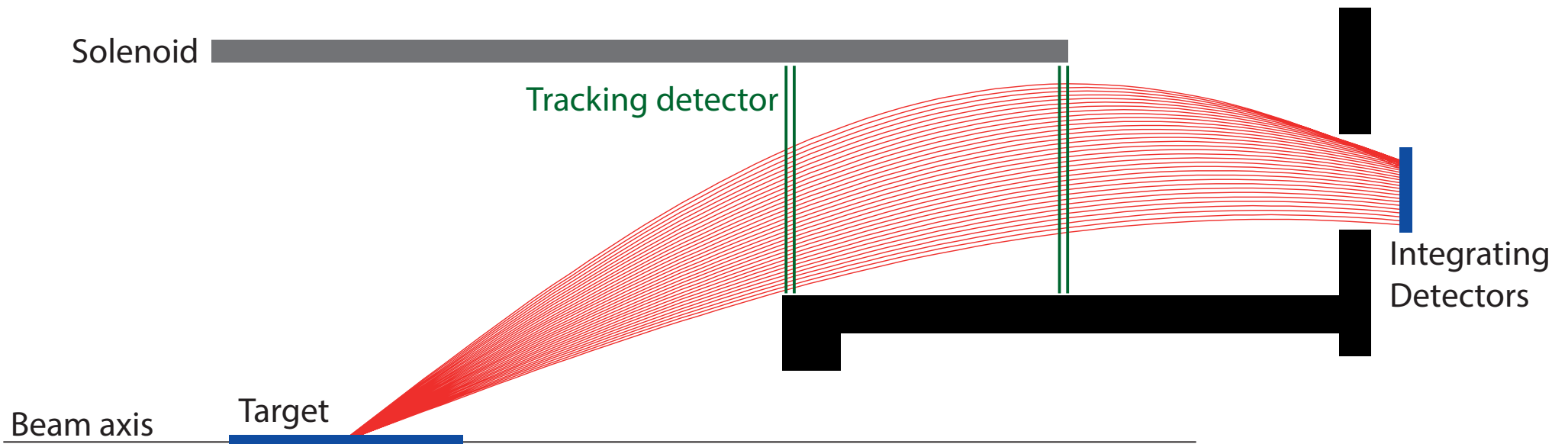


Measuring  $Q^2$ :

Tracking a  
lot of low momentum particles

# Tracker requirement

- Low momentum electrons:  
Thin detectors
- Very high rates:  
Fast and granular detectors

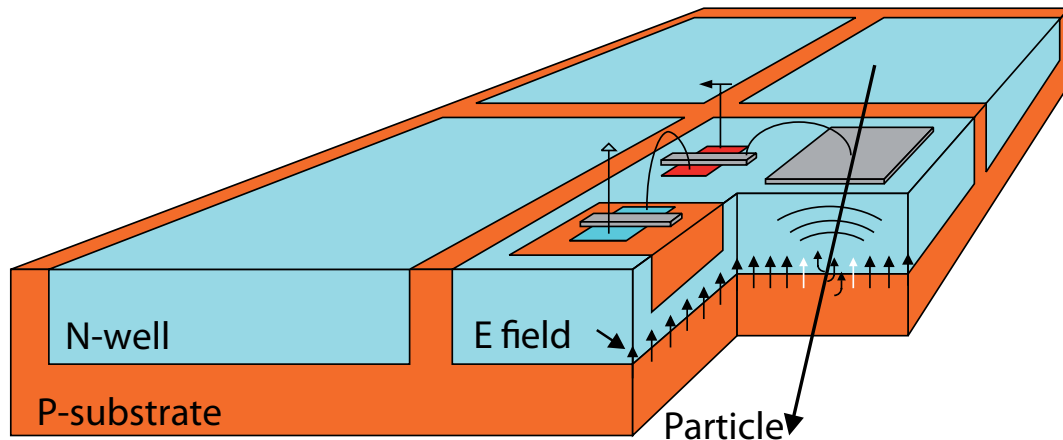


Fast, thin, cheap pixel sensors

High Voltage Monolithic Active Pixel Sensors

# Fast and thin sensors: HV-MAPS

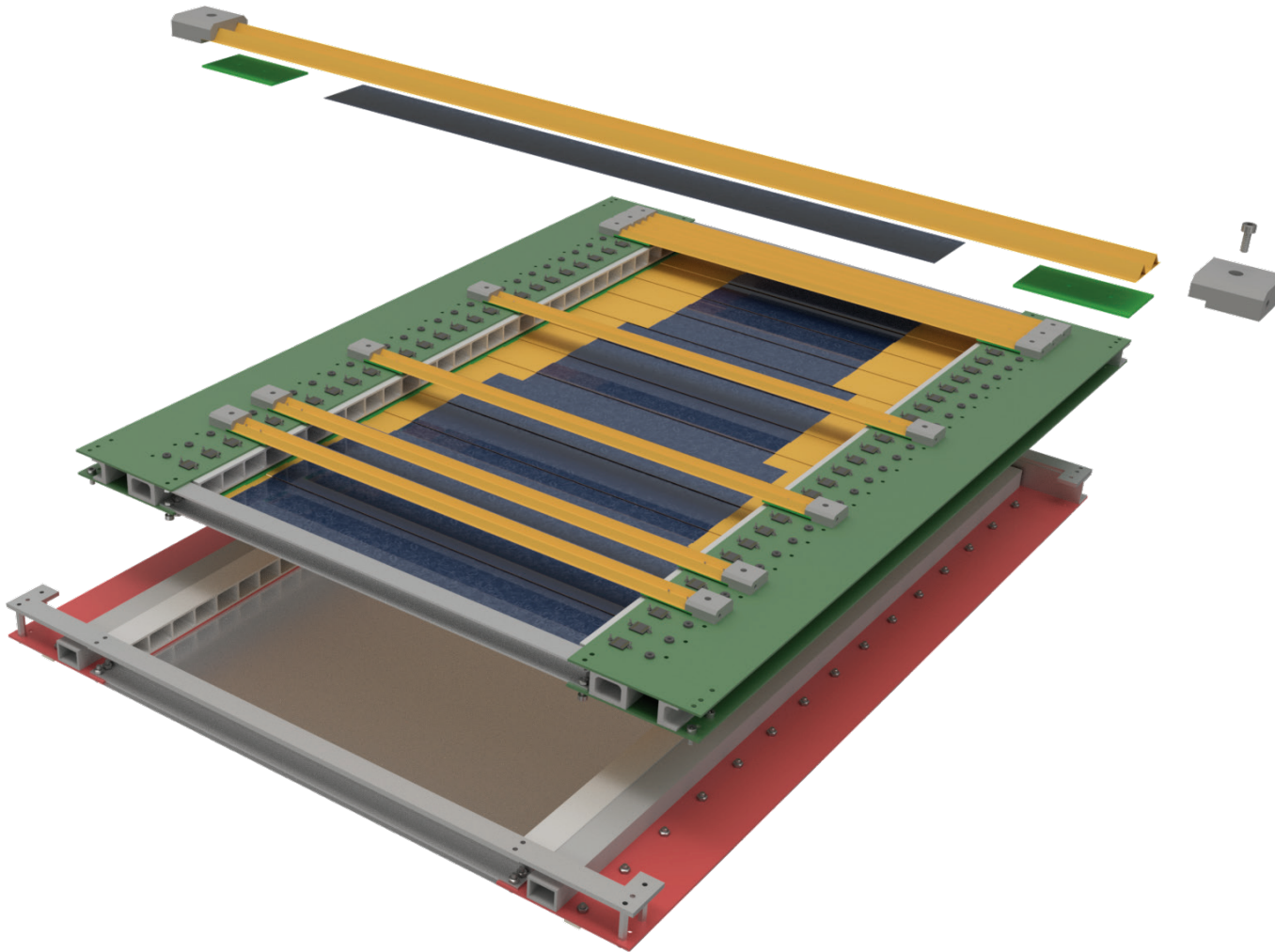
High voltage monolithic active pixel sensors - Ivan Perić (morning sessions!)



- Use a high voltage commercial process (automotive industry)
- Small active region, fast charge collection via drift
- Implement logic directly in N-well in the pixel - smart diode array
- Can be thinned down to  $< 50 \mu\text{m}$
- Logic on chip: Output are zero-suppressed hit addresses and timestamps

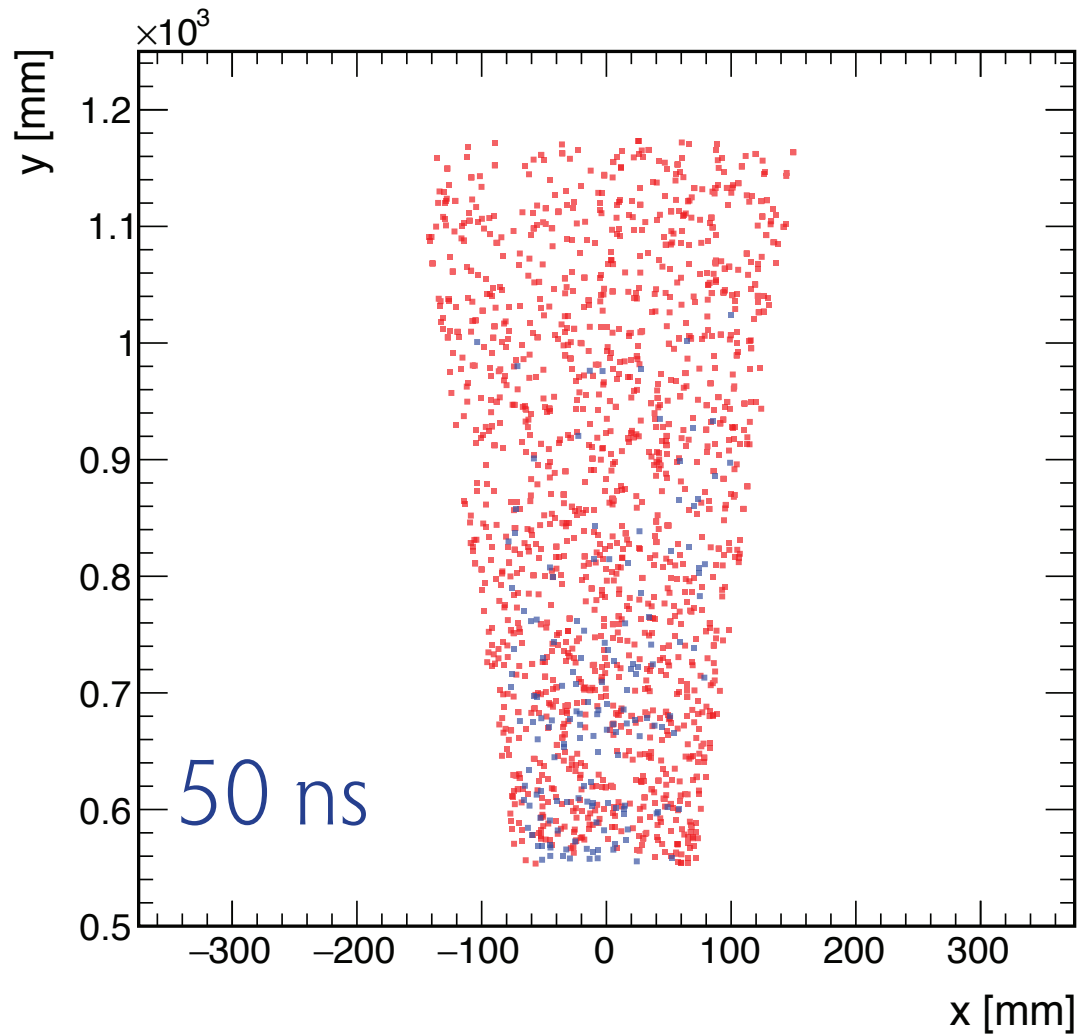
(I.Perić, P. Fischer et al., NIM A 582 (2007) 876 )

# Mechanics



- 50  $\mu\text{m}$  silicon
- 25  $\mu\text{m}$  Kapton™ flexprint with aluminium traces
- 25  $\mu\text{m}$  Kapton™ frame as support
- About 1‰ of a radiation length per layer

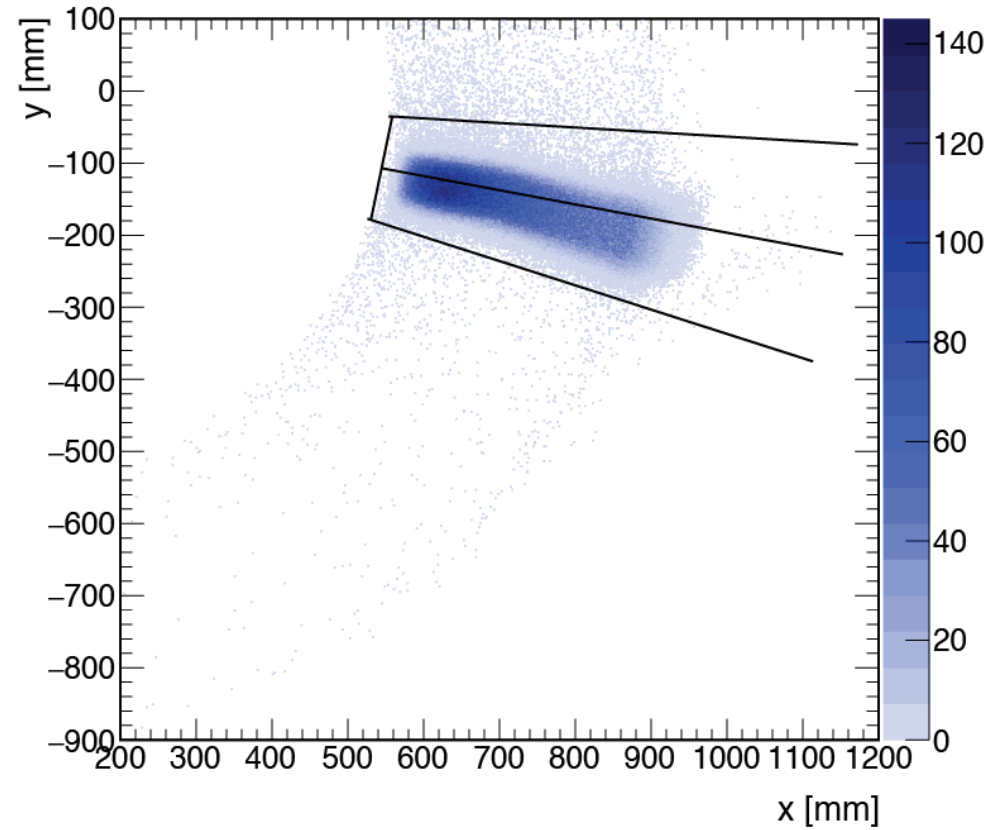
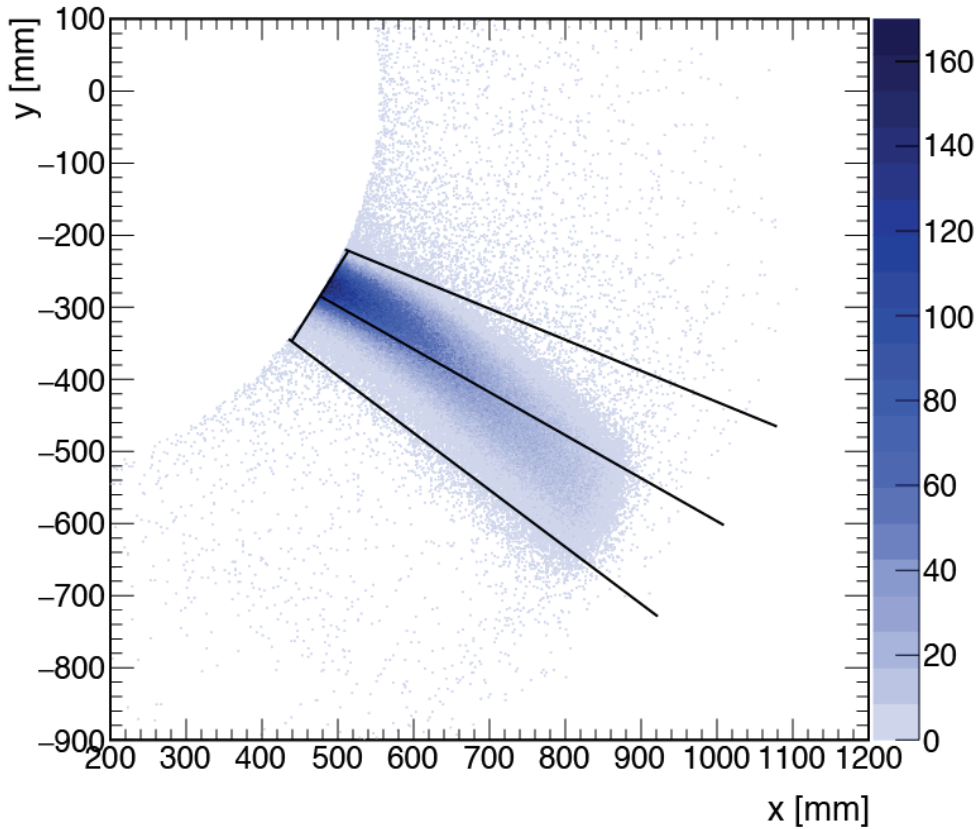
# The tracking challenge



- 100 GHz electrons
- about 1000 Bremsstrahlung photons per electron

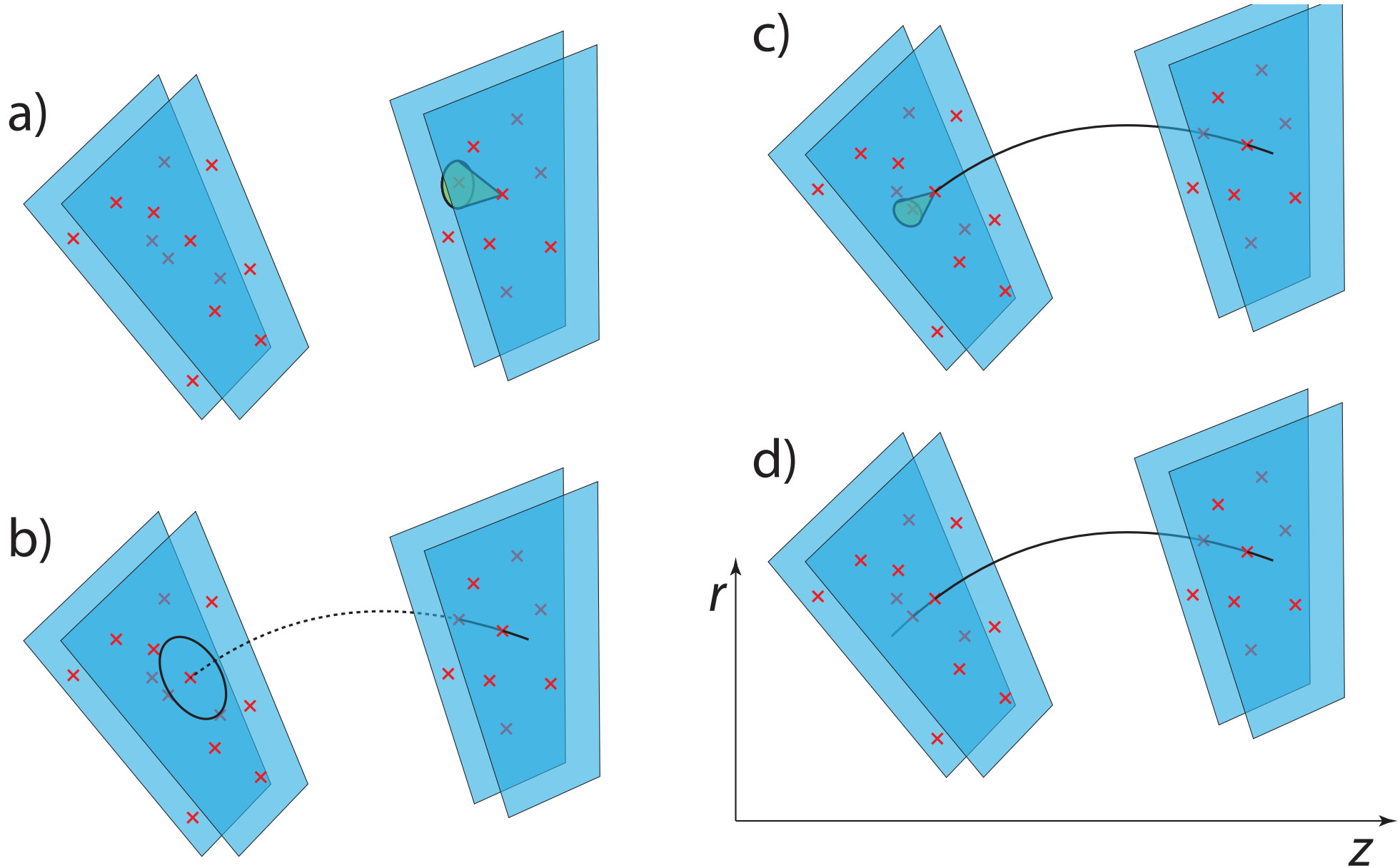


# No need for full coverage

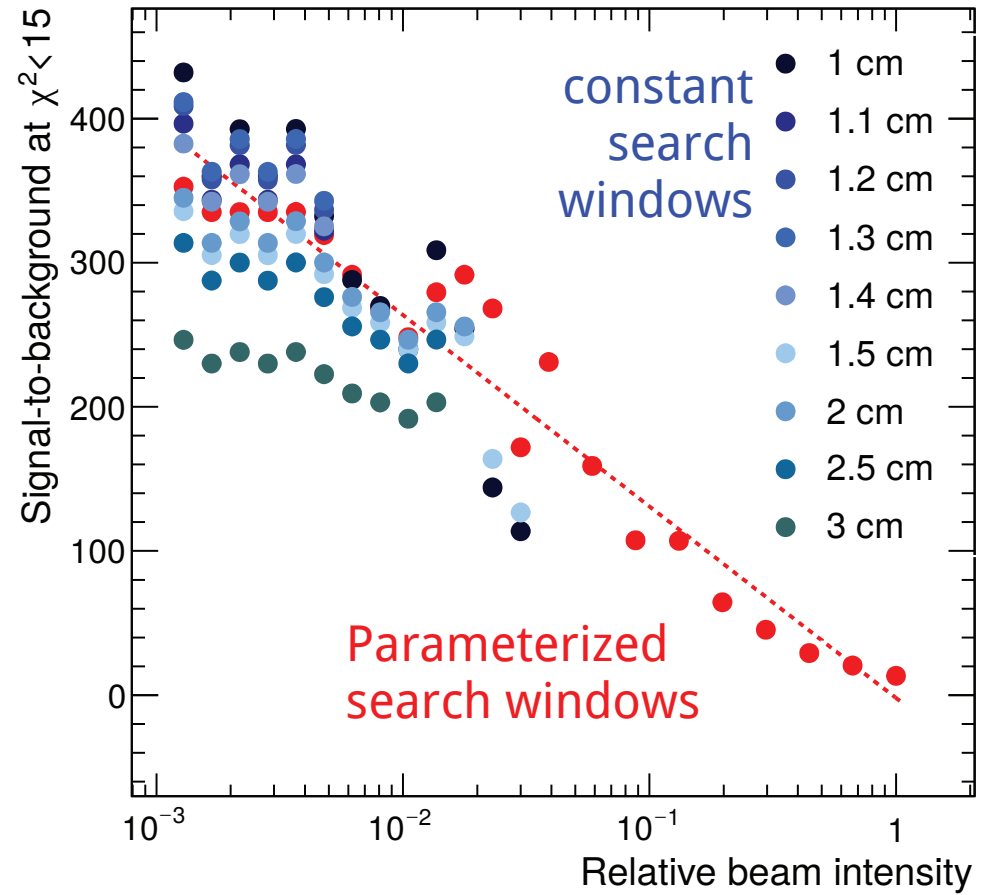
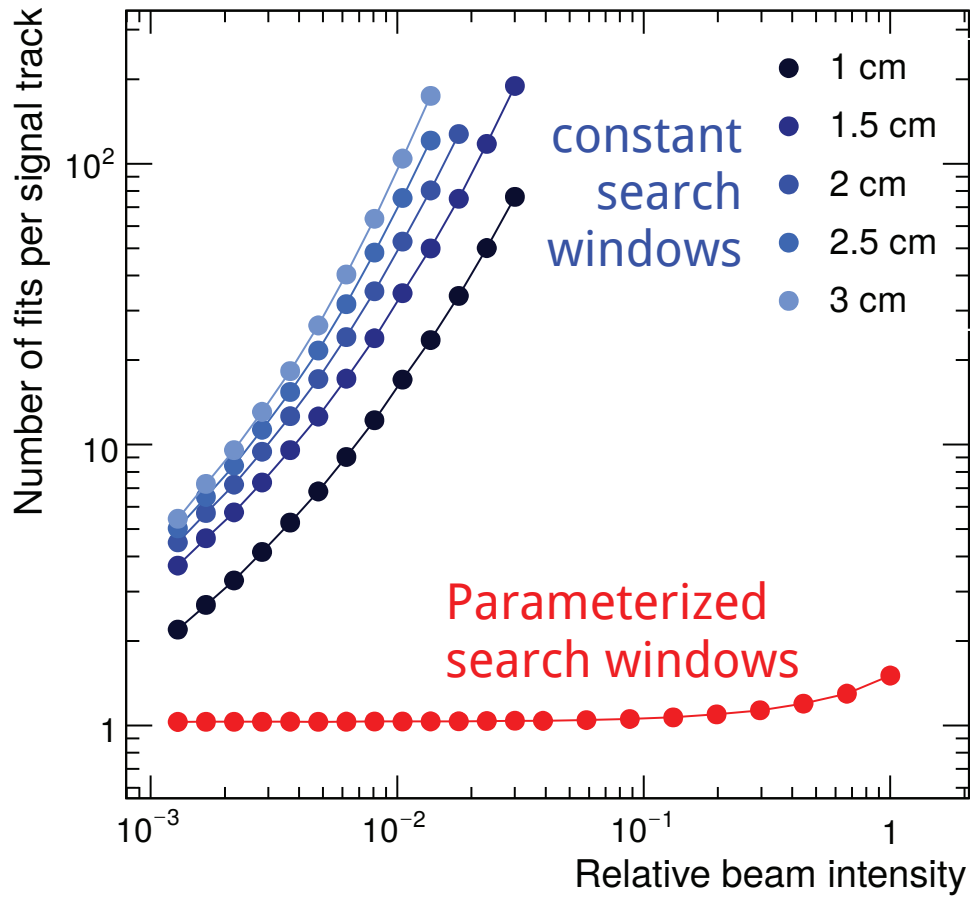


# Parametrization based tracking

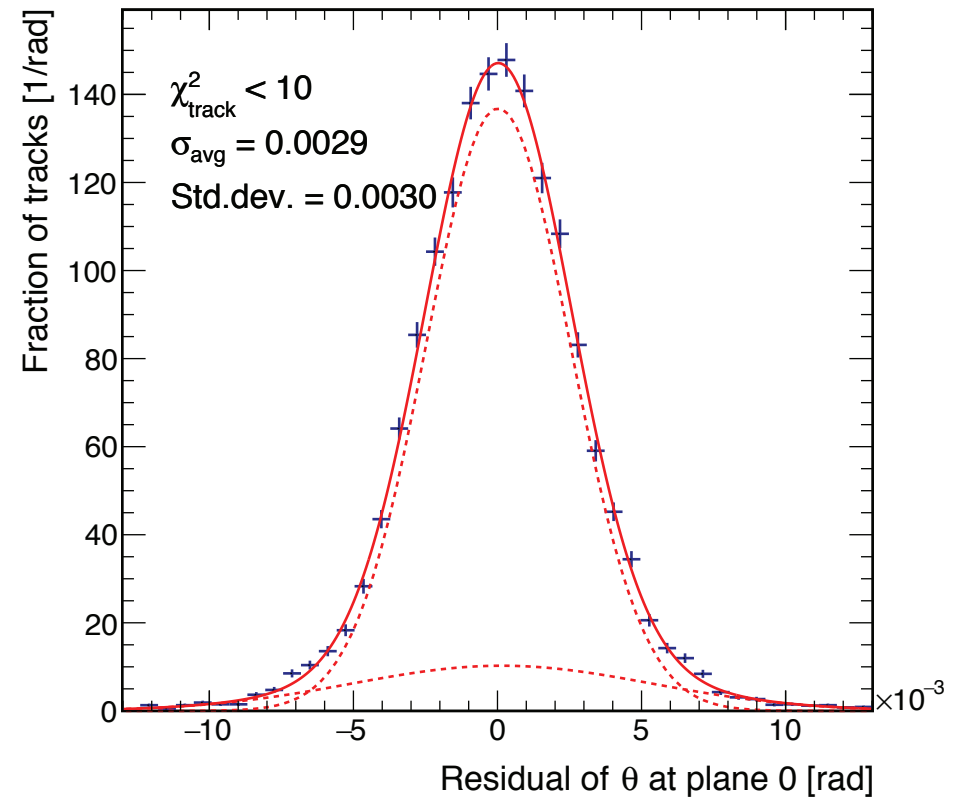
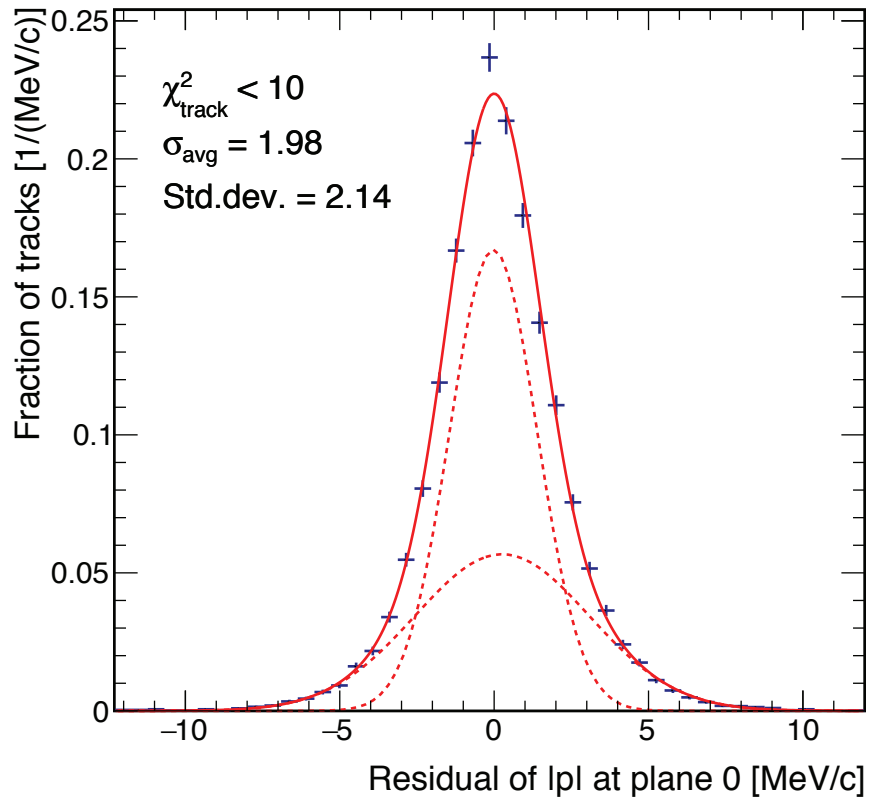
I. Sorokin



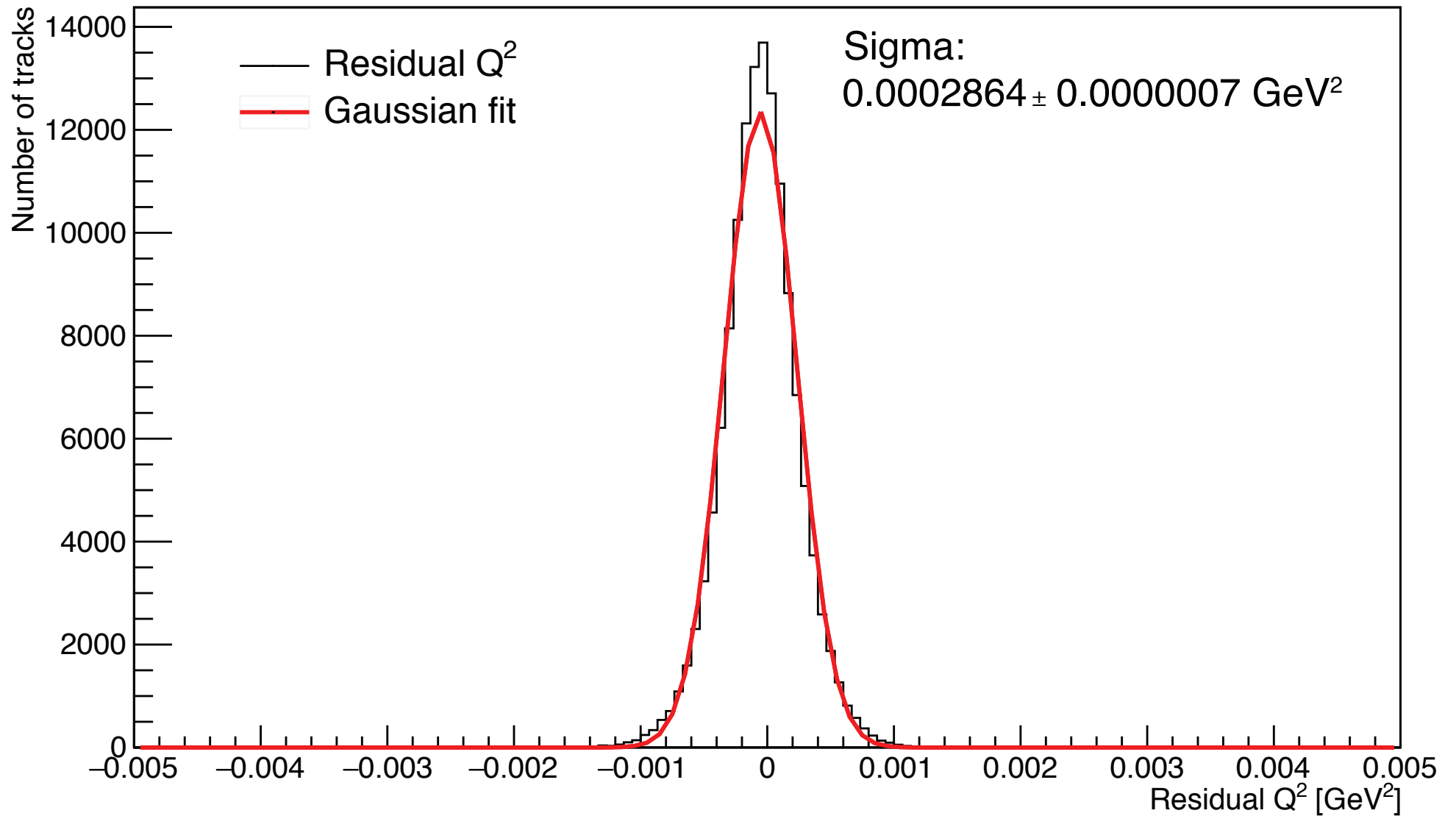
# Tracking performance

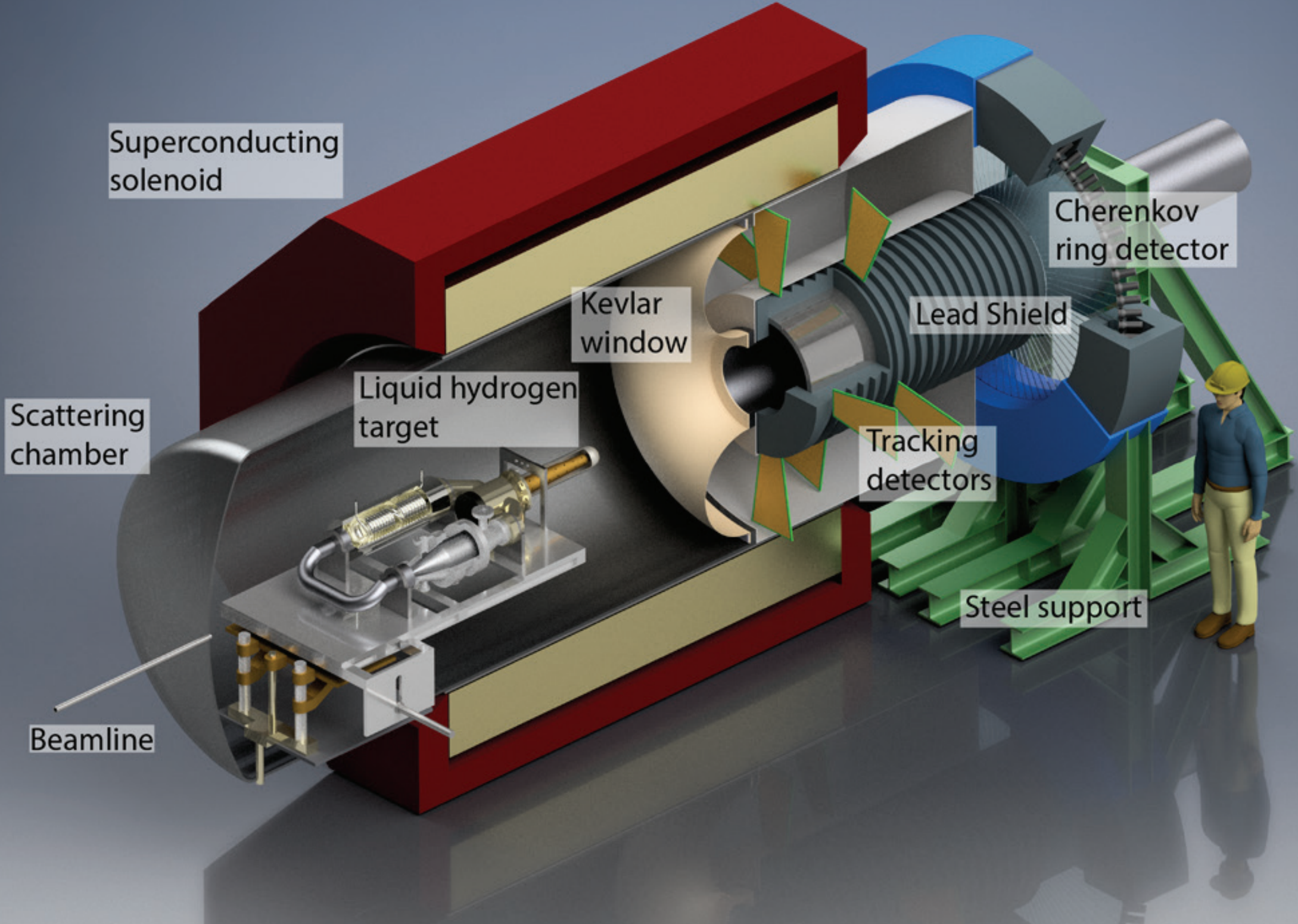


# Tracking performance



# $Q^2$ reconstruction





Superconducting solenoid

Scattering chamber

Beamline

Liquid hydrogen target

Kevlar window

Lead Shield

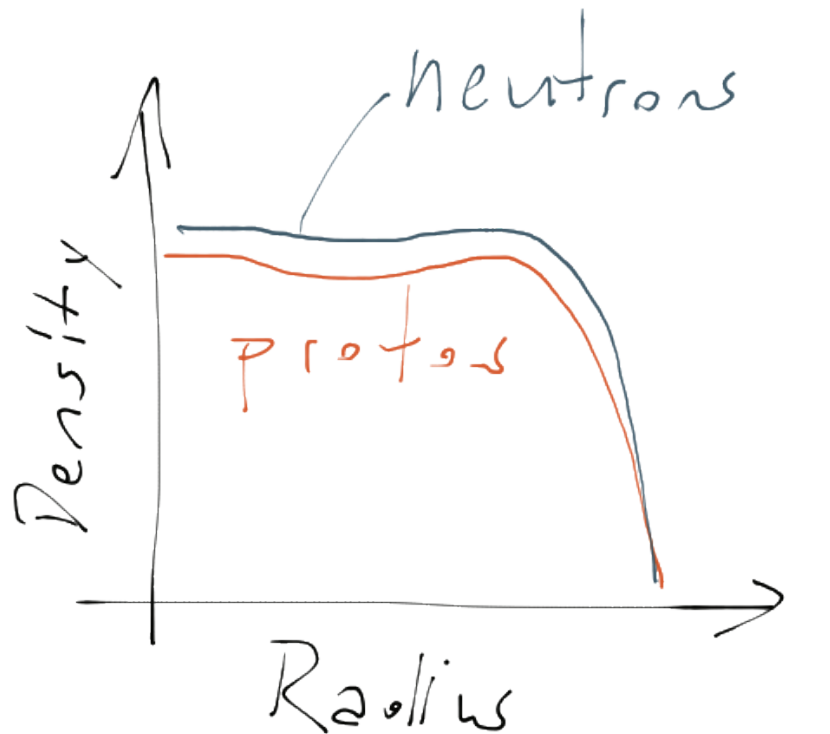
Tracking detectors

Cherenkov ring detector

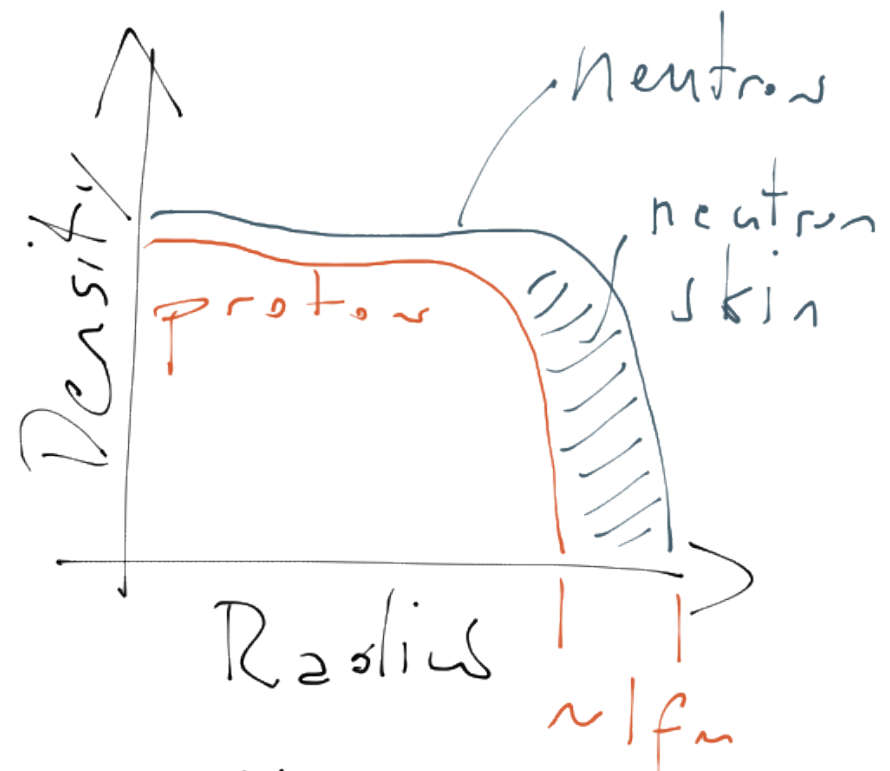
Steel support

# Neutron Skins

Where are the neutrons in the nucleus?



Balanced Nucleus

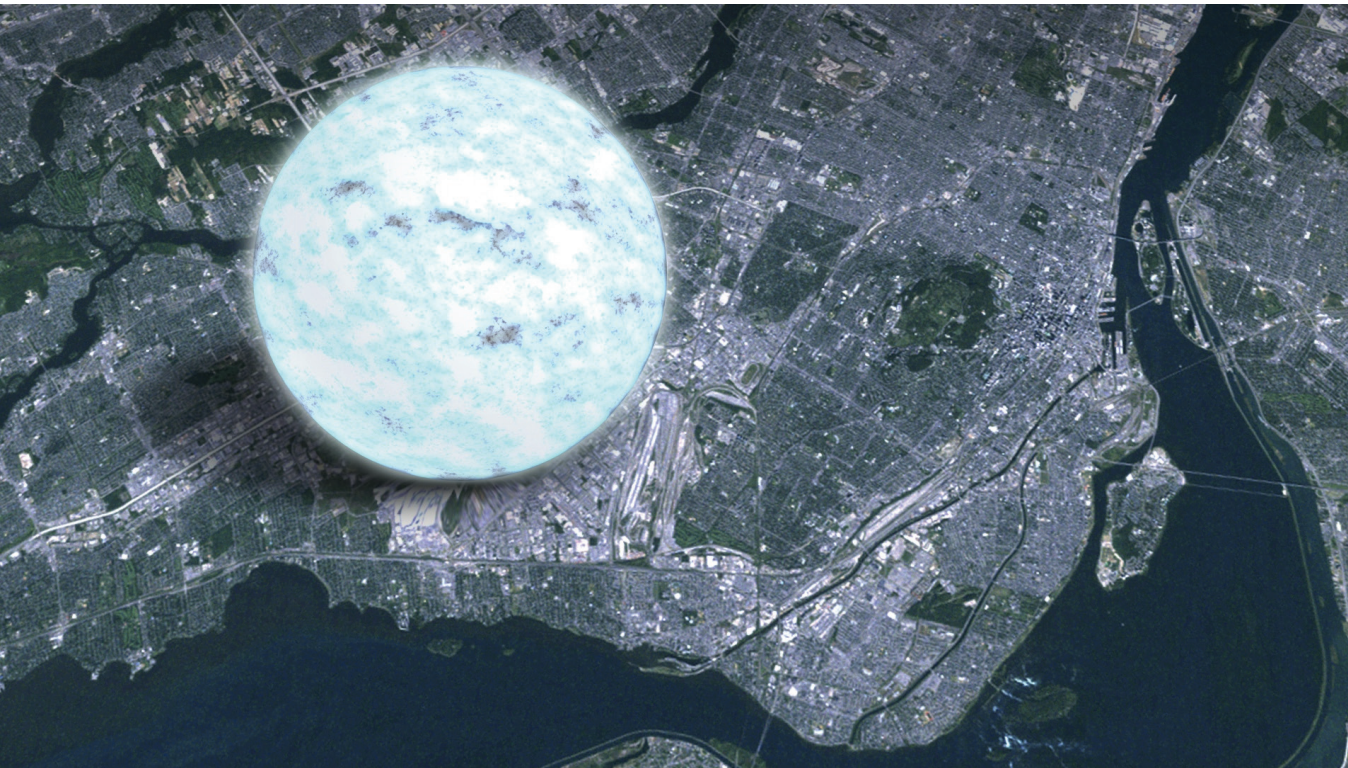
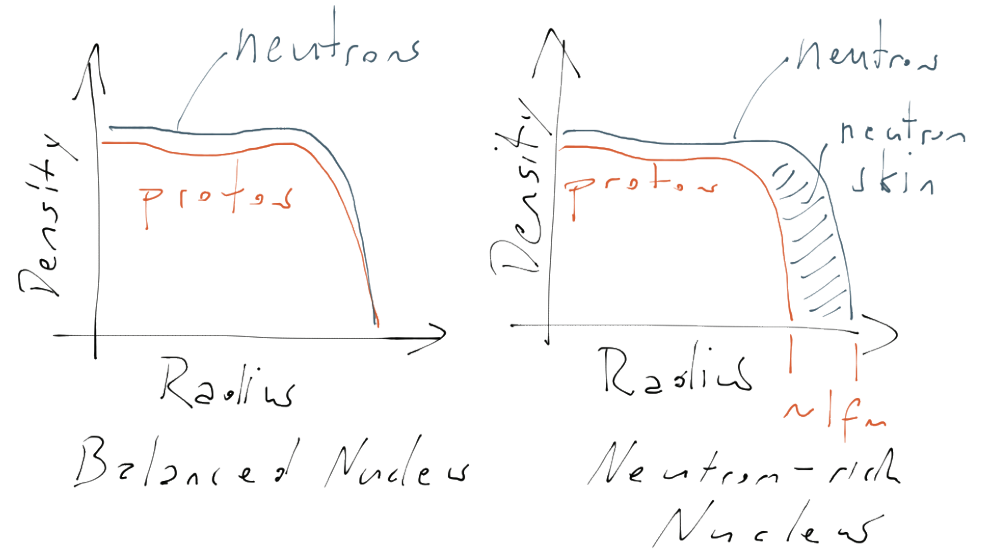


Neutron-rich Nucleus

# Neutron Skins

Where are the neutrons in the nucleus?

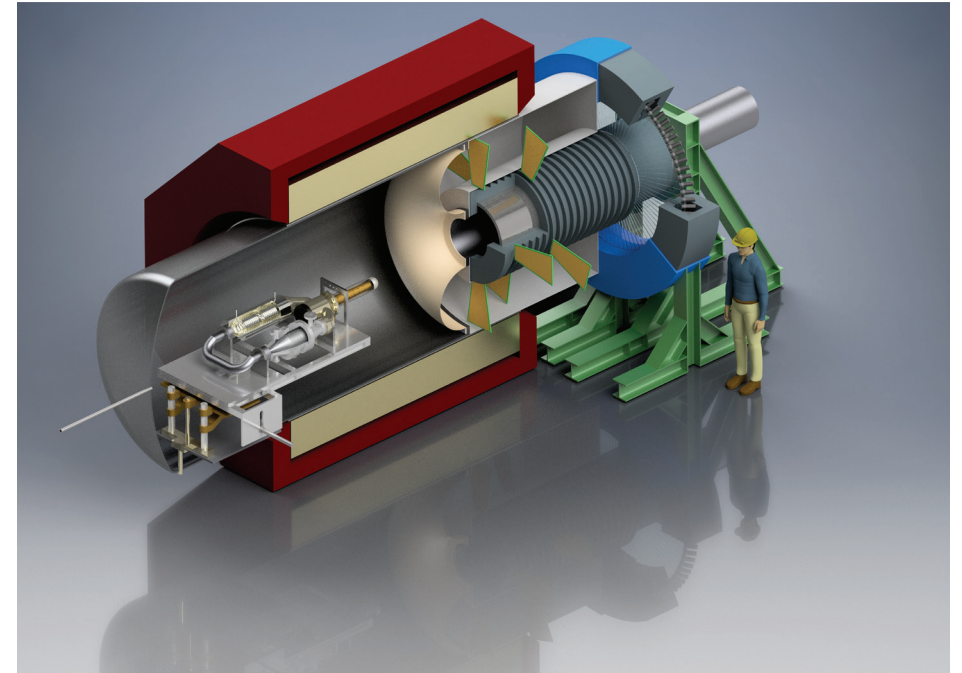
- Gives access to the equation of state of neutron matter
- Tells us how big/small neutron stars are





# How to see the neutrons?

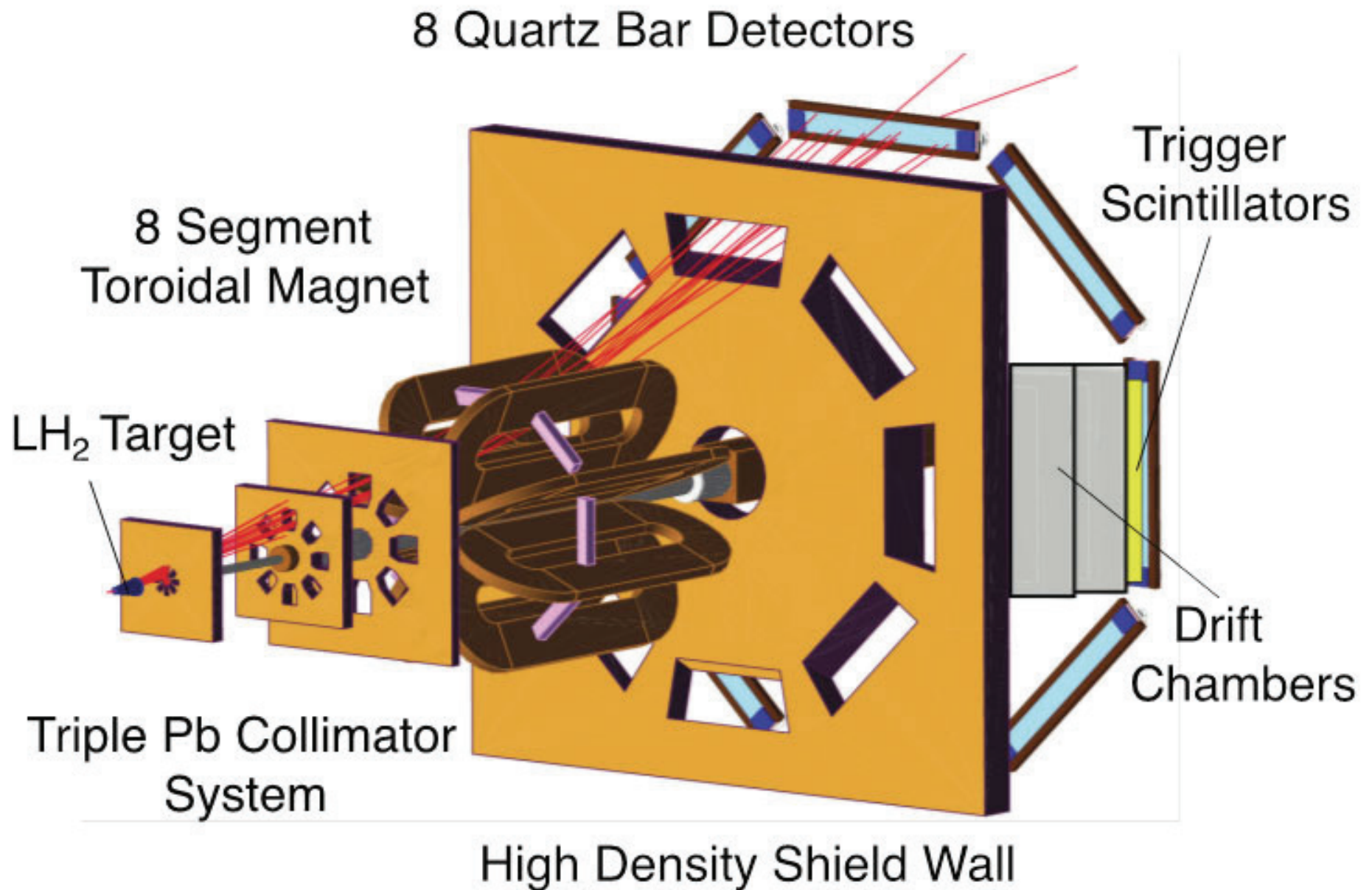
- Not charged: Photons not a good probe
- Use parity violating electron scattering:  
Proton weak charge is almost zero -  
see mostly neutrons

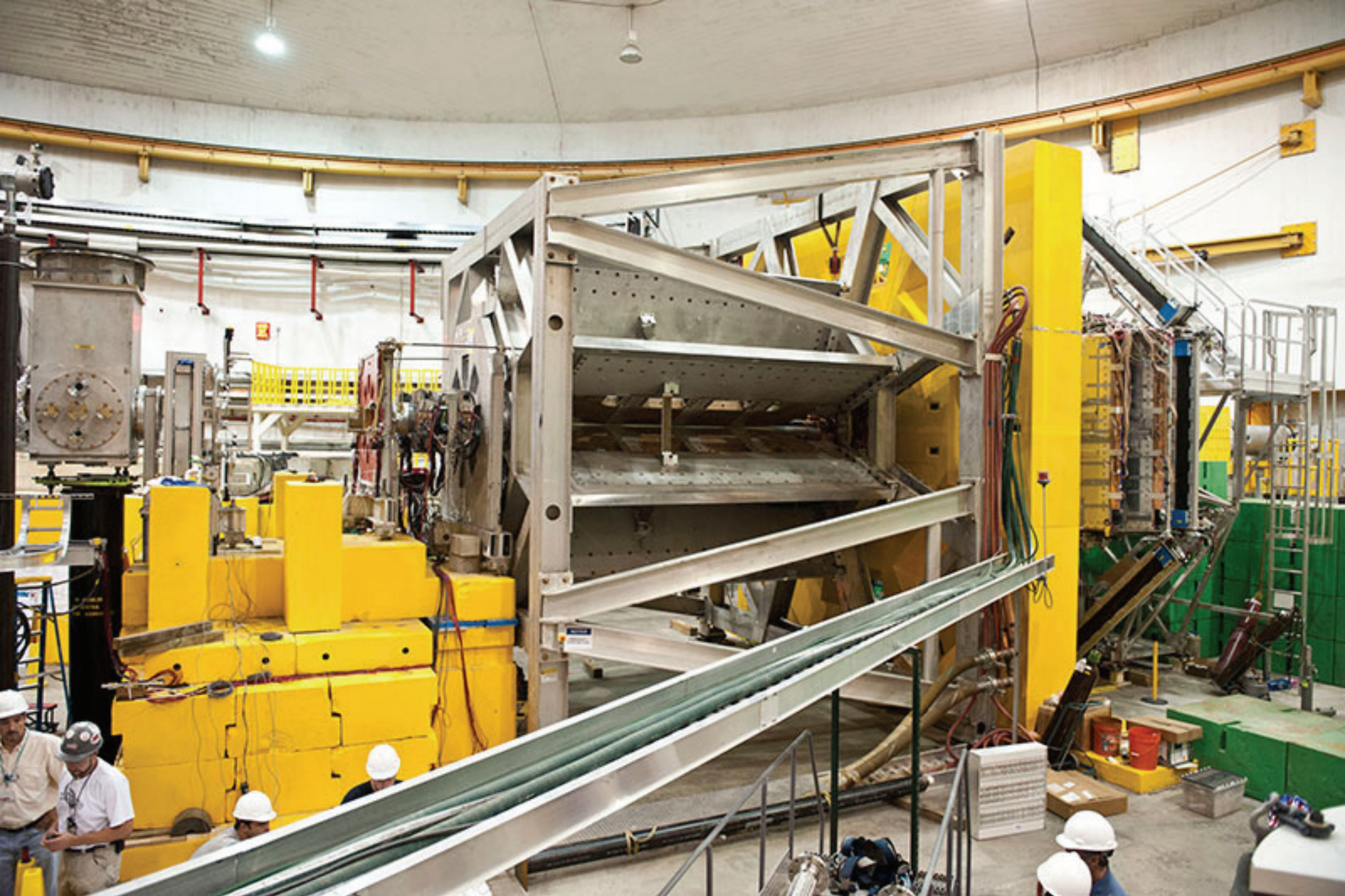


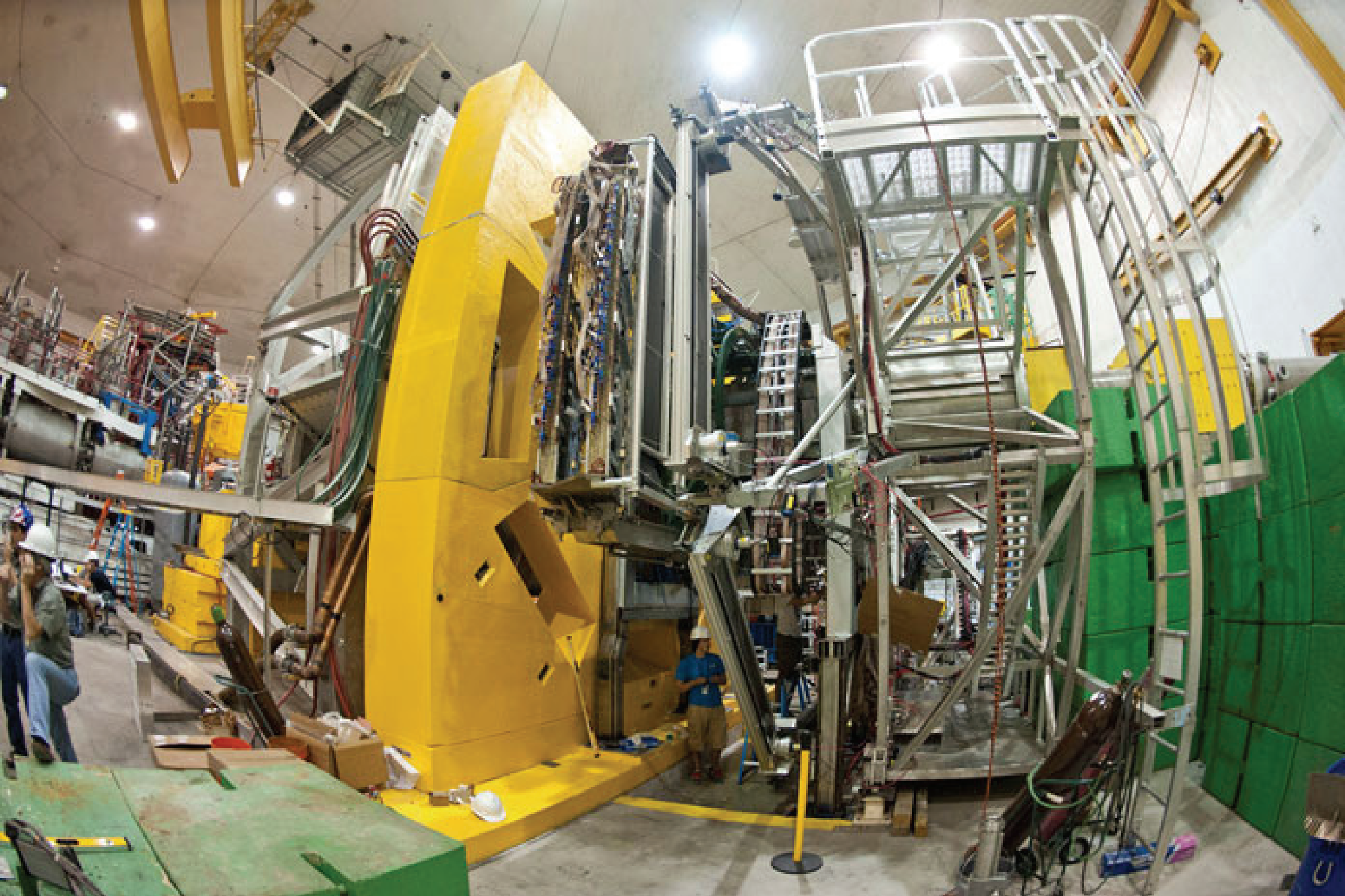
$$A_{PV} = \frac{G_F Q^2}{2\pi\alpha\sqrt{2}} \underbrace{\left( 1 - 4 \sin^2 \theta_W - \frac{F_n(Q^2)}{F_p(Q^2)} \right)}_{\approx 0}$$

# Other parity violation experiments

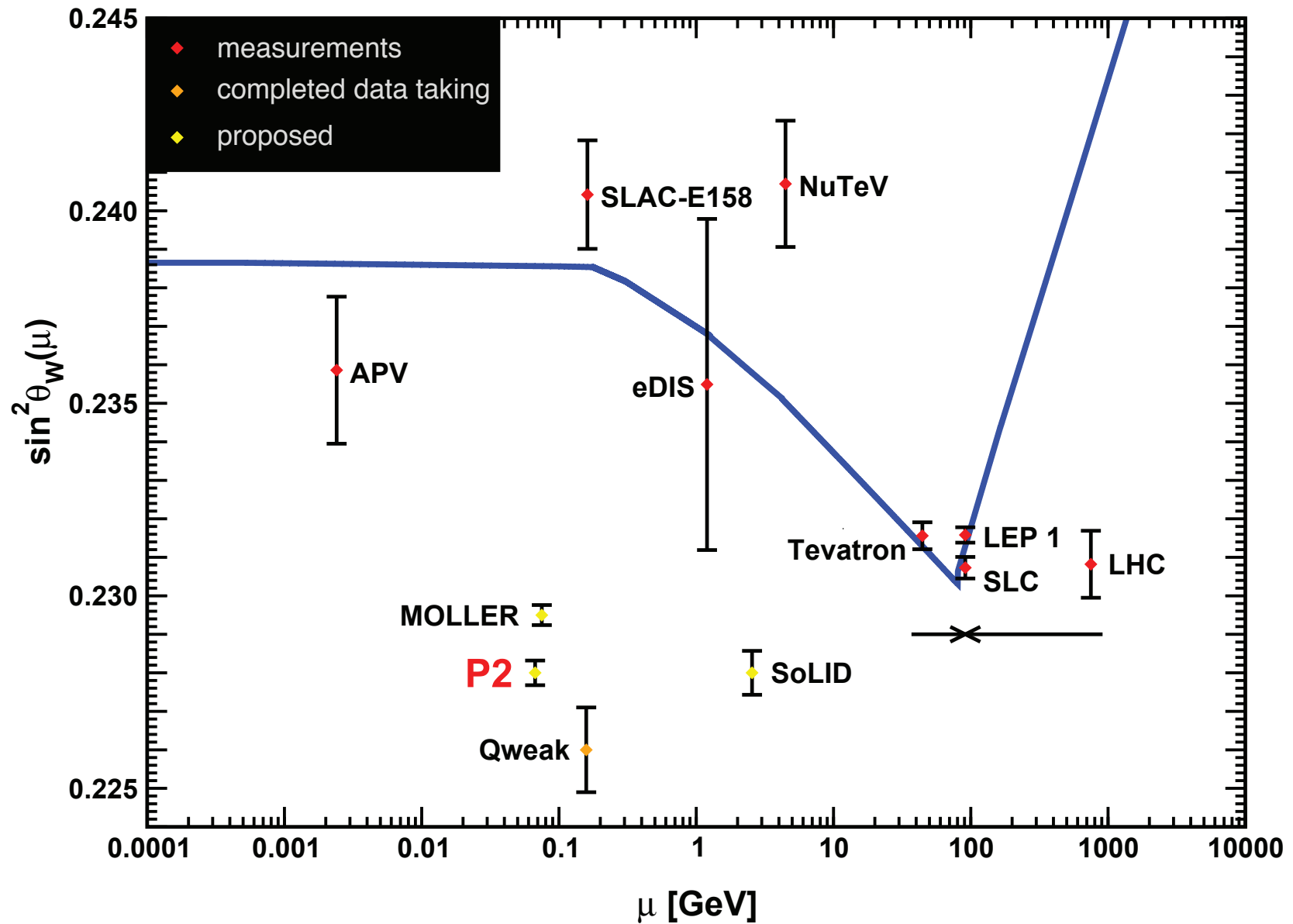
# QWeak (JLab)



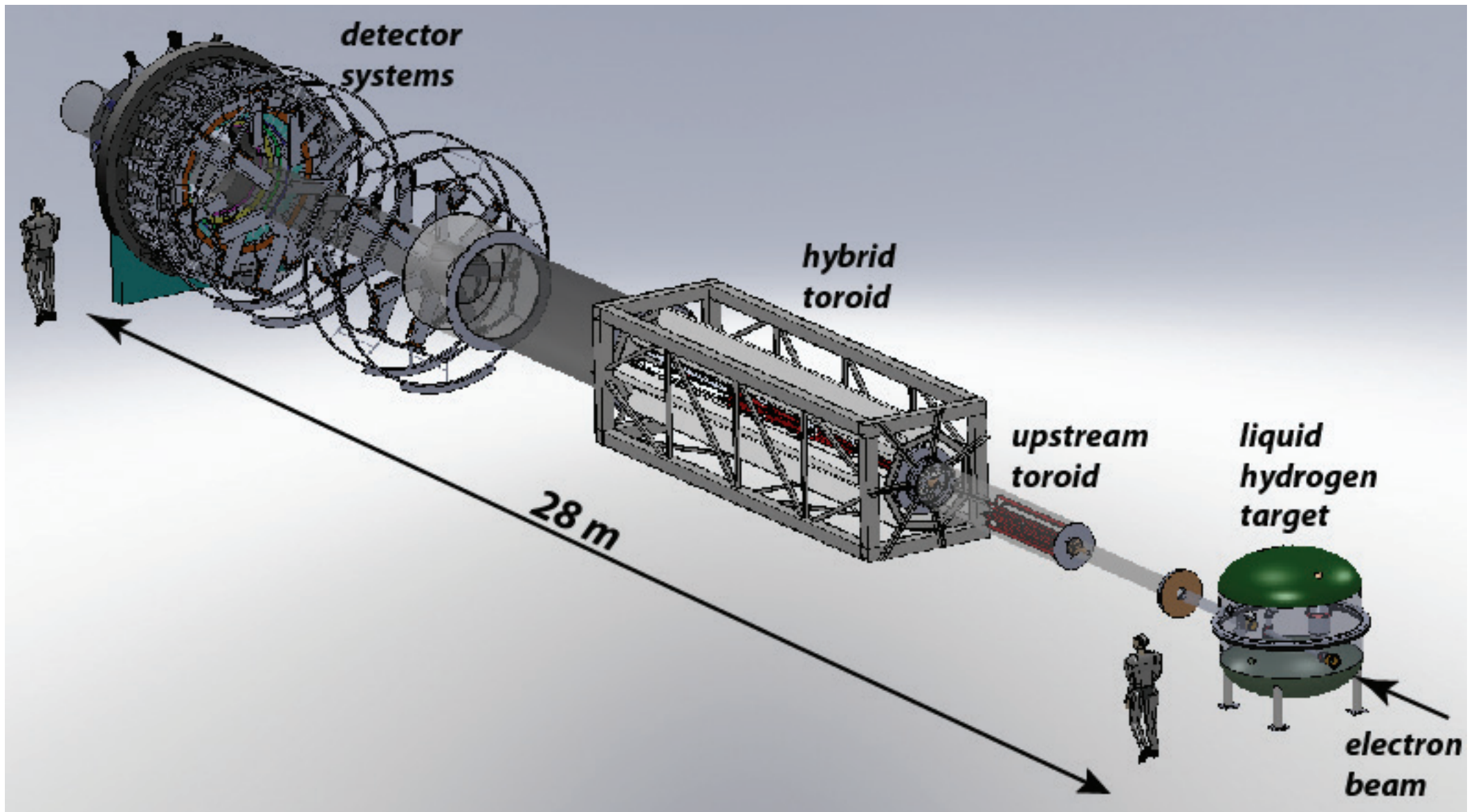




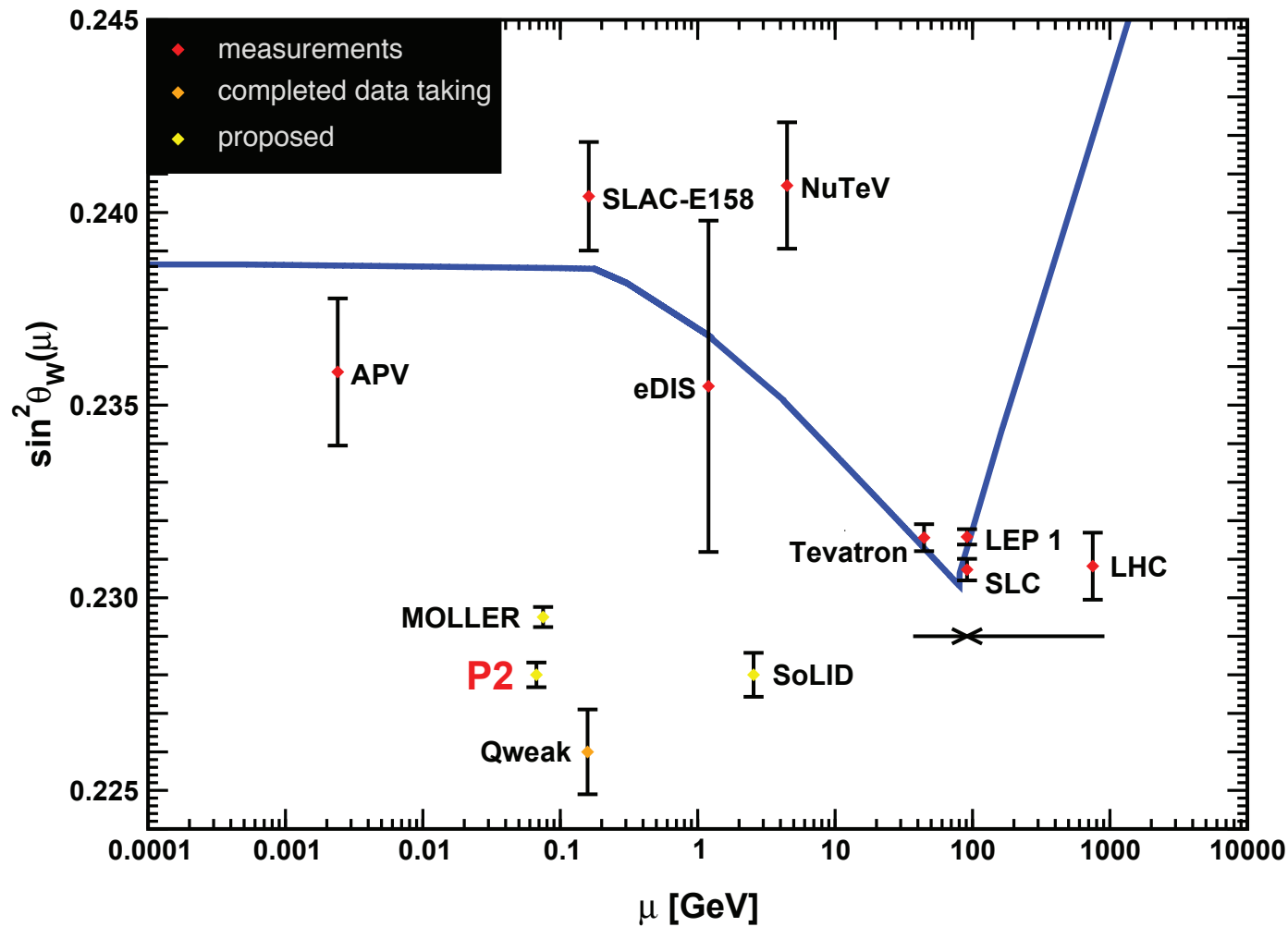
# QWeak results (not)



# Moller: $e^-e^-$ scattering at JLab



# More to come...



- Atomic parity violation in a single radium ion (Groningen)
- SoLID: Deep inelastic e-p scattering at JLab
- Much improved LHC measurements at the Z-pole



End of the preceding part